

Ultra-high frequency gravitational waves from inflaton decay

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Based on PLB 2211.02070 (A. Koshelev, A. Starobinsky, AT)
and ongoing work

Realization of inflation and reheating

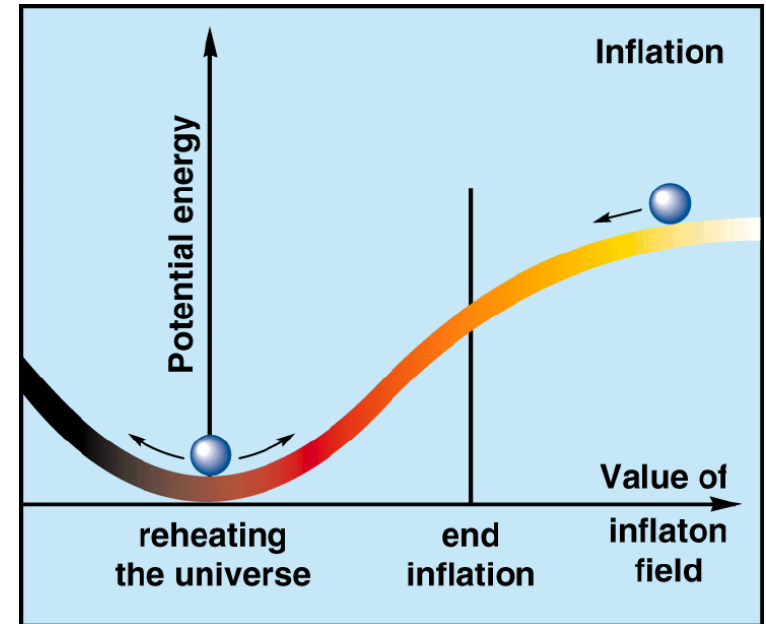
$$p = -\rho. \quad a(t) = \text{const} \cdot e^{H_{vac} t},$$

$$S = \int d^4x \sqrt{-g} \left(\frac{M_P^2}{2} R - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) \right)$$

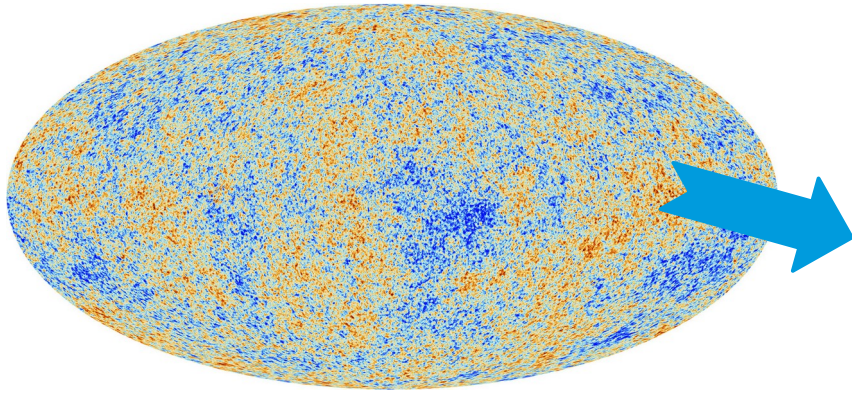
$$\rho = \frac{1}{2} \dot{\phi}^2 + V(\phi), \quad \text{Slowly rolling scalar field is a solution!}$$

$$p = \frac{1}{2} \dot{\phi}^2 - V(\phi).$$

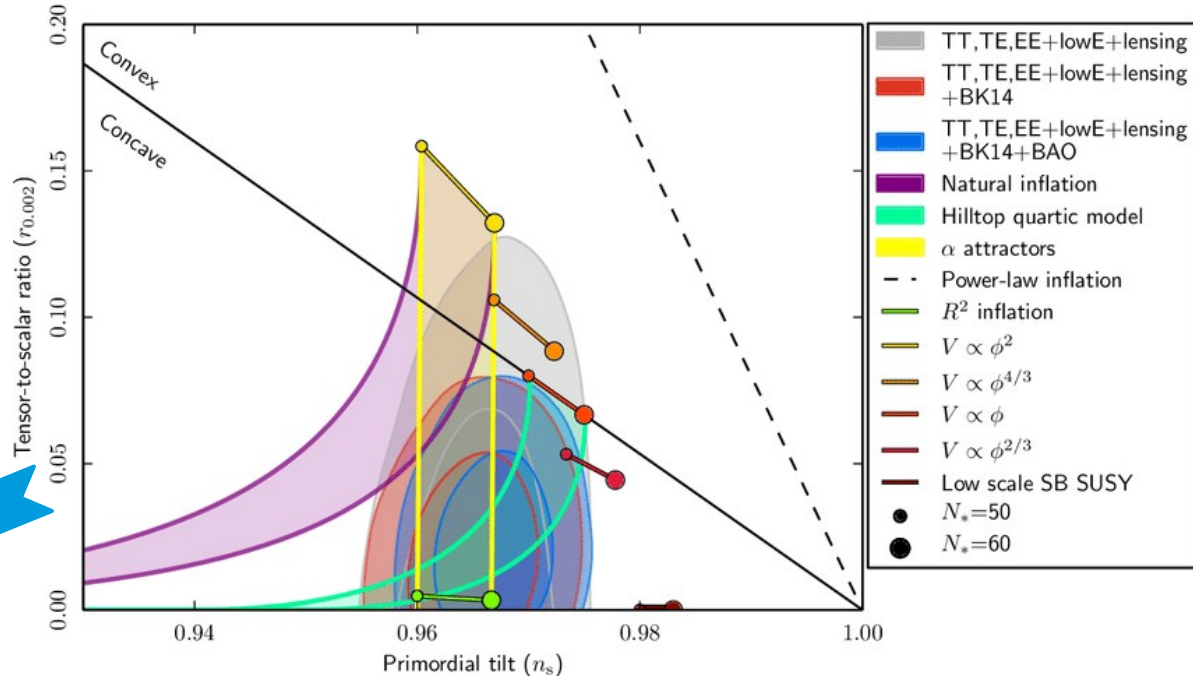
Oscillations after inflation decay to the SM particles \Rightarrow reheating of the Universe



Planck Constraints on the Potential



Serious QFT challenge – there is no renormalizable model left!



EFT of inflaton and gravity

Expansion around the flat space:

$$S = \int d^4x \sqrt{-g} \left(\frac{M_P^2}{2} R - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) \right)$$

$$S_{NR} = \int d^4x \sqrt{-g} \left(\frac{\phi}{\Lambda_1} R_{\mu\nu\lambda\rho} R^{\mu\nu\lambda\rho} + \frac{\phi}{\Lambda_2} R_{\mu\nu} R^{\mu\nu} + \frac{\phi}{\Lambda_3} R^2 + \frac{1}{\Lambda_4^2} G_{\mu\nu} \partial^\mu \phi \partial^\nu \phi \right)$$

$$S_{int}^{SM} = \int d^4x \sqrt{-g} \left(-|D_\mu H|^2 + \mu \phi H^\dagger H + \frac{1}{\Lambda_5^2} G_{\mu\nu} D^\mu H^\dagger D^\nu H \right)$$

Leading contribution to graviton production after inflation?

EFT of inflaton and gravity

Expansion around the flat space:

$$S = \int d^4x \sqrt{-g} \left(\frac{M_P^2}{2} R - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) \right)$$

$\Gamma = \frac{m^7}{32\pi M_p^4 \Lambda_1^2}$

$$S_{NR} = \int d^4x \sqrt{-g} \left(\frac{\phi}{\Lambda_1} R_{\mu\nu\lambda\rho} R^{\mu\nu\lambda\rho} + \frac{\phi}{\Lambda_2} R_{\mu\nu} R^{\mu\nu} + \frac{\phi}{\Lambda_3} R^2 + \frac{1}{\Lambda_4^2} G_{\mu\nu} \partial^\mu \phi \partial^\nu \phi \right)$$

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Decay to gravitons

reheating
bremsstrahlung

Other operators are suppressed by higher powers of Λ s

Inflaton decay to gravitons: selected results

A. Koshelev, A. Starobinsky, AT, PLB, arXiv:2211.02070

- Planck-suppressed operators **do matter** for low T_{reh} !

$$T_{reh} \lesssim \frac{m^{7/2}}{M_P^{3/2} \Lambda_1}$$

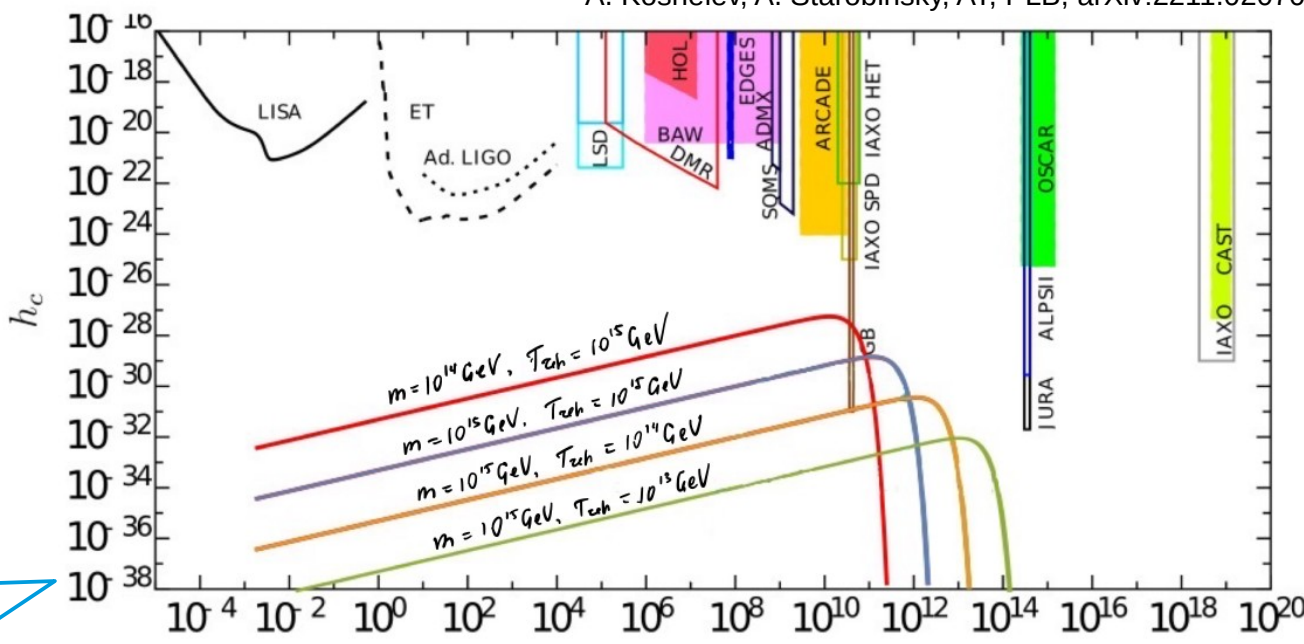
Overproduction of dark radiation

$$m = 10^{13} \text{ GeV} \quad T_{reh}^{min} = 1 \text{ GeV}$$

$$m = 10^{16} \text{ GeV} \quad T_{reh}^{min} = 10^{10} \text{ GeV}$$

- Larger inflaton mass - more HF GWs

$$\Lambda_1 = 10^8 \text{ GeV}$$



Detection proposals: G. Franciolini, A. Maharana, F. Muia, arxiv:2205.02153

$$\Delta N_{eff} \lesssim 0.2$$

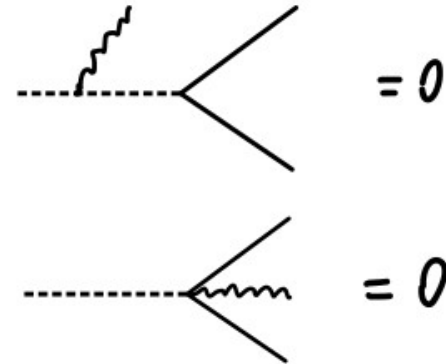
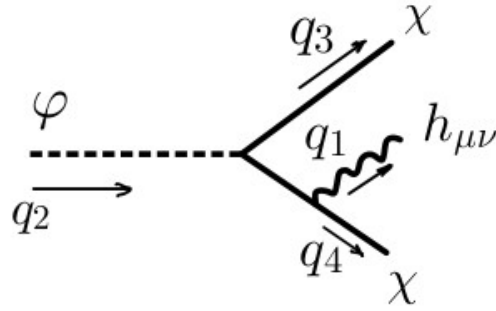
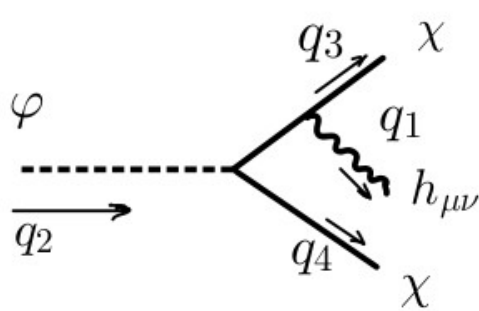
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$$\Delta N_{eff} = 2.85 \frac{\rho_{GW}}{\rho_{SM}} = 2.85 \frac{\Gamma_{GW}}{\Gamma_H}$$

$$\frac{d\Omega_{GW}}{d \log E} = \frac{16E^4}{M^4} \frac{\rho_{reh}}{\rho_0} \frac{\Gamma_{GW}}{H_{reh}} \frac{1}{\gamma(E)} e^{-\gamma(E)}$$

$$\gamma(E) = \left(\left(\frac{g_{reh}}{g_0} \right)^{1/3} \frac{T_{reh}}{T_0} \frac{2E}{M} \right)^{3/2}$$

Graviton bremsstrahlung during reheating



$$G(k) = \frac{\partial \Gamma}{\partial k} = A \frac{(m - 2k)^2}{m k}, \quad A = \frac{1}{64\pi^3} \frac{\mu^2}{3M_p^2} \left(1 + \frac{m^4}{\Lambda_5^4} \right)$$

$$\frac{d\rho_{GW}}{dk} = \int \frac{k dN}{a_0^3} = \int dt \frac{k n_\phi(t) a(t)^3}{a_0^3} G\left(k \frac{a_0}{a(t)}\right)$$

$$n_\phi = \frac{\rho_{reh}}{M} \left(\frac{a_{reh}}{a} \right)^3 e^{-\Gamma_{tot} t}$$

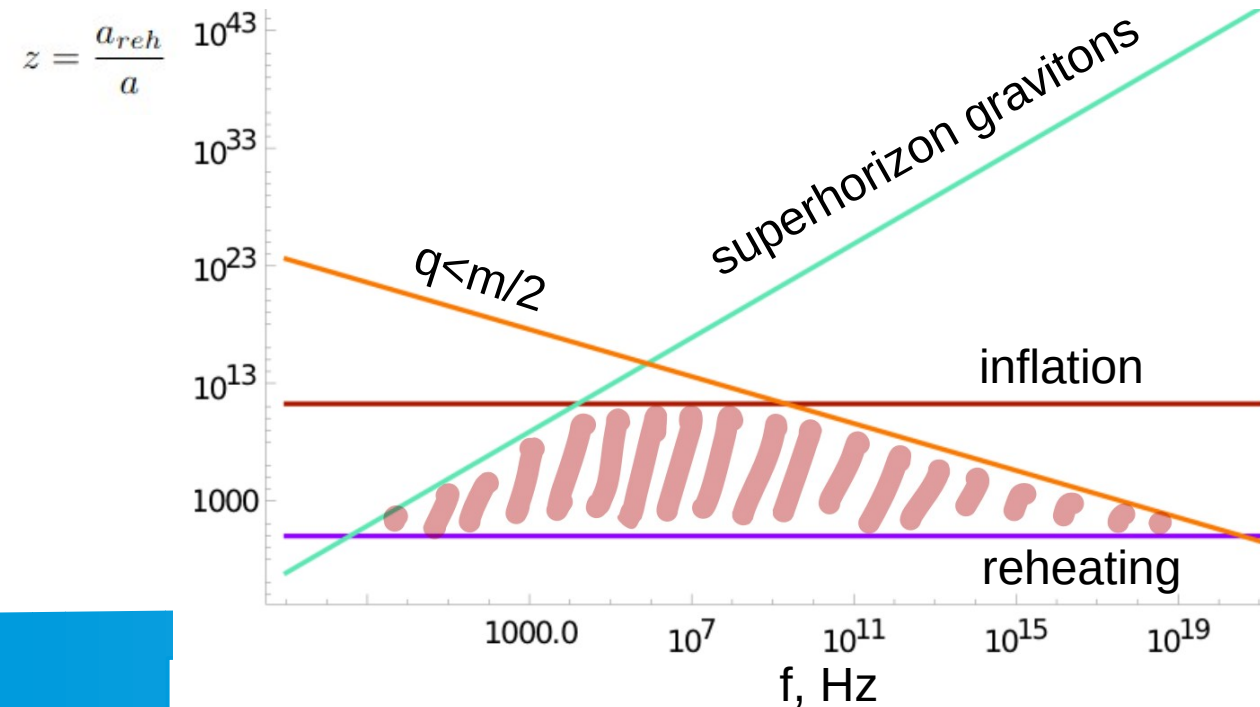
Not sensitive to inflaton-graviton coupling

Limits on GW frequencies

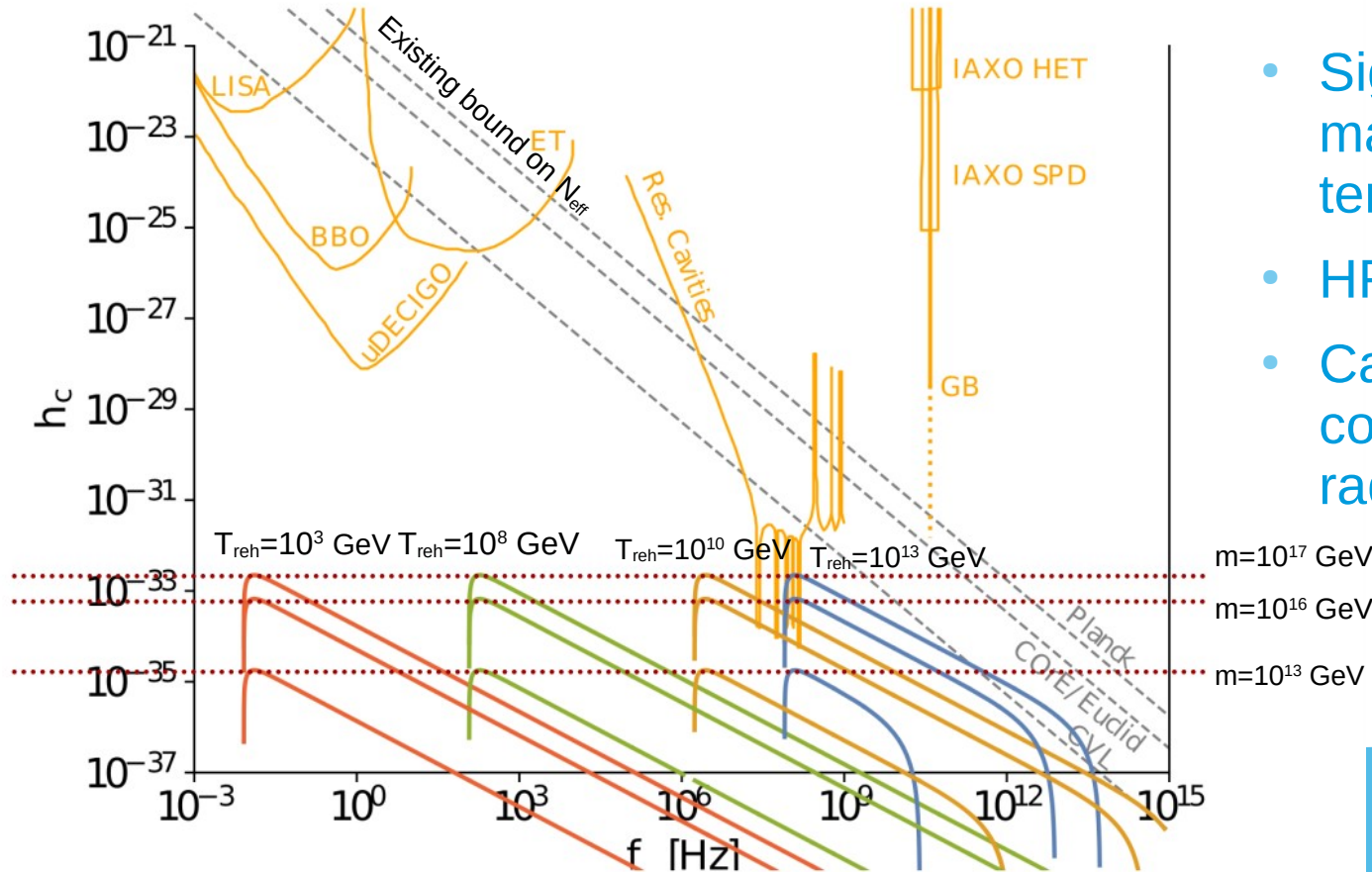
$$\frac{d\Omega_{GW}}{d \log k} = \frac{k^2}{M H_{reh}} \frac{a_{reh}^2}{a_0^2} \frac{\rho_{reh}}{\rho_0} \int_{z_{min}}^{z_{max}} dz G(kz \frac{a_0}{a_{reh}}) z^{-3/2} e^{-2z^{-3/2}/3}$$

Kinematic bound –
comoving momentum is
less than $m/2$

- Causality requirement - no superhorizon gravitons!
- Gravitons were emitted between inflation and reheating



Gravitational waves from bremsstrahlung: $\Lambda_5 = M_P$



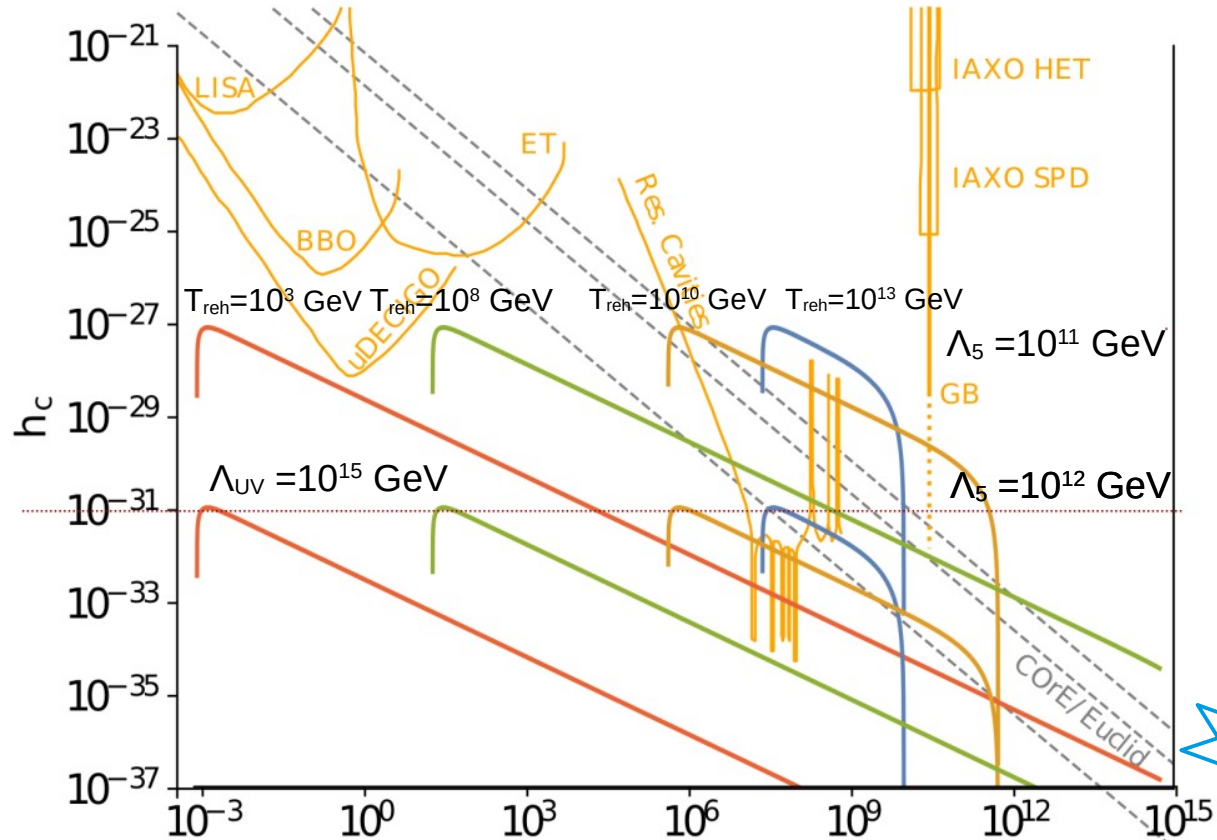
- Signals for different inflaton masses and reheating temperatures
- HF GW domain
- Can be also probed as contribution to the dark radiation

$$h_c(f) = \sqrt{\frac{3H_0^2}{\pi f^2} \frac{d\Omega_{GW}}{df}}$$

Detection prospects from Barman, Bernal, Xu, Zapata, 2301.11325

Results coincide with 2301.11325, except the IR cutoff

What if the quantum gravity scale is lower?



- GW signals for inflaton mass $m=10^{13}$ GeV
- The shape does not change, the amplitude is becoming higher
- The unitarity breaking scale is $\Lambda_{\text{UV}}=(\Lambda_5 M_{\text{P}})^{1/2} > m$
- From $\Lambda_{\text{UV}}=10^{15}$ GeV – tension with N_{eff} bound

Reheating-dependent bounds on quantum gravity scale!

Conclusions

- High frequency gravitational waves can be sensitive to the quantum gravity effects
- Perturbative decay of inflation to gravitons can be non-negligible for low reheating temperatures → high frequency GWs
- Graviton bremsstrahlung during reheating can provide a sizable HF GW signal → constraints on EFT
- Reheating-dependent constraints on quantum gravity scale from gravitational waves !

Thank you!