

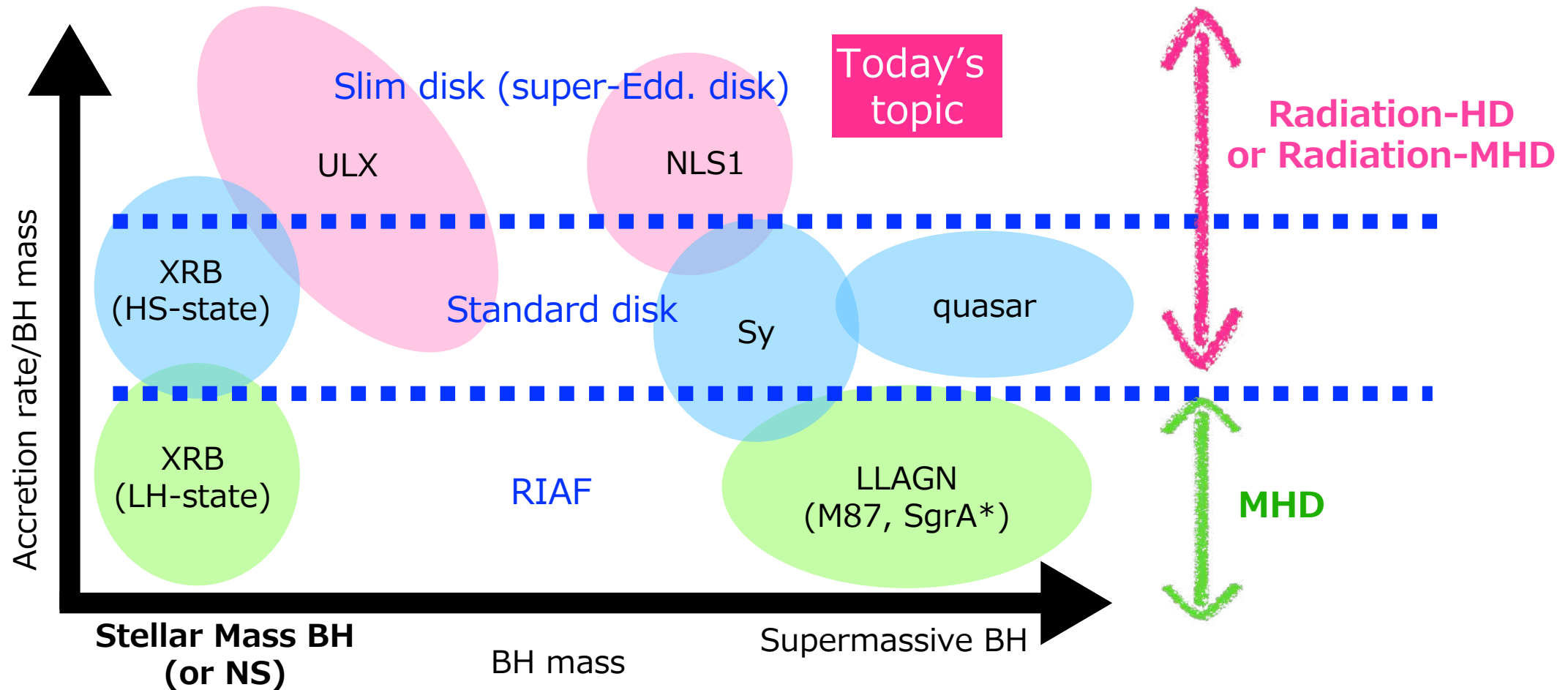
32<sup>nd</sup> Texas Symposium on Relativistic Astrophysics

# Numerical simulations of super-Eddington accretion flows

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# Three Accretion Modes



# Development of research of Super-Edd. flows

1988~

## **1D approach**

Slim disk model has been established ([Abramowicz et al. 1988](#))

2005~

## **Radiation-HD sim.**

## **Radiation-MHD sim.**

### **Multi-dimensional Simulations**

Quasi steady inflow-outflow structure has been revealed.

([Ohsuga et al. 2005, 2009](#), [Ohsuga & Mineshige 2011](#), [Jiang et al. 2014, 2019](#))

2014~

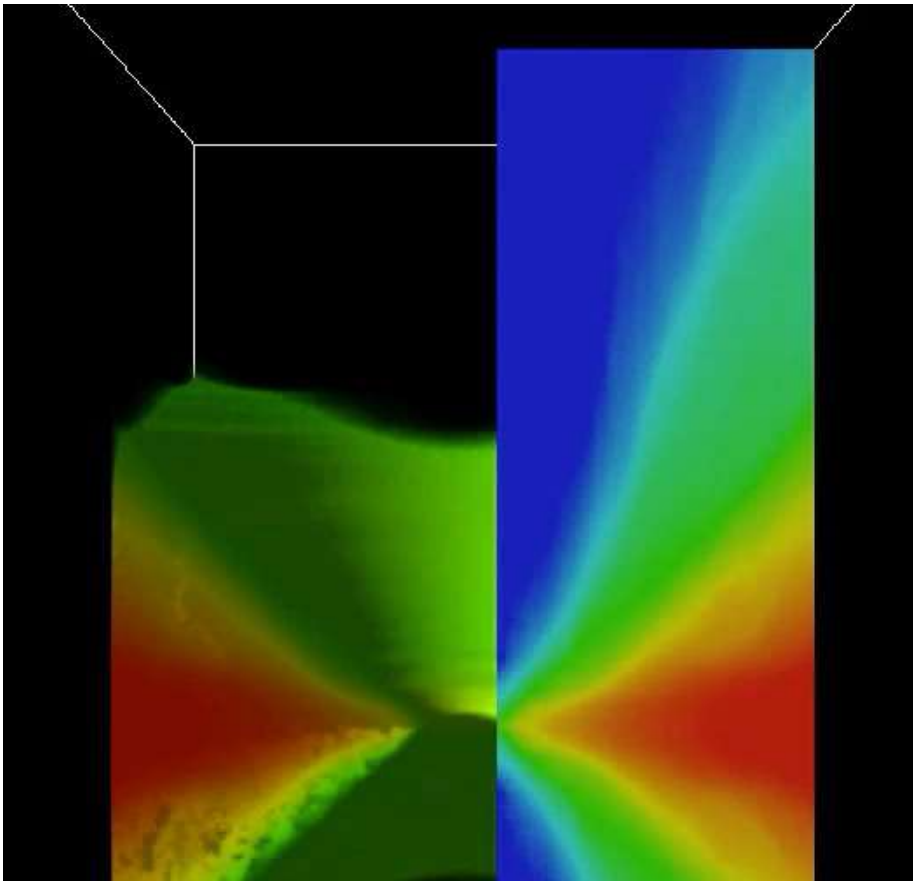
## **General Relativistic Radiation-MHD Sim.**

General relativistic effects (e.g., BZ effect, LT precession) has been studied

([Sadowski et al. 2014, 2016](#), [Takahashi, Ohsuga et al. 2016](#), [Utsumi, Ohsuga et al. 2022](#), [Brandon 2023](#), [Asahina & Ohsuga submitted](#)).

# Radiation-HD/MHD simulations Super-Edd. Flows

Ohsuga et al. 2009



## Setup

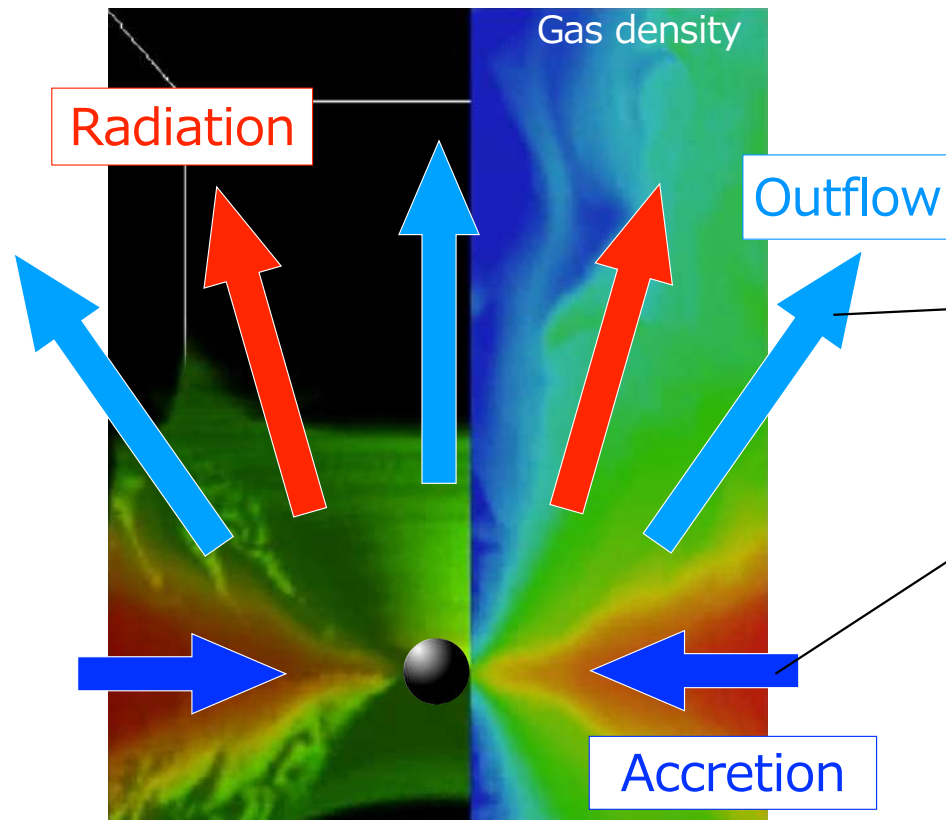
- BH mass:  $10M_{\text{sun}}$
- Initial condition: equilibrium torus with embedded poloidal magnetic field (plasma-beta=100)

## Quasi-steady structure

- The super-Eddington disks ( $\dot{M} \sim \text{a few } 100L_{\text{Edd}}/c^2$ ,  $L_{\text{disk}} \gg L_{\text{Edd}}$ )
- Radiatively-driven outflows

see also Ohsuga et al. 2005; 2011, Ohsuga 2007,  
Jiang et al. 2014, 2019

# Why is super-Eddington accretion feasible?



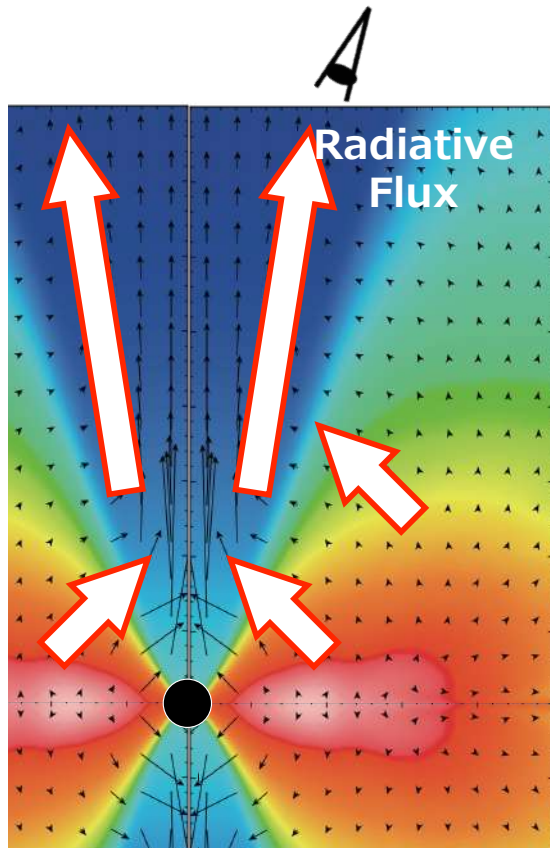
## Radiatively driven outflows:

Strong radiation pressure supports the thick disk and generates the outflows above the disk.

## Accretion:

Photons mainly escape through the less-dense region above the disk. The radiation pressure cannot prevent the accreting motion within the disk.

# Apparent Luminosity



Ohsuga, Mineshige 2011

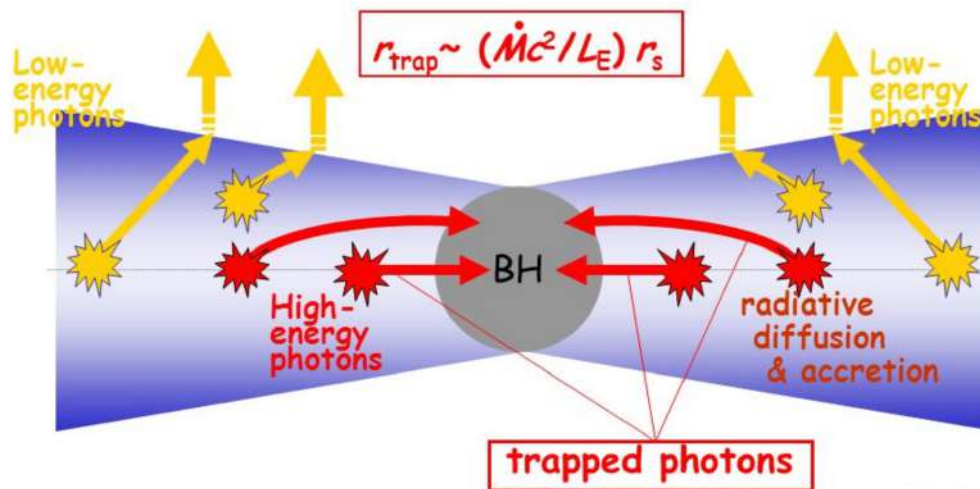
- The radiative flux is mildly collimated since the disk is optically and geometrically thick.
- Thus, observed luminosity is very sensitive to the observer's viewing angle.
- The apparent luminosity becomes highly super-Eddington for the face-on observers.

ex:  $22L_{\text{Edd}}$  for  $\lesssim 20^\circ$   
when  $\dot{M} \sim 100L_{\text{Edd}}/c^2$  &  $L_{\text{disk}} \sim 3L_{\text{Edd}}$ .

**Large luminosity of ULXs ( $>10^{39-40}\text{erg/s}$ ) can be explained for the face-on case.**

# Photon trapping

Ohsug et al. 2002, 20023



## Photon trapping:

Huge amount of radiation energy is swallowed by the black hole with accreting matter.

Radiation diffusion timescale

> Matter accretion timescale

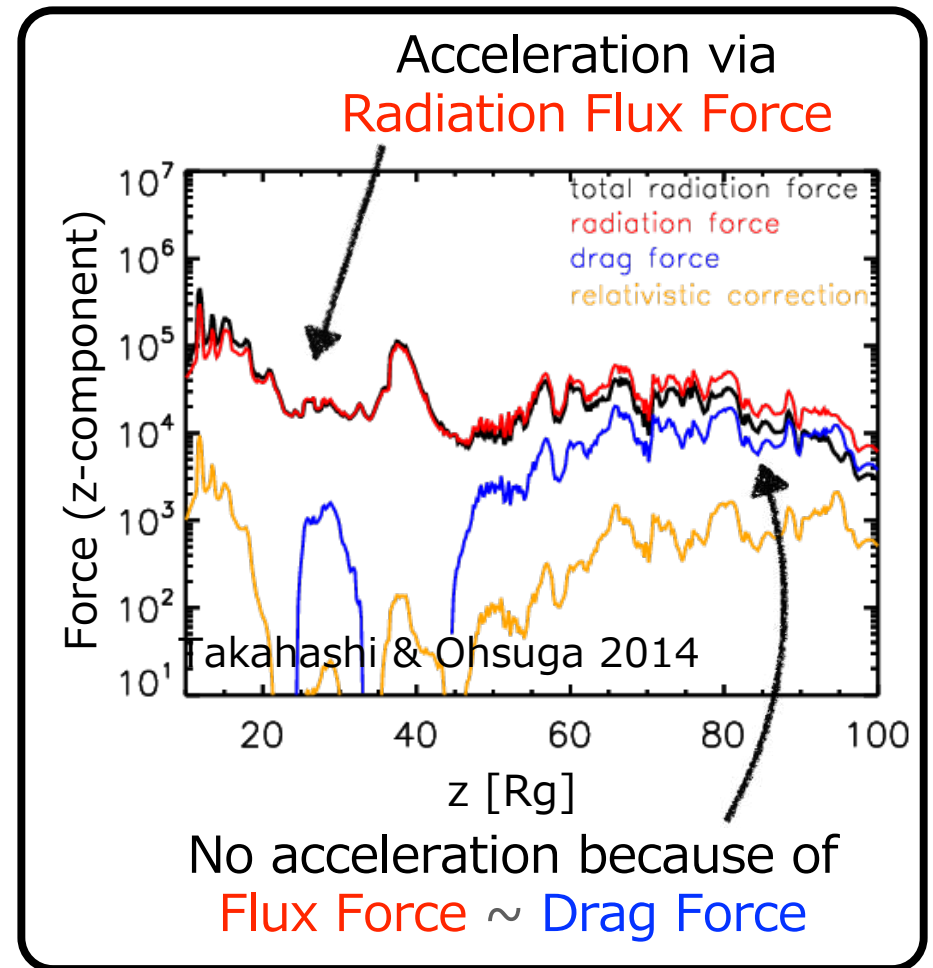
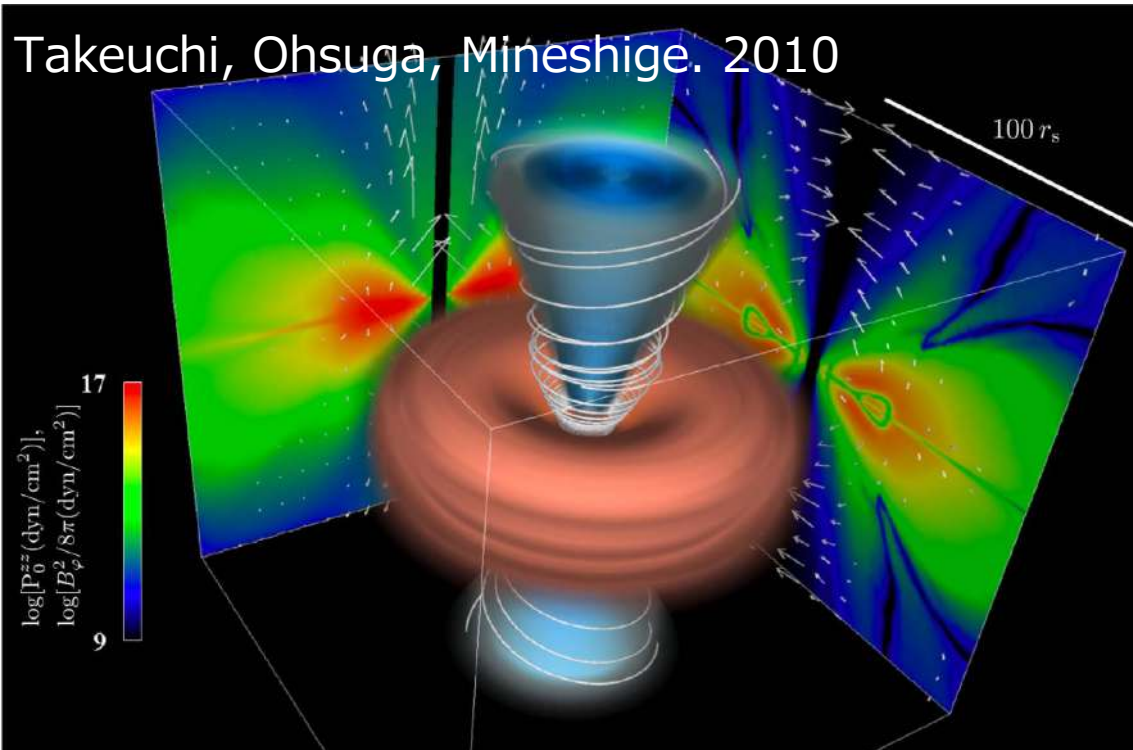
$$\rightarrow r_{\text{trap}} < \frac{\dot{M} c^2}{L_{\text{Edd}}} r_s$$

Photon-trapping occur  
in super-Eddington accretion flows.

# Radiatively-driven Jets

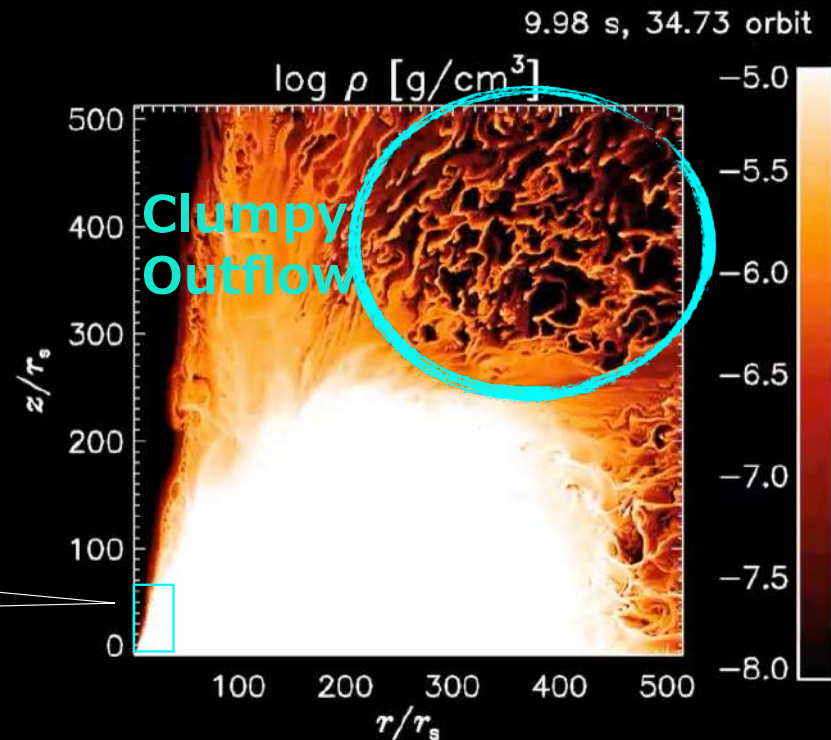
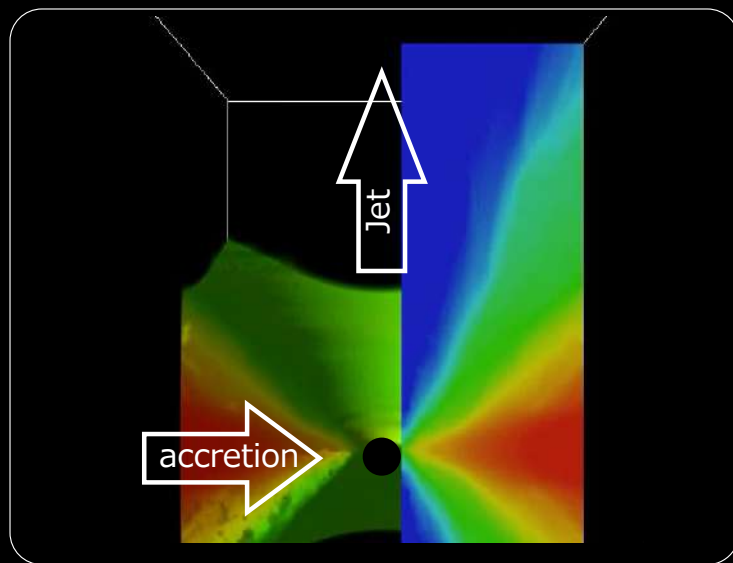
Resulting jet velocity ( $\sim 0.3-0.5c$ ) is roughly consistent with the jets in SS433.

Takeuchi, Ohsuga, Mineshige. 2010



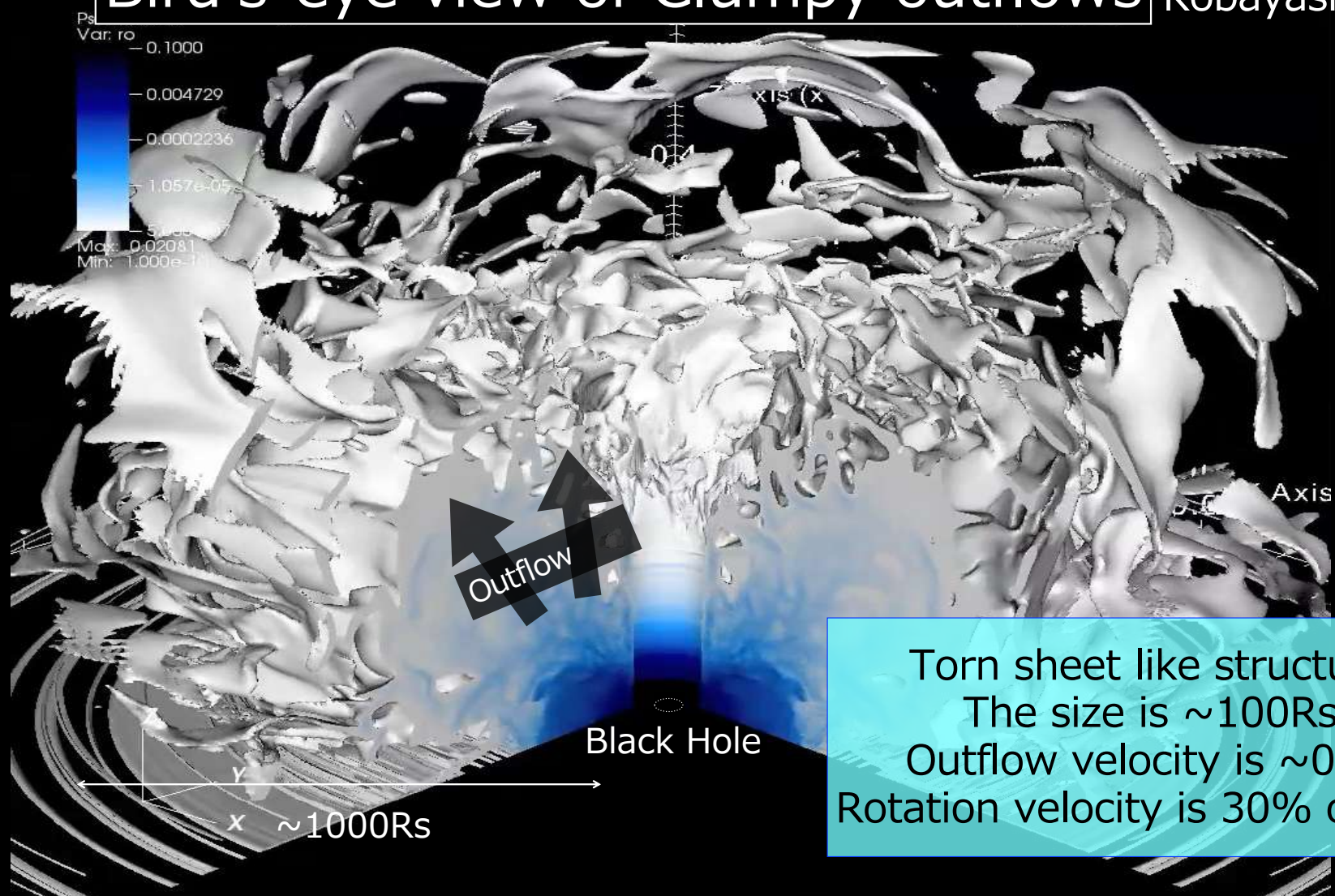
# Clumpy Outflows

Takeuchi, Ohsuga, Mineshige 2013



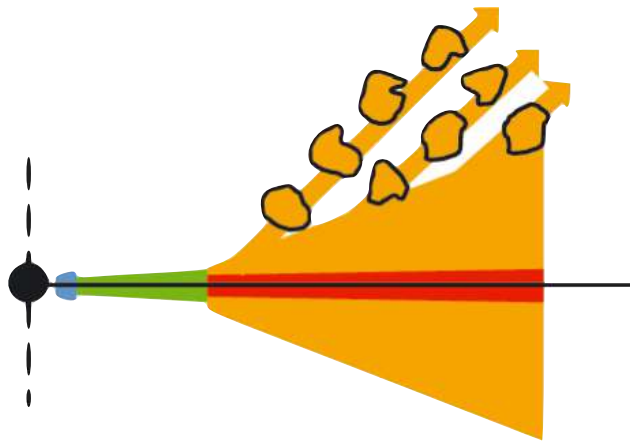
Clumpy outflows: Wind outflows fragment into many gas clouds

# Bird's-eye view of Clumpy outflows Kobayashi+18



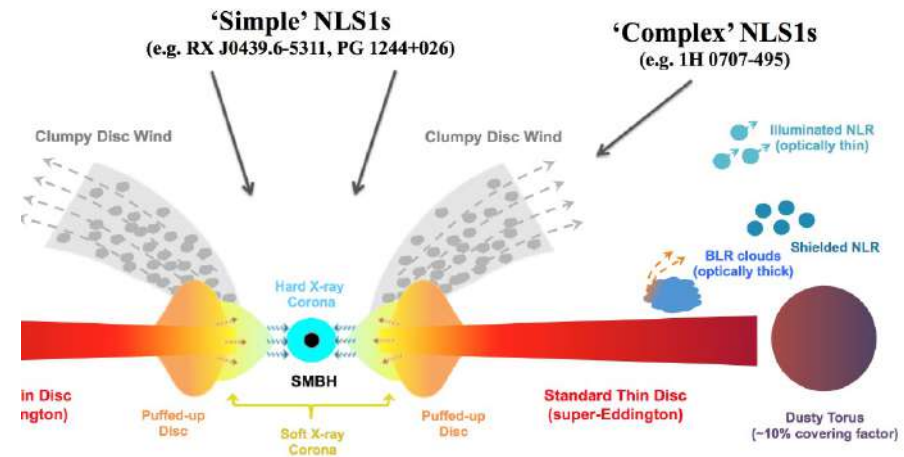
# Observations of Clumpy outflows

Some ULXs exhibit the time variations of X-ray luminosity, implying the launching of clumpy outflows.



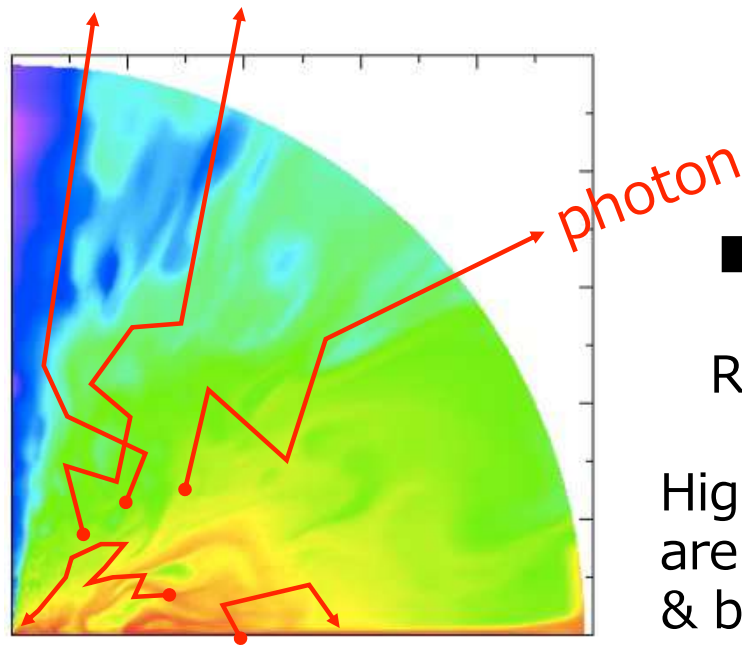
Middleton+11

Launching of clumpy winds is also reported by observations of NLS1s or V404 Cyg.



Jin+17 see also Motta+17

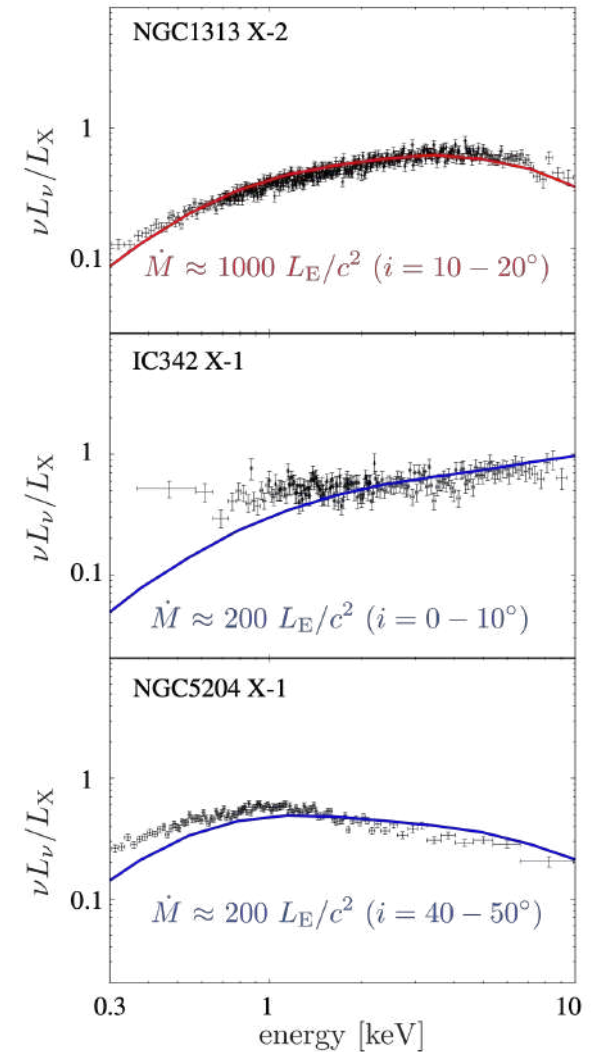
# Comparison with ULXs



Monte Carlo  
Radiation Transfer:

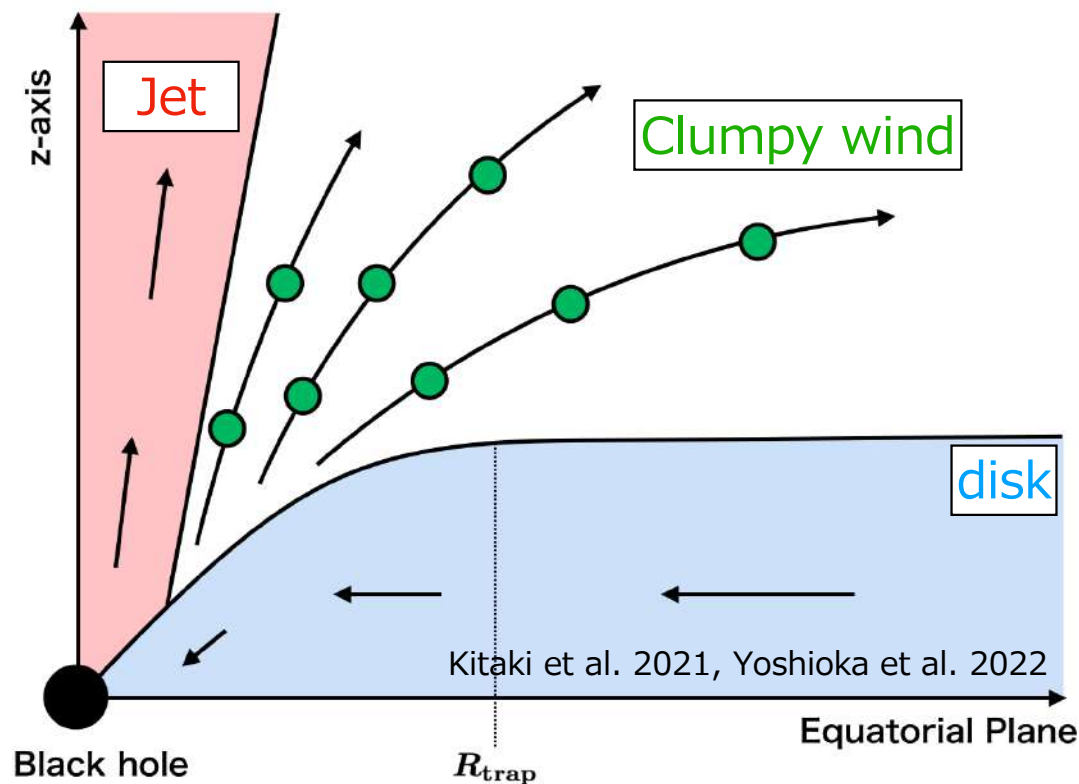
High-energy X-ray photons  
are produced by the thermal  
& bulk Comptonization.  
Simulated spectra nicely fit  
the observations of ULXs.

Kawashima et al. 2012  
(data; Gladstone 2009)



# Overall structure of the super-Edd. flows

Schematic picture of the overall structure

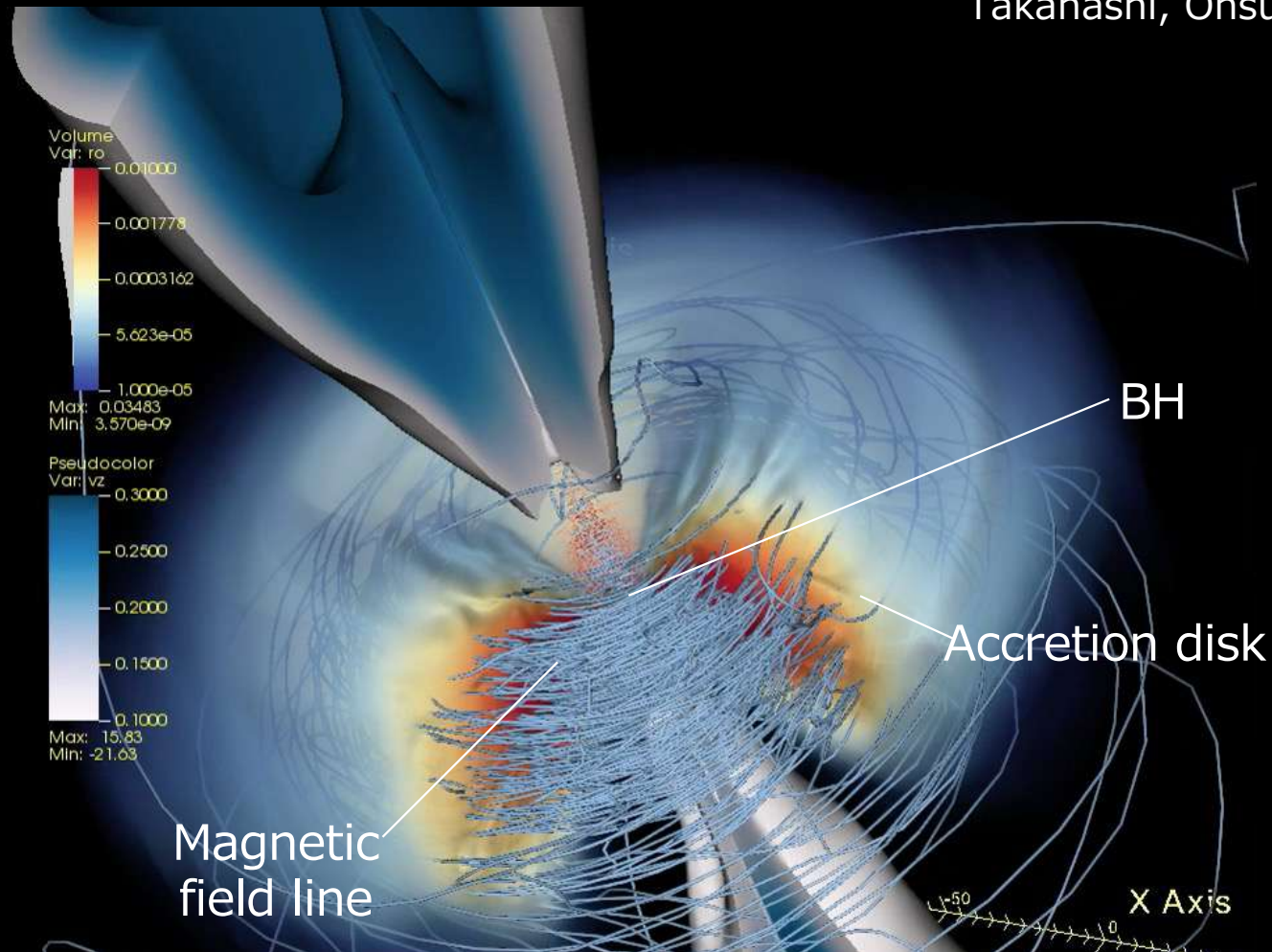


Super-Eddington flows consist of three components;

- radiation pressure-dominated **disk**
- radiatively-driven high-velocity outflow around the rotation axis (**jet**)
- radiatively-driven **clumpy wind**.

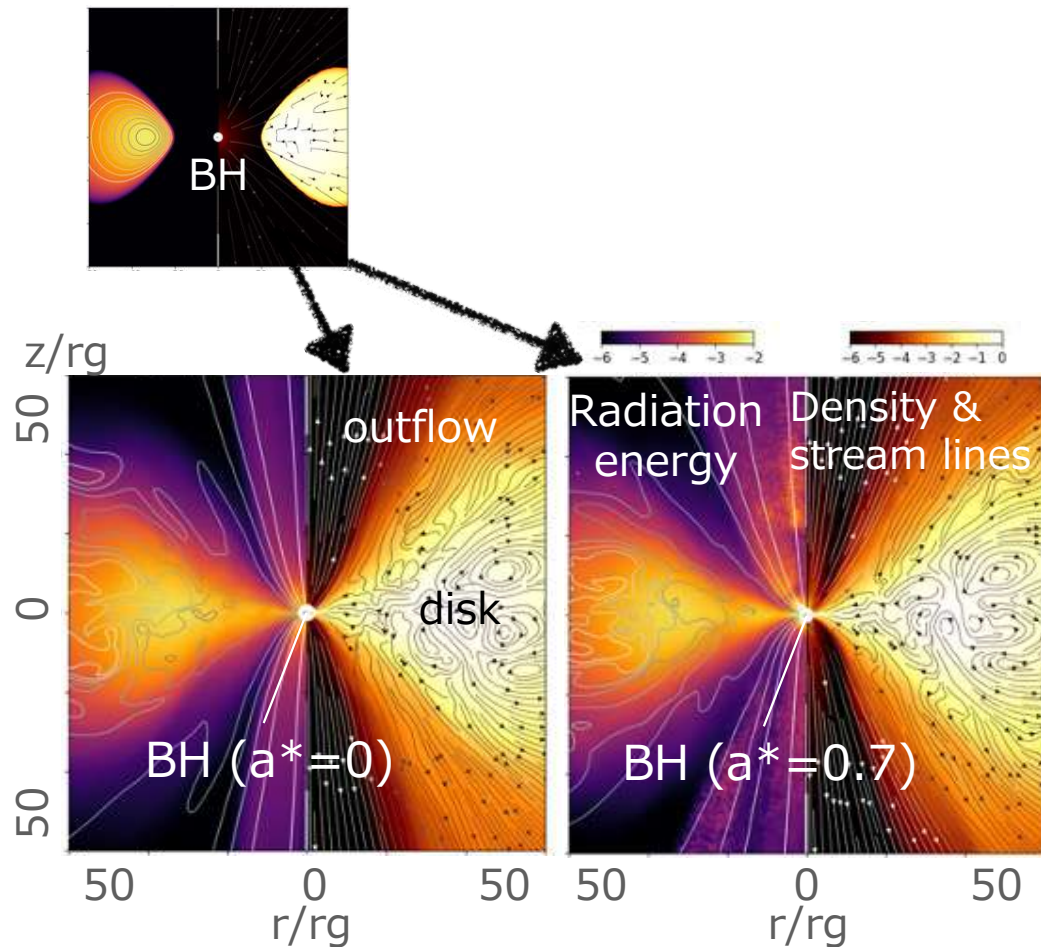
# General Relativistic Radiation-MHD simulations

Takahashi, Ohsuga et al. 2016



# BZ mechanism in Super-Edd. disk

Utsumi, Ohsuga et al. 2022



## Setup

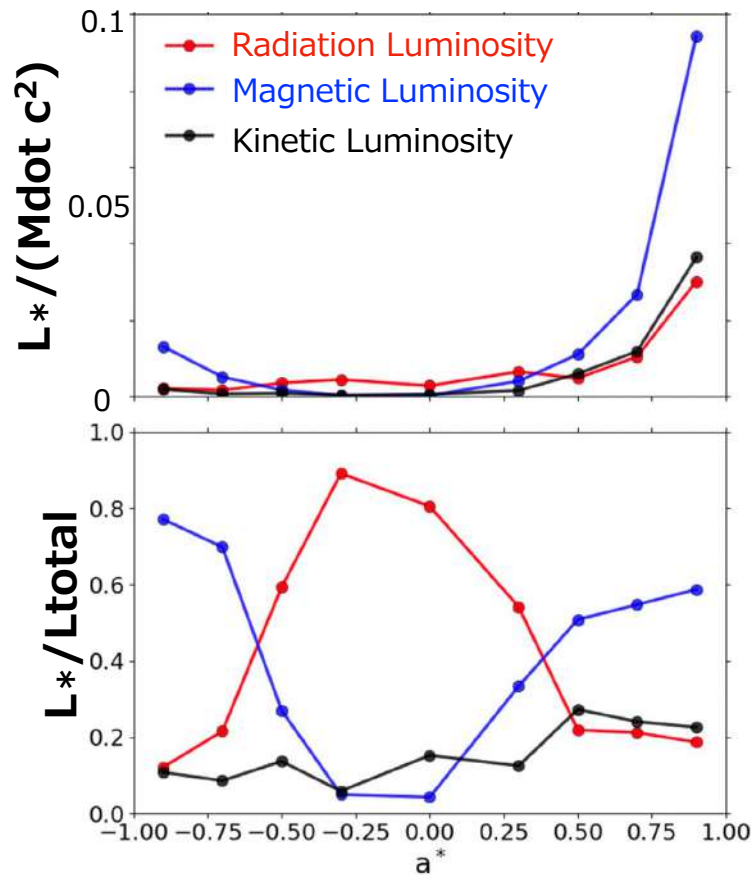
- BH mass:  $10M_{\text{sun}}$
- Initial condition: equilibrium torus with embedded poloidal magnetic field (plasma-beta=100)
- Spin parameter: -0.9, -0.7, -0.5, -0.3, 0, 0.3, 0.5, 0.7, 0.9

## Quasi-steady structure

- In all models, the super-Eddington disks ( $\dot{M} \sim \text{a few } 100L_{\text{Edd}}/c^2$ ) and strong outflows are formed.

\* Magnetic field is not so strong (SANE)

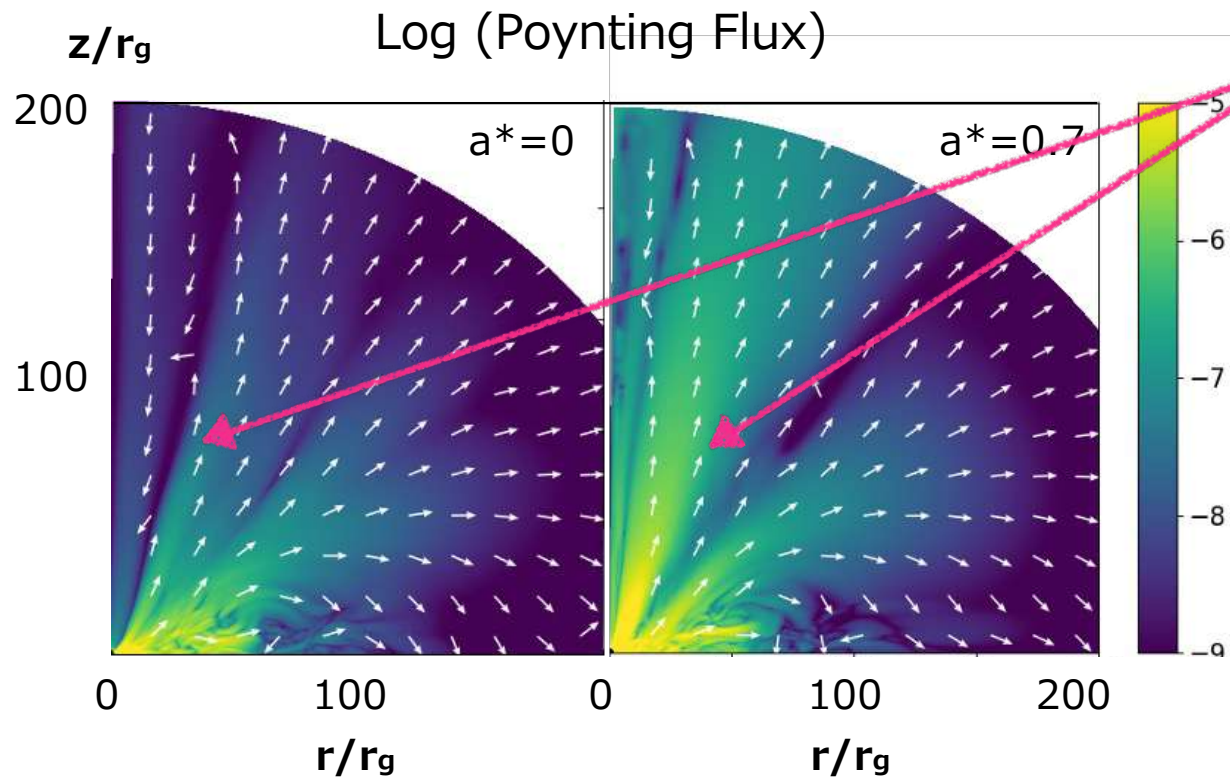
# Energy Conversion Efficiency



For the case of  $a^* \sim 0$ , energy is mainly released by the radiation. When  $|a^*|$  is large, the energy released by the Poynting flux (Magnetic Luminosity) exceeds the Radiation Luminosity.

Radiation luminosity accounts for 80% when  $a^* \sim 0$ . But the magnetic luminosity is three times larger than the radiation luminosity for the case of  $a^* > 0.5$ .

# Enhancement of Poynting Flux

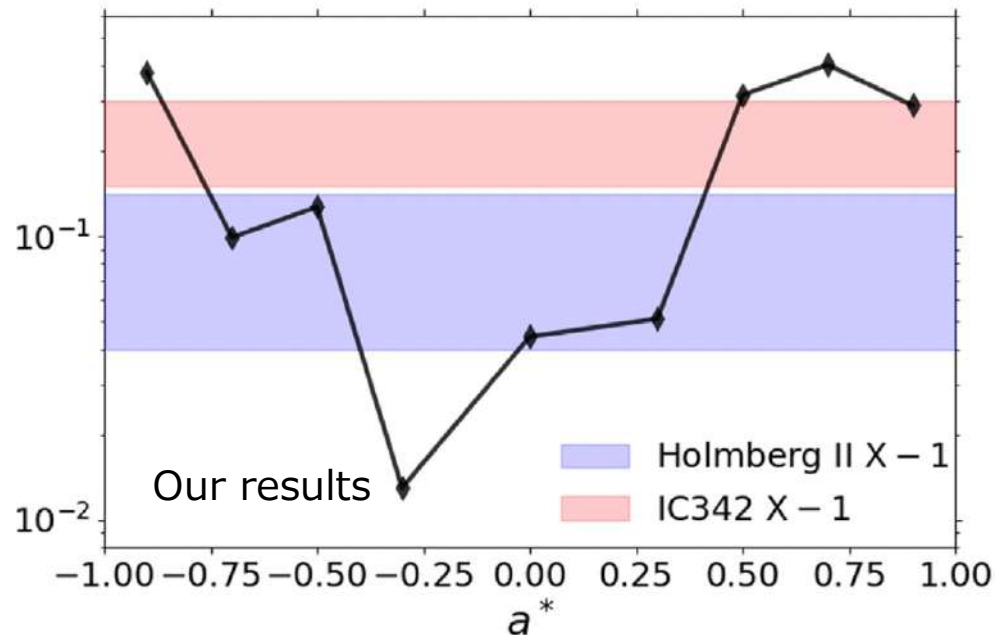


The Poynting flux around the rotation axis is stronger for larger  $|a^*|$ . This is probably caused by **Blandford-Znajek (BZ) effect**.

# Are black holes in ULXs rotating?

Kinetic Luminosity/Isotropic X-ray Luminosity

\*Isotropic X-ray Luminosity:  
Radiation luminosity observed  
by face-on observer.

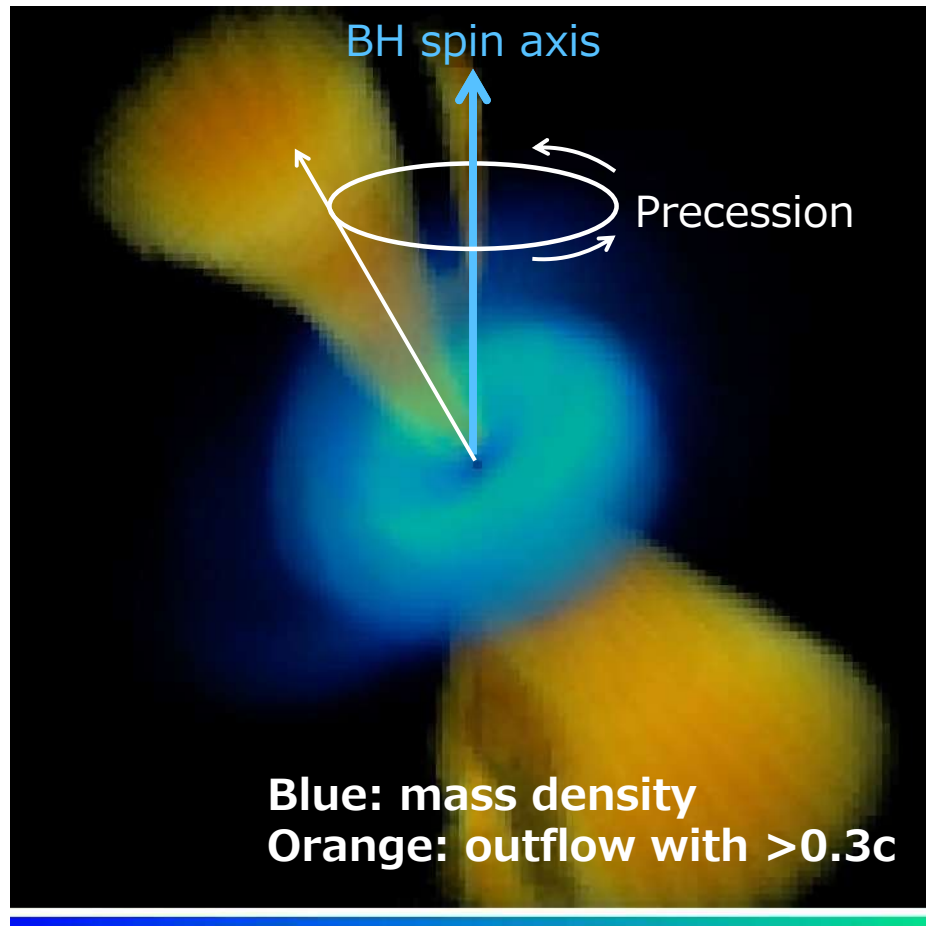


In our results, the ratio of the kinetic luminosity to isotropic X-ray luminosity tends to increase with  $|a^*|$ .

Thus, rapidly (slowly) rotating black hole probably exist in IC342 X-1 (Holmberg II X-1).

# Lense-Thirring Precession of Super-Edd. disk

Asahina & Ohsuga submitted



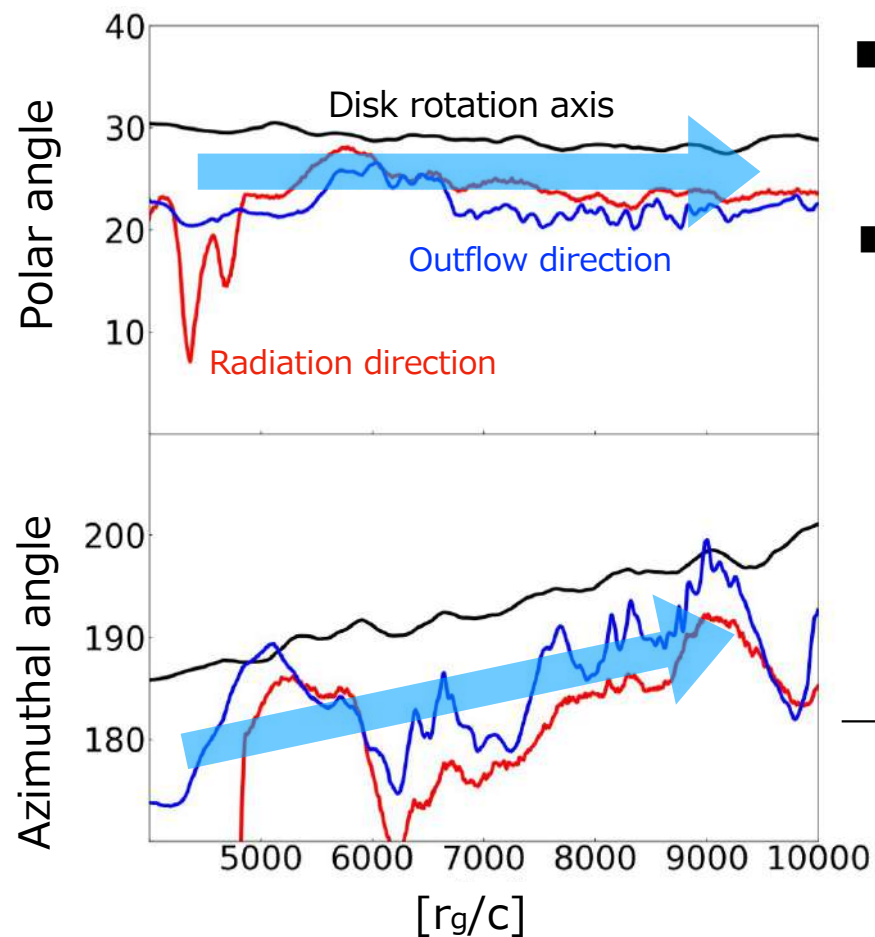
## Setup

- BH mass:  $10M_{\text{sun}}$
- Initial condition: equilibrium torus with embedded poloidal magnetic field (plasma-beta=100) tilted **30 degree**.
- Spin parameter: **0.9**

## Inflow-outflow structure

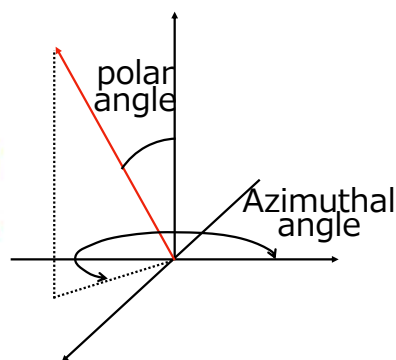
- The super-Eddington disk, which is tilted and twisted, forms.
- Strong outflows are also formed.
- Accretion rate: several  **$100 L_{\text{Edd}}/c^2$**
- Radiation Luminosity: **several  $L_{\text{Edd}}$**
- Kinetic Luminosity: **several  $L_{\text{Edd}}$**

# Precession of disk, outflow, radiation



- The super-Eddington disk exhibits the precession motion.

- The gas and radiation is mainly ejected around the rotation axis of the disk ( $\sim 30^\circ$ ), rather than around the spin axis of the BH ( $0^\circ$ ).



- The direction of outflow and radiation also changes according to the precession motion of the disk. However, there is about  $10^\circ$  delay from the rotational axis of the disk.

# Comparison with observations

[1] Quasi periodic oscillations of ULXs:

The typical timescale of the precession is  $\sim 9$ sec for the case of stellar mass BH.

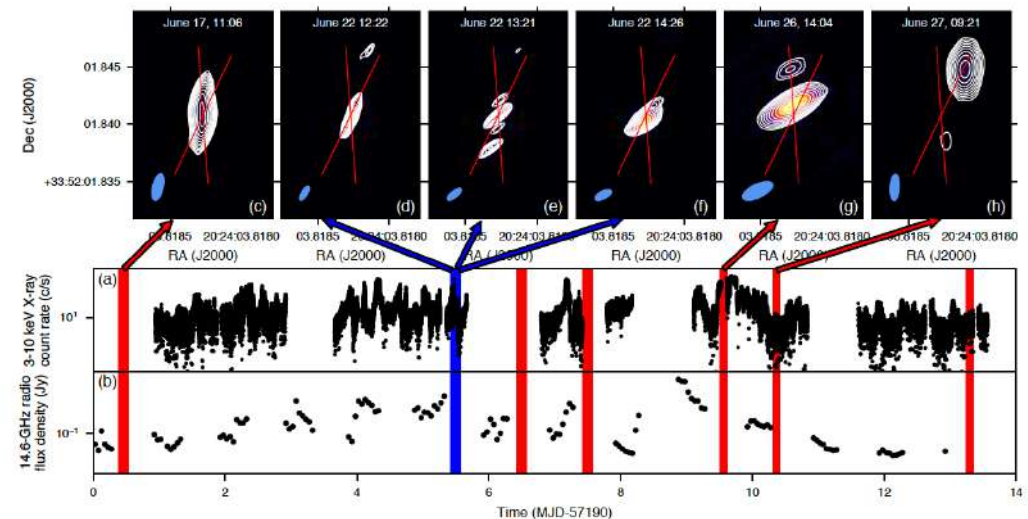
This timescale is consistent with the QPOs observed in some ULXs.

[2] Precession of jets in V404 Cygni:

The direction of jet is changing with time in V404 Cygni

(Miller-Jones et al. 2019).

Such behavior is consistent with our calculations.



# Summary

- Basic features of super-Eddington flows are revealed by Radiation-HD/Radiation-MHD simulations. The super-Eddington flows consist of the (1) radiation pressure-supported disk, (2) the radiatively-driven high-velocity outflows around the rotation axis, (3) and clumpy disk winds with wide opening angle.
- Our simulation results can explain the basic features of ULXs, Large X-ray luminosity, X-ray spectra, and so on.
- BZ effect and Lense-Thirring precession in the super-Eddington accretion has been confirmed by our GR-Radiation-MHD simulations.

For neutron stars; Takahashi & Ohsuga 2017, Takahashi, Ohsuga, et al. 2018, Inoue, Ohsuga et al. 2023