



Sphaleron in the Higgs Triplet Model

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in collaboration with Jiahang Hu and Shun Zhou

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D Theoretical Setup and Sphaleron Ansatz

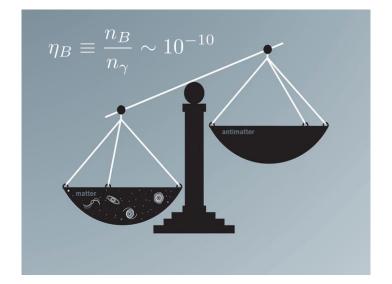
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Motivation

Two important **unsolved** problems in the Standard Model





origin of neutrino masses

matter-antimatter asymmetry

Can we accommodate **both** in a unified framework?

A successful attempt: **thermal leptogenesis**, SM + 3 singlet fermions

(naturally, new physics @ GUT scale) Fukugita, Yanagida, 1986

Motivation

Motivation

Higgs triplet model (HTM)

Alternatives: extend the SM by one complex triplet scalar

$$\mathcal{L} \supset -\frac{1}{2} \overline{\ell_{\mathrm{L}}} Y_{\Delta} \Delta \mathrm{i} \sigma^2 \ell_{\mathrm{L}}^c + \mathrm{h.c.} \quad \Rightarrow \quad M_{\nu} = Y_{\Delta} v_{\Delta}$$
(naturally, new physics @ TeV scale)

type-II seesaw mechanism

$\Delta \sim (1, 3, -1)$

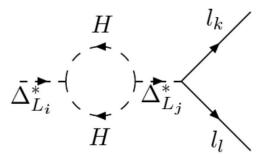
Konetschny et al., 1977 Magg *et al.*, 1980 Schechter et al., 1980 Cheng et al., 1980 Lazarides et al., 1981 Mohapatra et al., 1981

Can the HTM explain the baryon asymmetry of the Universe?

triplet leptogenesis

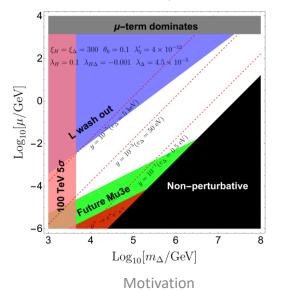
Affleck-Dine mechanism

Ma et al., 1998 Hambye et al., 2003



At least two triplet scalars are needed!

Barrie, Han, Murayama, 2021



electroweak baryogenesis

- The pattern of EW phase transition is modified by the triplet
- New sources of CP violation come from the leptonic sector

Can a successful EW baryogenesis be fulfilled in the HTM?

Prerequisite: working out the sphaleron configuration in the presence of the triplet scalar

Sphaleron in the Standard Model

 Vacuum structure of non-Abelian gauge theories is characterized by integer numbers (Chern-Simons numbers N_{CS})
 G. 't Hooft, 1976, Callan *et al.*, 1976

$$N_{
m CS} = - \, rac{g^2}{16 \pi^2} \! \int \! d^3x \, \, 2 \epsilon^{ijk} \, \, {
m Tr} igg[(\partial_i W_j) W_k - rac{2}{3} i g W_i W_j W_k igg]$$

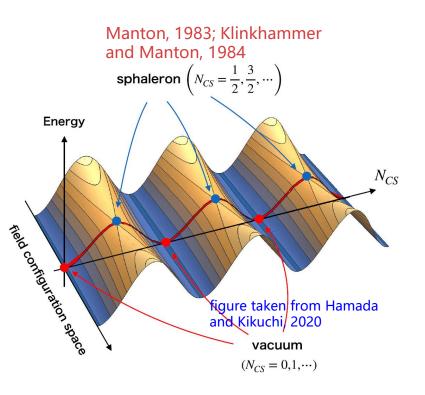
\Box Chiral anomaly: $\Delta B = \Delta L = N_f \times \Delta N_{\rm CS}$

D Instanton rate:
$$\exp(-16\pi^2/g^2) \sim 10^{-160}$$

- Sphaleron configuration: saddle-point solution of the energy functional
- □ Sphaleron rate (*B*-violating rate):

$$\begin{split} \Gamma_{\rm sph} &\sim \exp\left(-E_{\rm sph}/T\right) , \quad T < T_{\rm EW} \\ \Gamma_{\rm sph} &\sim \alpha_{\rm W}^5 T^4 , \qquad T > T_{\rm EW} \end{split}$$

 $\alpha_{\rm W} \equiv g^2/(4\pi)$ $E_{\rm sph} \sim 4\pi v/g \sim 5 \,\,{\rm TeV}$



Sphaleron in the Higgs Triplet Model

Assumptions:

Neglect the contributions from fermion fields

□ Neglect the effects from the finite Weinberg angle

Lagrangian in the HTM:

$$\phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}$$
$$\Delta = \begin{pmatrix} \Delta^- & -\sqrt{2}\Delta^0 \\ \sqrt{2}\Delta^{--} & -\Delta^- \end{pmatrix}$$

$$\mathcal{L} = -\frac{1}{2} \operatorname{Tr} \left(F_{\mu\nu} F^{\mu\nu} \right) + \left(D_{\mu} \phi \right)^{\dagger} \left(D^{\mu} \phi \right) + \frac{1}{2} \operatorname{Tr} \left[\left(D^{\mu} \Delta \right)^{\dagger} \left(D_{\mu} \Delta \right) \right] - V \left(\phi, \Delta \right)$$
$$\langle \phi \rangle = v_{\phi} / \sqrt{2} , \qquad \langle \Delta \rangle = v_{\Delta} , \qquad \sqrt{v_{\phi}^2 + 2v_{\Delta}^2} = v \approx 246 \text{ GeV}$$

Energy functional:

$$\begin{aligned} \mathcal{H}\left[W_{\mu},\phi,\Delta\right] &= \frac{1}{2}g^{ik}g^{jl}\mathrm{Tr}\left(F_{ij}F_{kl}\right) + g^{ij}\left(D_{i}\phi\right)^{\dagger}\left(D_{j}\phi\right) + \frac{1}{2}g^{ij}\left[\left(D_{i}\Delta\right)^{\dagger}\left(D_{j}\Delta\right)\right] + V\left(\phi,\Delta\right) \\ E\left[W_{\mu},\phi,\Delta\right] &= \int_{0}^{2\pi}\mathrm{d}\varphi\int_{0}^{\pi}\mathrm{d}\theta\sin\theta\int_{0}^{\infty}\mathrm{d}r\,r^{2}\,\mathcal{H}\left[W_{\mu},\phi,\Delta\right] \end{aligned}$$

Sphaleron in the Higgs Triplet Model $\mu \in [0, \pi]$

Non-contractible loops: $U(\mu, \theta, \varphi) = \begin{pmatrix} e^{i\mu} (\cos \mu - i \sin \mu \cos \theta) & e^{i\varphi} \sin \mu \sin \theta \\ -e^{-i\varphi} \sin \mu \sin \theta & e^{-i\mu} (\cos \mu + i \sin \mu \cos \theta) \end{pmatrix}$

$$U(\mu, \theta = 0, \varphi) = U(\mu = 0, \theta, \varphi) = U(\mu = \pi, \theta, \varphi) = \mathbf{1}$$

 $U(\mu, \theta, \varphi)$ defines a map: $S^3 \to S^3 \qquad \pi_3(S^3) = \mathbb{Z}$

Minmax procedure: $\mu = 0$ or $\mu = \pi$: vacuum configuration (minimal energy) $\mu = \pi/2$: sphaleron configuration (maximal energy)

Sphaleron ansatz:

$$\begin{split} W_{j}\left(\mu,r,\theta,\varphi\right) &= -\frac{1}{g}f(r)\underline{\partial_{j}U\left(\mu,\theta,\varphi\right)U^{-1}\left(\mu,\theta,\varphi\right)}, \quad j=\theta,\varphi \ , \\ \phi\left(\mu,r,\theta,\varphi\right) &= \frac{v_{\phi}}{\sqrt{2}}h(r)U\left(\mu,\theta,\varphi\right)\begin{pmatrix}0\\1\end{pmatrix}, \\ \text{fundamental rep.} \\ \Delta\left(\mu,r,\theta,\varphi\right) &= v_{\Delta}h_{\Delta}(r)U\left(\mu,\theta,\varphi\right)\begin{pmatrix}0-1\\0&0\end{pmatrix}U^{-1}\left(\mu,\theta,\varphi\right) \ , \end{split}$$

$$f(0) = h(0) = h_{\Delta}(0) = 0$$

$$f(\infty) = h(\infty) = h_{\Delta}(\infty) = 1$$

adjoint rep.

Sphaleron in the Higgs Triplet Model

Kinematic terms:

$$\frac{1}{2}g^{jk}g^{jl}\operatorname{Tr}\left(F_{ij}F_{kl}\right) = \frac{4}{g^2r^4}\sin^2\mu\left[2f^2\left(1-f\right)^2\sin^2\mu+r^2f'^2\right] \qquad \text{spherically symmetric} \\ g^{ij}\left(D_i\phi\right)^{\dagger}\left(D_j\phi\right) = \frac{v_{\phi}^2}{2r^2}\left[2\left(1-f\right)^2h^2\sin^2\mu+r^2h'^2\right] \qquad \text{spherically symmetric} \\ \frac{1}{2}g^{ij}\left[\left(D_i\Delta\right)^{\dagger}\left(D_j\Delta\right)\right] = \frac{v_{\Delta}^2}{2r^2}\left[\left(5-\cos 2\theta\right)\left(1-f\right)^2h_{\Delta}^2\sin^2\mu+r^2h_{\Delta}'^2\right] \qquad \begin{array}{c} \text{non-spherically} \\ \text{symmetric} \end{array}$$

Scalar potential:

$$V(\phi, \Delta) = \lambda \left(\phi^{\dagger}\phi\right)^{2} - \kappa^{2}\phi^{\dagger}\phi + \frac{1}{2}M_{\Delta}^{2}\operatorname{Tr}\left(\Delta^{\dagger}\Delta\right) - \left(\lambda_{\Delta}M_{\Delta}\phi^{\mathrm{T}}\epsilon\Delta\phi + \mathrm{h.c.}\right) + \frac{\lambda_{1}}{4}\left[\operatorname{Tr}\left(\Delta^{\dagger}\Delta\right)\right]^{2} + \frac{\lambda_{2}}{4}\operatorname{Tr}\left[\left(\Delta^{\dagger}\Delta\right)^{2}\right] + \lambda_{3}\left(\phi^{\dagger}\phi\right)\operatorname{Tr}\left(\Delta^{\dagger}\Delta\right) + \lambda_{4}\phi^{\dagger}\Delta\Delta^{\dagger}\phi$$

8 real parameters

but not all of them are relevant to the sphaleron configuration

VEVs determined from:

$$\frac{\partial}{\partial v_{\phi}} V\left(v_{\phi}, v_{\Delta}\right) = \left(-\kappa^{2} + \lambda v_{\phi}^{2} - 2\lambda_{\Delta}M_{\Delta}v_{\Delta} + \lambda_{3}v_{\Delta}^{2}\right)v_{\phi} = 0$$
$$\frac{\partial}{\partial v_{\Delta}} V\left(v_{\phi}, v_{\Delta}\right) = -\lambda_{\Delta}M_{\Delta}v_{\phi}^{2} + M_{\Delta}^{2}v_{\Delta} + (\lambda_{1} + \lambda_{2})v_{\Delta}^{3} + \lambda_{3}v_{\phi}^{2}v_{\Delta} = 0$$

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Theoretical Setup and Sphaleron Ansatz

Sphaleron Equations of Motion

$$\begin{aligned} \text{Sphaleron energy:} \quad & \xi \equiv gvr \approx 8.1 \times \left(\frac{r}{10^{-15} \text{ cm}}\right) \qquad \beta \equiv v^2/v_{\phi}^2 = 1 + 2\varrho_3 \\ E_{\text{sph}} &= \frac{4\pi v}{g} \int_0^\infty d\xi \left\{ 4f'^2 + \frac{8}{\xi^2} f^2 (1-f)^2 + \frac{1}{\beta} (1-f)^2 h^2 + \frac{1}{2\beta} \xi^2 h'^2 \\ & \quad + \frac{\xi^2}{4\beta^2} \left[(\varrho_1 - \varrho_2) (1-h^2)^2 + \varrho_2 (h^2 - h_{\Delta})^2 \right] + \frac{\varrho_3}{6\beta} \left[3\xi^2 h_{\Delta}'^2 + 16h_{\Delta}^2 (1-f)^2 \right] \\ & \quad + \frac{\xi^2}{4\beta^2} \left[2 (\varrho_4 - \varrho_1 + \varrho_2) (1-h^2) - (2\varrho_3 \varrho_5 - \varrho_2) (1-h_{\Delta}^2) \right] \\ & \quad + \frac{\xi^2}{2\beta^2} \left(\sqrt{2\varrho_2 \varrho_3 \varrho_5} - \varrho_2 \right) (1-h^2 h_{\Delta}) + \frac{\xi^2}{2\beta^2} \left(\varrho_1 - \varrho_4 - \sqrt{2\varrho_2 \varrho_3 \varrho_5} \right) (1-h^2 h_{\Delta}^2) \\ & \quad - \frac{\xi^2}{4\beta^2} \left(\varrho_1 - \varrho_4 - \varrho_3 \varrho_5 - \sqrt{\varrho_2 \varrho_3 \varrho_5/2} \right) (1-h_{\Delta}^4) \right\} . \end{aligned}$$

$$\begin{array}{rcl} \varrho_1 &\equiv& \lambda/g^2 \\ \varrho_2 &\equiv& 2\lambda_{\Delta}^2/g^2 \\ \varrho_3 &\equiv& v_{\Delta}^2/v_{\phi}^2 \\ \varrho_4 &\equiv& \kappa^2/\left(g^2 v_{\phi}^2\right) \\ \varrho_5 &\equiv& M_{\Delta}^2/\left(g^2 v_{\phi}^2\right) \end{array}$$

sphaleron configuration in the HTM is determined by 5 independent parameters

Equations of motion:

$$\begin{aligned} \xi^{2}f'' &= 2f(1-f)(1-2f) - \frac{\xi^{2}}{4\beta}(1-f)h^{2} - \frac{2\varrho_{3}}{3\beta}\xi^{2}(1-f)h_{\Delta}^{2} \\ (\xi^{2}h')' &= 2(1-f)^{2}h - \frac{\xi^{2}}{\beta}\left[(\varrho_{1}-\varrho_{2})h(1-h^{2}) - \varrho_{2}h(h^{2}-h_{\Delta}) \\ &+ (\varrho_{4}-\varrho_{1}+\varrho_{2})h + \left(\sqrt{2\varrho_{2}\varrho_{3}\varrho_{5}} - \varrho_{2}\right)hh_{\Delta} + \left(\varrho_{1}-\varrho_{4}-\sqrt{2\varrho_{2}\varrho_{3}\varrho_{5}}\right)hh_{\Delta}^{2}\right] \\ \varrho_{3}\left(\xi^{2}h'_{\Delta}\right)' &= \frac{16}{3}\varrho_{3}(1-f)^{2}h_{\Delta} - \frac{\varrho_{2}\xi^{2}}{2\beta}(h^{2}-h_{\Delta}) + \frac{\xi^{2}}{2\beta}\left[(2\varrho_{3}\varrho_{5}-\varrho_{2})h_{\Delta} - \left(\sqrt{2\varrho_{2}\varrho_{3}\varrho_{5}} - \varrho_{2}\right)h^{2} \\ &- 2\left(\varrho_{1}-\varrho_{4}-\sqrt{2\varrho_{2}\varrho_{3}\varrho_{5}}\right)h^{2}h_{\Delta} + 2\left(\varrho_{1}-\varrho_{4}-\varrho_{3}\varrho_{5}-\sqrt{\varrho_{2}\varrho_{3}\varrho_{5}/2}\right)h_{\Delta}^{3}\right] \end{aligned}$$
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Constraints on Parameters

EW precision measurements on triplet VEV:

$$\varrho_3 = v_{\Delta}^2 / v_{\phi}^2 \lesssim 10^{-3} \qquad \qquad \varrho_4 \approx \varrho_1 - \frac{1}{2} \varrho_3 \varrho_5 \left(2 + \sqrt{\frac{2\varrho_2}{\varrho_3 \varrho_5}} \right)$$

$$\begin{array}{rcl} \varrho_1 &\equiv& \lambda/g^2\\ \varrho_2 &\equiv& 2\lambda_{\Delta}^2/g^2\\ \varrho_3 &\equiv& v_{\Delta}^2/v_{\phi}^2\\ \varrho_4 &\equiv& \kappa^2/\left(g^2v_{\phi}^2\right)\\ \varrho_5 &\equiv& M_{\Delta}^2/\left(g^2v_{\phi}^2\right)\\ \end{array}$$
$$\frac{\lambda_3}{g^2} \approx \sqrt{\frac{\varrho_2\varrho_5}{2\varrho_3}} - \varrho_5 \end{array}$$

$$\begin{split} 0 &< \varrho_1 \leqslant \frac{4\pi}{g^2} \ , \quad -\sqrt{\frac{4\pi}{g^2}} \varrho_1 < \sqrt{\frac{\varrho_2 \varrho_5}{2\varrho_3}} - \varrho_5 \leqslant \frac{4\pi}{g^2} \ , \quad \varrho_1 - \varrho_3 \varrho_5 - \sqrt{\varrho_2 \varrho_3 \varrho_5/2} > 0 \\ &-\sqrt{\frac{4\pi}{g^2}} \varrho_1 < \frac{\lambda_3 + \lambda_4}{g^2} \leqslant \frac{4\pi}{g^2} \ , \quad |2\lambda_3 + 3\lambda_4| \leqslant 8\pi \ , \quad |2\lambda_3 - \lambda_4| \leqslant 8\pi \end{split}$$

D Higgs mass $m_h \approx 125 \text{ GeV}$

$$\frac{\lambda_4}{g^2} \approx \varrho_5 \pm \frac{1}{\left(2\varrho_3\right)^{3/4}} \sqrt{\left(\varrho_1 - \frac{m_h^2}{2g^2 v_\phi^2}\right) \left(\sqrt{\varrho_2 \varrho_5} - \frac{\sqrt{2\varrho_3}m_h^2}{g^2 v_\phi^2}\right)}$$

Collider constraints

$$m_{H^{\pm\pm}} \gtrsim 350 \text{ GeV}$$

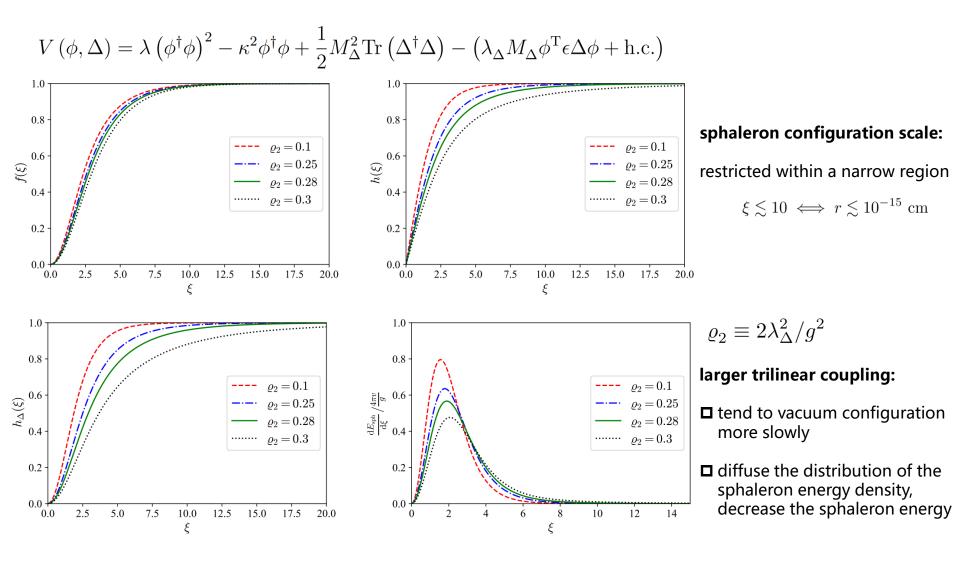
□ Charged lepton flavor violation

$$m_{H^{\pm\pm}}^2 \approx g^2 v_{\phi}^2 \left(\sqrt{\frac{\varrho_2 \varrho_5}{2\varrho_3}} - \frac{\lambda_4}{g^2} \right) \quad \Rightarrow \quad \sqrt{\frac{\varrho_2 \varrho_5}{2\varrho_3}} - \frac{\lambda_4}{g^2} \gtrsim 4.8$$

 $M_{\Delta} v_{\Delta} \gtrsim 10^2 \text{ GeV} \cdot \text{eV} \quad \Rightarrow \quad \varrho_3 \varrho_5 \gtrsim 10^{-24}$

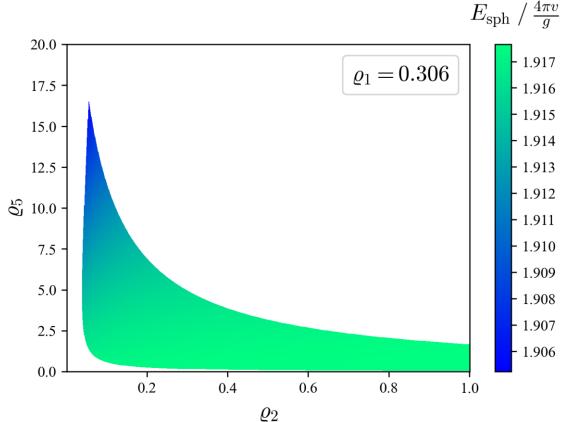
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Sphaleron Field Configuration



Sphaleron Energy

allowed parameter space in the HTM:



$$\rho_3 = 10^{-3}$$
$$\varrho_4 \approx \varrho_1 - \frac{1}{2} \varrho_3 \varrho_5 \left(2 + \sqrt{\frac{2\varrho_2}{\varrho_3 \varrho_5}} \right)$$

In the SM:

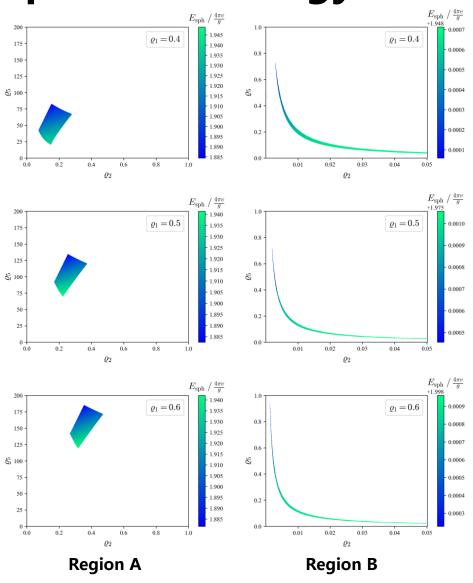
 $\varrho_1^{\rm SM} = m_h^2 / \left(2g^2 v_\phi^2\right) \approx 0.306$

$$E_{\rm sph}^{\rm SM} \approx 1.92 \times \frac{4\pi v}{g} \approx 9 {
m TeV}$$

In the HTM:

larger trilinear coupling or heavier triplet scalar ↔ smaller sphaleron energy

Sphaleron Energy



$$\begin{split} \varrho_1^{\rm SM} &= m_h^2 / \left(2g^2 v_\phi^2 \right) \approx 0.306 \\ E_{\rm sph}^{\rm SM} &\approx 1.92 \times 4\pi v / g \end{split}$$

parameter space starts to split into two parts as $\rho_1 > 0.306$

Region A: large trilinear coupling and heavy triplet scalar $M_{\Delta} \gtrsim 1 \text{ TeV}$

Region B: small trilinear coupling and light triplet scalar $M_{\Delta} \lesssim 1 \text{ TeV}$ (narrow band, testable by future collider searches)

sphaleron energy in **Region A**:

 $1.88\times 4\pi v/g \lesssim E_{\rm sph} \lesssim 1.97\times 4\pi v/g$

sphaleron energy in Region B:

 $1.92 \times 4\pi v/g \lesssim E_{\rm sph} \lesssim 2.48 \times 4\pi v/g$

enhanced up to 30% compared with the SM case!

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Sphaleron Solution

Summary

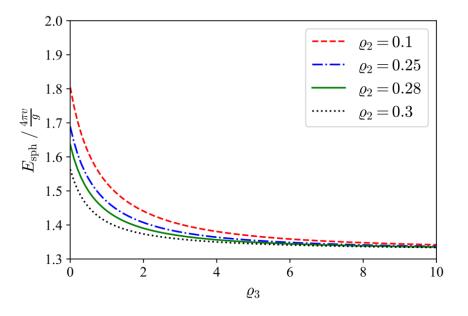
- To carry out a consistent EW baryogenesis in the HTM, we calculate the sphaleron configuration in the presence of a triplet.
- The sphaleron configuration in the HTM depends on 5 parameters: the quartic coupling parameter ϱ_1 would increase the sphaleron energy; the trilinear coupling parameter ϱ_2 , the VEV ratio parameter ϱ_3 , and the triplet mass parameter ρ_5 would decrease the sphaleron energy.
- Basically, the difference of the sphaleron energy between the SM and the HTM is suppressed by the small triplet VEV. Interestingly, there still exists some narrow parameter space where the sphaleron energy could be enhanced by 30% compared with the SM case. Such narrow space can be tested by future collider searches.

Future Extensions

Ginite temperature effects:

$$\begin{split} E_{\rm sph}(T) &= E_{\rm sph} \frac{v(T)}{v} \\ v(T) &= \sqrt{v_{\phi}^2(T) + 2v_{\Delta}^2(T)} \end{split}$$

 $\varrho_3 = v_{\Delta}^2 / v_{\phi}^2$ is not suppressed by experiments at finite temperature



□ Effects of leptonic CP violation on EW baryogenesis:

The triplet coupling violates the lepton number and brings about new sources of CP violation.

This may have sizable effects on EW baryogenesis (collision between the triplet and the bubble wall will propagate new CP asymmetries).

Thank you!

Q&A