

Higgs boson pair production and decay to $b\bar{b}\gamma\gamma$ at NLO in QCD

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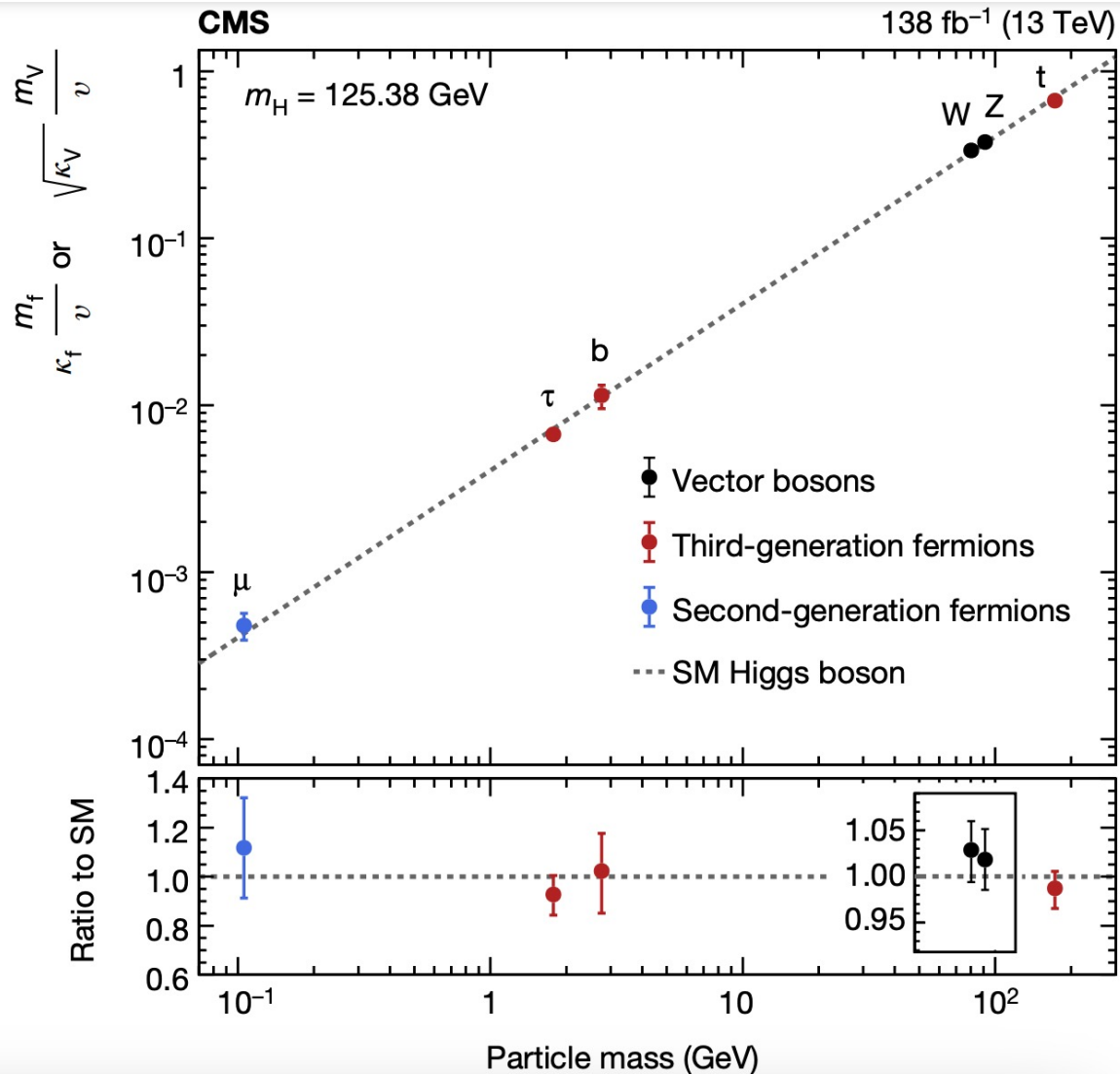
in collaboration with Haitao Li, Zongguo Si, Jian Wang, Xiao Zhang

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OUTLINE

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2. Framework
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Introduction

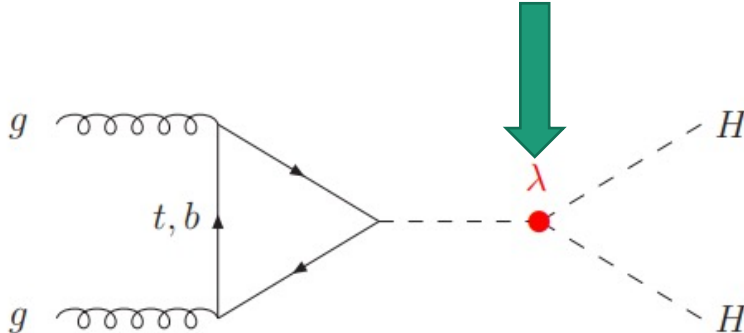


[CMS, Nature 607, 60–68 (2022)]

Why we study Higgs boson pair production?

Higgs pair production is a key to probing Higgs self-coupling!

$$V(h) = \frac{1}{2} m_h^2 h^2 + \boxed{\lambda v h^3} + \frac{1}{4} \lambda h^4$$

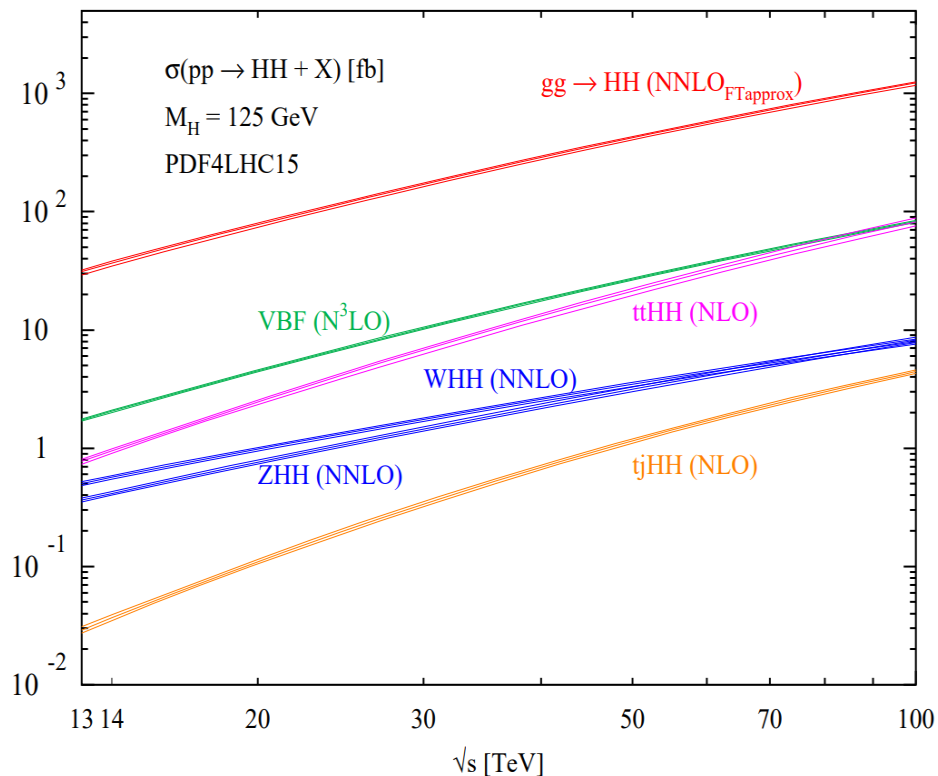


The trilinear Higgs self-coupling λ is important for examining the shape of Higgs potential and understanding EWSB.

Introduction

Higgs pair production by ggF: Theoretical (partial) Status

Main production channel



[Biagio Di Micco *et al*, arXiv:1910.00012]

Heavy top limit:

N3LO+N3LL: [Ajjath, Shao, arXiv:2209.03914]

N3LO: [Chen, Li, Shao, Wang, arXiv:1909.06808]

NNLO: [de Florian, Grazzini, Hanga, Kallweit, Lindert, Maierhöfer, Mazzitelli, Rathlev, arXiv:1606.09519]

NLO+NNLL: [Shao, Li, Li, Wang, arXiv:1301.1245]

Finite m_t dependence:

NNLO_{FTapprox} +NNLL: [De Florian, Mazzitelli, arXiv:1807.03704]

NNLO_{FTapprox}: [Grazzini, Heinrich, Jones, Kallweit, Kerner, arXiv:1803.02463] (applied by CMS and ATLAS)

NLO+NLL: [Feerera, Pires, arXiv:1609.01691]

NLO: [Borowka, *et al*, arXiv:1608.14798, arXiv:1604.06477]

NLO EW: [Davies, Schonwald, Steinhauser, Zhang, arXiv:2308.01355]

NNLO: [Mazzitelli, arXiv:2206.14667]

[Czakon, Harlander, Klappert, Niggetiedt, arXiv:2105.04436]

NLO: [Baglio, Campanario, Glaus, Muhlleitner, Spira, Streicher, arXiv:1811.05692]

Introduction

$$\kappa_\lambda = \frac{\lambda_{HHH}^{EXP}}{\lambda_{HHH}^{SM}} \quad \boxed{\kappa - \text{framework}} \quad [\text{CERN-2013-004 (CERN, 2013)}]$$

$$\sigma(gg \rightarrow HH) \sim \frac{\sigma(gg \rightarrow H)}{1000}$$

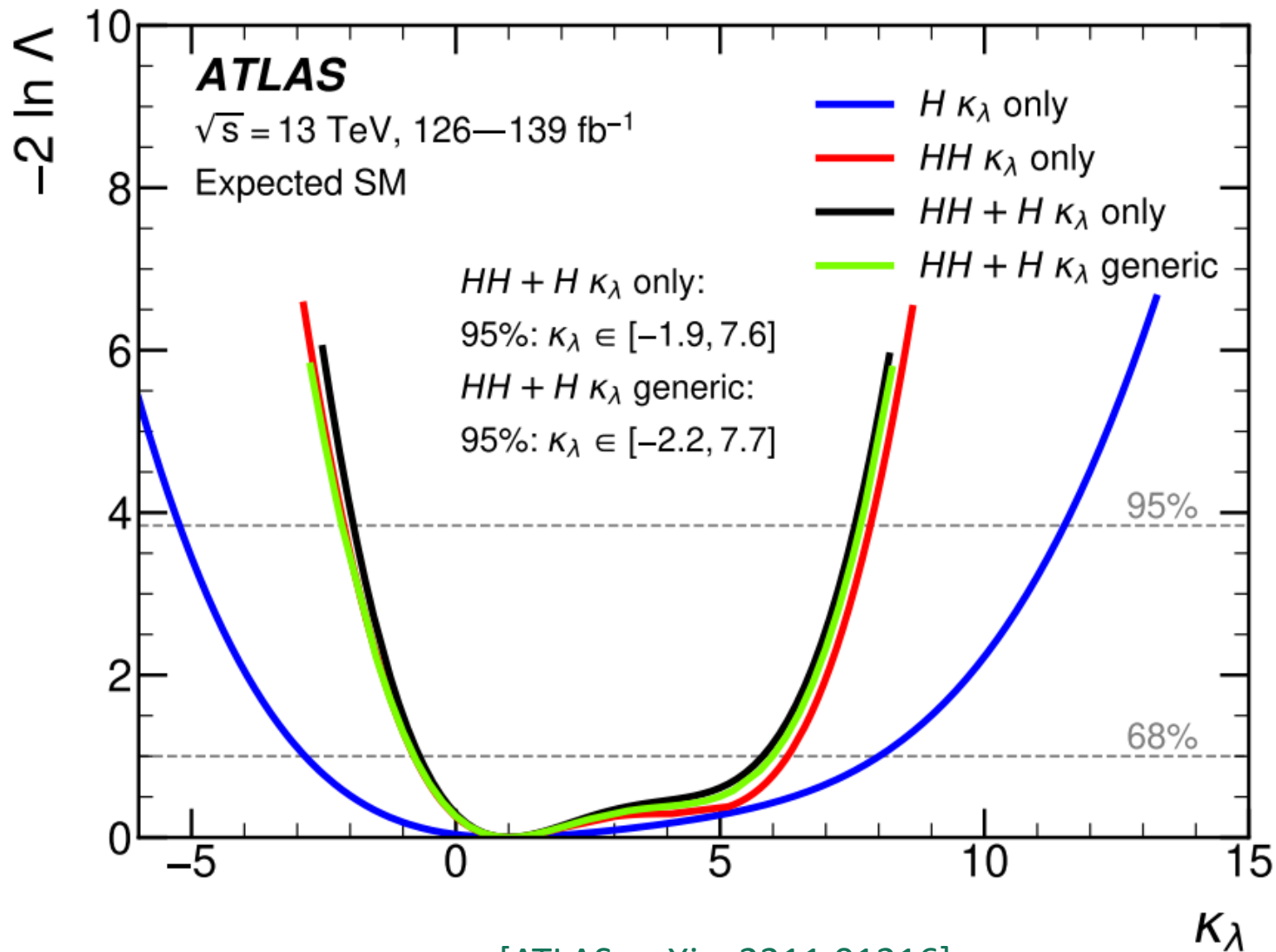
$HH \rightarrow b\bar{b}\gamma\gamma$ has the highest S/B.

$HH \rightarrow b\bar{b}b\bar{b}$ has the highest branching ratio.

$HH \rightarrow b\bar{b}\tau\tau$ has middle S/B and branching ratio.

ATLAS:	Final states	Obs.	Exp.	
	$b\bar{b}b\bar{b}$	[-3.5,11.3]	[-5.4,11.4]	[ATLAS,arXiv:2301.03212]
	$b\bar{b}\gamma\gamma$	[-1.5,6.7]	[-2.4,7.7]	[ATLAS,arXiv:2112.11876]
	$b\bar{b}\tau\tau$	[-2.4,9.2]	[-2.0,9.0]	[ATLAS-CONF-2021-052]
	Combine	[-0.6,6.6]	[-2.1,7.8]	[ATLAS,PLB 843(2023)137745]
CMS:	Final states	Obs.	Exp.	
	$b\bar{b}b\bar{b}$	[-2.3,9.4]	[-5.0,12.0]	[CMS,arXiv:2202.09617]
	$b\bar{b}\gamma\gamma$	[-3.3,8.5]	[-2.5,8.2]	[CMS,arXiv:2011.12373]
	$b\bar{b}\tau\tau$	[-1.7,8.7]	[-2.9,9.8]	[CMS,arXiv:2206.09401]
	Combine	[-1.24,6.49]	[-2.28,7.94]	[Higgs 2022, N. De Filippis's talk]

Introduction



[ATLAS, arXiv: 2211.01216]

$$\kappa_\lambda = \frac{\lambda_{HHH}}{\lambda_{HHH}^{SM}}$$

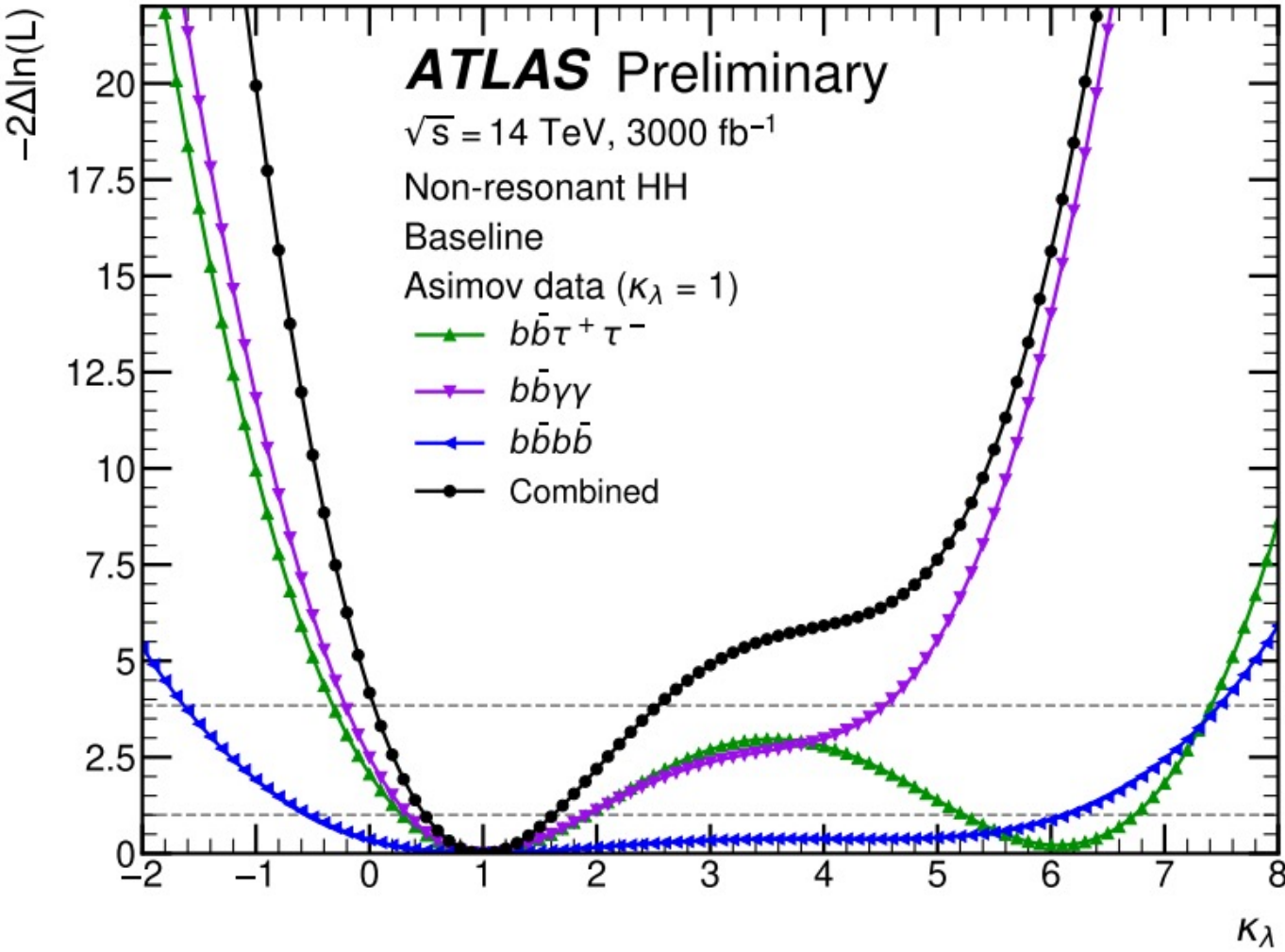
Expected self-coupling modifier at 95% C.L.

Single- and di-Higgs combination:

Run 2: $\kappa_\lambda \in [-1.9, 7.6]$

Constraint is still loose.

Introduction



$$\kappa_\lambda = \frac{\lambda_{HHH}}{\lambda_{HHH}^{SM}}$$

Expected self-coupling modifier at 95% C.L.

HL-LHC projection : $\kappa_\lambda \in [0.0, 2.5]$

Need more precise theoretical prediction.

Framework

- For the specific Higgs pair decay model $H_1 \rightarrow X_1, H_2 \rightarrow X_2$, in the frame of narrow width approximation, we have,

$$\sigma_{pro+dec}(X_1, X_2) = \sigma_{pro} \frac{\int d\Gamma_{H_1 \rightarrow X_1}}{\Gamma_{H_1 \rightarrow X_1}} \frac{\int d\Gamma_{H_2 \rightarrow X_2}}{\Gamma_{H_2 \rightarrow X_2}} \times R(H_1 \rightarrow X_1) R(H_2 \rightarrow X_2).$$

The cross-section can be expanded in a series of strong coupling,

$$\sigma_{pro+dec(\bar{b}b, \gamma\gamma)}^{(0)} = \sigma_{pro}^{(0)} \frac{\int d\Gamma_{H_1 \rightarrow b\bar{b}}^{(0)}}{\Gamma_{H_1 \rightarrow b\bar{b}}^{(0)}} \frac{\int d\Gamma_{H_2 \rightarrow \gamma\gamma}^{(0)}}{\Gamma_{H_2 \rightarrow \gamma\gamma}^{(0)}} \times R(H_1 \rightarrow b\bar{b}) R(H_2 \rightarrow \gamma\gamma). \quad LO$$

$$\sigma_{pro+dec(\bar{b}b, \gamma\gamma)}^{pro(1)} = \sigma_{pro}^{(1)} \frac{\int d\Gamma_{H_1 \rightarrow b\bar{b}}^{(0)}}{\Gamma_{H_1 \rightarrow b\bar{b}}^{(0)}} \frac{\int d\Gamma_{H_2 \rightarrow \gamma\gamma}^{(0)}}{\Gamma_{H_2 \rightarrow \gamma\gamma}^{(0)}} \times R(H_1 \rightarrow b\bar{b}) R(H_2 \rightarrow \gamma\gamma). \quad \delta NLO^{pro} \times LO^{dec}$$

Framework

$$\begin{aligned} \sigma_{pro+dec(b\bar{b},\gamma\gamma)}^{dec(1)} &= \sigma_{pro}^{(0)} \frac{1}{\Gamma_{H_1 \rightarrow b\bar{b}}^{(0)}} \frac{1}{\Gamma_{H_2 \rightarrow \gamma\gamma}^{(0)}} \int d\Gamma_{H_1 \rightarrow b\bar{b}}^{(1)} \int d\Gamma_{H_2 \rightarrow \gamma\gamma}^{(0)} \times R(H_1 \rightarrow b\bar{b}) R(H_2 \rightarrow \gamma\gamma) \\ &\quad - \sigma_{pro}^{(0)} \frac{\Gamma_{H_1 \rightarrow b\bar{b}}^{(1)}}{(\Gamma_{H_1 \rightarrow b\bar{b}}^{(0)})^2} \frac{1}{\Gamma_{H_2 \rightarrow \gamma\gamma}^{(0)}} \int d\Gamma_{H_1 \rightarrow b\bar{b}}^{(0)} \int d\Gamma_{H_2 \rightarrow \gamma\gamma}^{(0)} \times R(H_1 \rightarrow b\bar{b}) R(H_2 \rightarrow \gamma\gamma) \end{aligned}$$

Numerator expanding

Denominator expanding

$$\sigma_{pro+dec(b\bar{b},\gamma\gamma)}^{dec(1)} = \sigma_{pro}^{(0)} \frac{\int d\Gamma_{H_1 \rightarrow b\bar{b}}^{(0)} \int d\Gamma_{H_2 \rightarrow \gamma\gamma}^{(0)}}{\Gamma_{H_1 \rightarrow b\bar{b}}^{(0)} \Gamma_{H_2 \rightarrow \gamma\gamma}^{(0)}} \left(\frac{\int d\Gamma_{H_1 \rightarrow b\bar{b}}^{(1)}}{\int d\Gamma_{H_1 \rightarrow b\bar{b}}^{(0)}} - \frac{\Gamma_{H_1 \rightarrow b\bar{b}}^{(1)}}{\Gamma_{H_1 \rightarrow b\bar{b}}^{(0)}} \right) \times R(H_1 \rightarrow b\bar{b}) R(H_2 \rightarrow \gamma\gamma)$$

$LO^{pro} \times \delta NLO^{dec}$

This contribution is **0** in the total cross-section.

Why we consider its correction effects ?

Why we **ignore** the $H_2 \rightarrow \gamma\gamma$ NLO QCD correction ?

$$\dots \left(\frac{\int d\Gamma_{H_2 \rightarrow \gamma\gamma}^{(1)}}{\int d\Gamma_{H_2 \rightarrow \gamma\gamma}^{(0)}} - \frac{\Gamma_{H_2 \rightarrow \gamma\gamma}^{(1)}}{\Gamma_{H_2 \rightarrow \gamma\gamma}^{(0)}} \right) \dots \stackrel{\text{Cut}}{=} 0$$

Numerical result

	without decays	with decays but no cuts		with decays and cuts	
		LO ^{dec}	δ NLO ^{dec}	LO ^{dec}	δ NLO ^{dec}
LO _{∞} ^{pro}	17.07 ^{+31%} _{-22%}	0.02257 ^{+31%} _{-22%}	0	0.01257 ^{+30%} _{-22%}	-0.001750 ^{+42%} _{-28%}
LO _{m_t} ^{pro}	19.85 ^{+28%} _{-21%}	0.02624 ^{+28%} _{-20%}	0	0.01395 ^{+27%} _{-20%}	-0.002613 ^{+39%} _{-27%}
δ NLO _{∞} ^{pro}	14.86 ^{+6%} _{-7%}	0.01964 ^{+6%} _{-7%}	—	0.01064 ^{+6%} _{-7%}	—
δ NLO _{m_t} ^{pro}	13.08 ^{+4%} _{-8%}	0.01729 ^{+4%} _{-8%}	—	0.009142 ^{+4%} _{-8%}	—
Full NLO result					
NLO _{∞}	31.93 ^{+18%} _{-15%}	0.04221 ^{+18%} _{-15%}		0.02146 ^{+15%} _{-14%}	
NLO _{m_t}	32.93 ^{+14%} _{-13%}	0.04354 ^{+14%} _{-13%}		0.02047 ^{+10%} _{-11%}	

$\sqrt{s} = 14$ TeV
 $m_t = 173$ GeV
 $m_H = 125$ GeV
 PDF4LHC15_nlo_100_pdfas

$\mu_R, \mu_F \in \frac{m_{HH}}{2} \left(\frac{1}{2}, 1, 2 \right)$ and $\frac{\mu_R}{\mu_F} = \frac{1}{4}$ or 4 are excluded (7-point).

$$LO_{m_t} / LO_{\infty} = 16\%$$



$$NLO_{m_t} / NLO_{\infty} = 3\%$$

$$\delta NLO_{\infty} / LO_{\infty} = 87\%$$



$$\delta NLO_{m_t} / LO_{m_t} = 66\%$$

Agree with [arXiv:1909.06808] (HTL) and [arXiv:1608.14798] (SM)

Numerical result

	without decays	with decays but no cuts		with decays and cuts	
		LO ^{dec}	δNLO ^{dec}	LO ^{dec}	δNLO ^{dec}
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In numerical data:

$$R(H \rightarrow b\bar{b}) \times R(H \rightarrow \gamma\gamma) = 58.24\% \times 0.227\% = 0.13\%$$

$$LO^{pro} LO^{dec} = LO^{pro} \times R(H \rightarrow b\bar{b}) \times R(H \rightarrow \gamma\gamma)$$

$$\delta NLO^{pro} LO^{dec} = \delta NLO^{pro} \times R(H \rightarrow b\bar{b}) \times R(H \rightarrow \gamma\gamma)$$

$$LO^{pro} \delta NLO^{dec} = \dots \left(\frac{\int d\Gamma_{H_1 \rightarrow b\bar{b}}^{(1)}}{\int d\Gamma_{H_1 \rightarrow b\bar{b}}^{(0)}} - \frac{\Gamma_{H_1 \rightarrow b\bar{b}}^{(1)}}{\Gamma_{H_1 \rightarrow b\bar{b}}^{(0)}} \right) \dots = 0$$

Confirm the validity of our numerical program.

Numerical result

	without decays	with decays but no cuts		with decays and cuts	
		LO ^{dec}	δNLO ^{dec}	LO ^{dec}	δNLO ^{dec}
LO _∞ ^{pro}	17.07 ^{+31%} _{-22%}	0.02257 ^{+31%} _{-22%}	0	0.01257 ^{+30%} _{-22%}	-0.001750 ^{+42%} _{-28%}
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$p_T^j, p_T^\gamma > 25 \text{ GeV}$
 $|\eta^j|, |\eta^\gamma| < 2.5$
 $\Delta R_{jj, j\gamma, \gamma\gamma} > 0.4$
 $90 < m_{jj} < 190 \text{ GeV}$

$$\dots \left(\frac{\int d\Gamma_{H_1 \rightarrow b\bar{b}}^{(1)}}{\int d\Gamma_{H_1 \rightarrow b\bar{b}}^{(0)}} - \frac{\Gamma_{H_1 \rightarrow b\bar{b}}^{(1)}}{\Gamma_{H_1 \rightarrow b\bar{b}}^{(0)}} \right) \dots = 0 \quad \xrightarrow{\text{CUT}} \quad \dots \left(\frac{\int d\Gamma_{H_1 \rightarrow b\bar{b}}^{(1)}}{\int d\Gamma_{H_1 \rightarrow b\bar{b}}^{(0)}} - \frac{\Gamma_{H_1 \rightarrow b\bar{b}}^{(1)}}{\Gamma_{H_1 \rightarrow b\bar{b}}^{(0)}} \right) \dots \neq 0$$

$$\left(LO_{m_t}^{pro} \delta NLO^{dec} / LO_{m_t}^{pro} LO^{dec} \right) \xrightarrow{\text{Cut}} -19\%$$

$$\left(LO_{\infty}^{pro} \delta NLO^{dec} / LO_{\infty}^{pro} LO^{dec} \right) \xrightarrow{\text{Cut}} -14\%$$

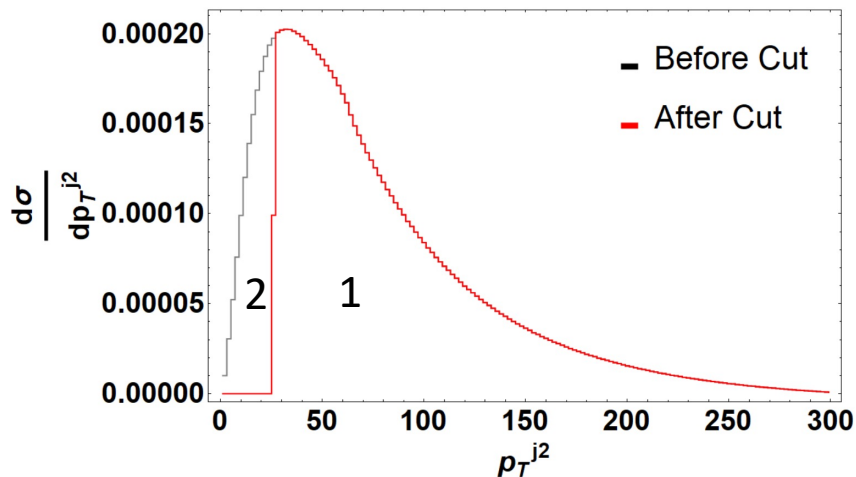
Larger than N³LO QCD correction (~+6%)

[Chen, Li, Shao, Wang, arXiv:1910.00012]

Numerical result

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δ NLO _{m_t} ^{pro}	13.08 ^{+4%} _{-8%}	0.01729 ^{+4%} _{-8%}	—	0.009142 ^{+4%} _{-8%}	—
Full NLO result					
NLO _{∞}	31.93 ^{+18%} _{-15%}	0.04221 ^{+18%} _{-15%}	-49%	0.02146 ^{+15%} _{-14%}	
NLO _{m_t}	32.93 ^{+14%} _{-13%}	0.04354 ^{+14%} _{-13%}	-53%	0.02047 ^{+10%} _{-11%}	

$p_T^j, p_T^\gamma > 25 \text{ GeV}$
 $|\eta^j|, |\eta^\gamma| < 2.5$
 $\Delta R_{jj, j\gamma, \gamma\gamma} > 0.4$
 $90 < m_{jj} < 190 \text{ GeV}$

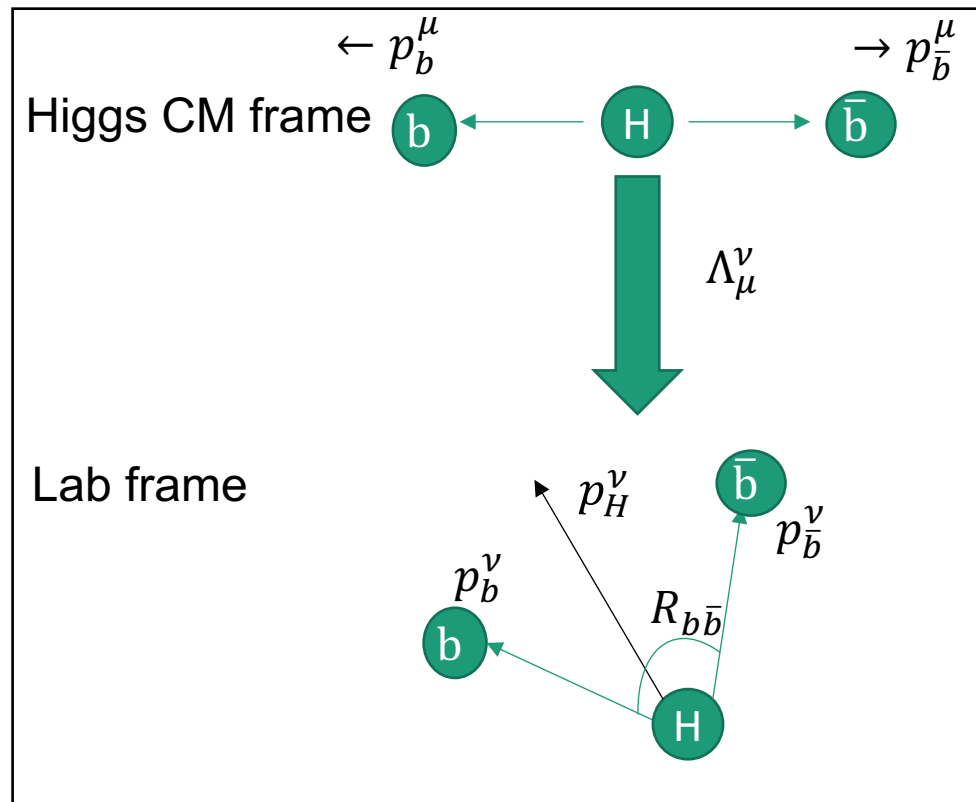
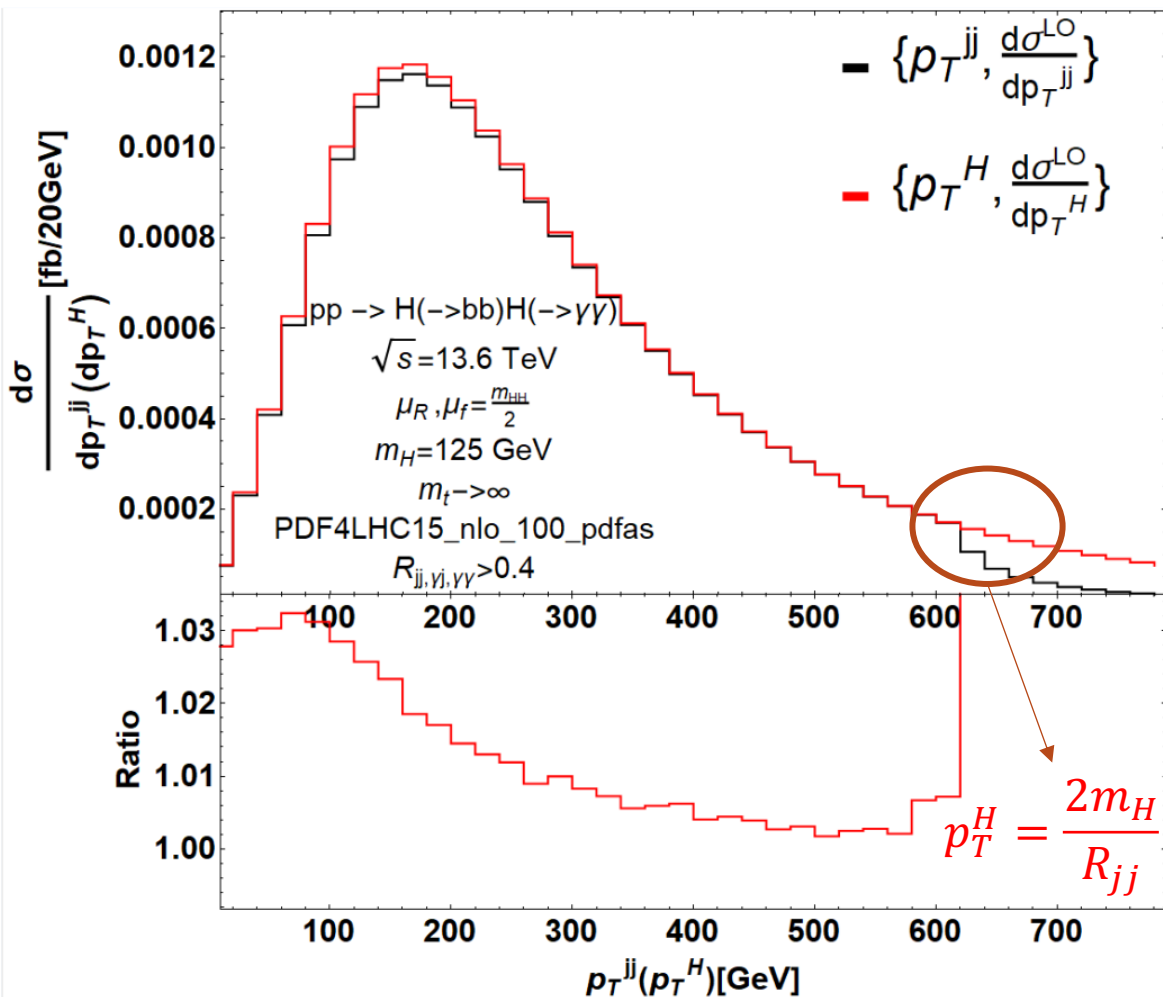


$$\frac{\text{Region1}}{\text{Region2}} \sim \frac{4}{1}$$

The cross-section with low p_T is cut.

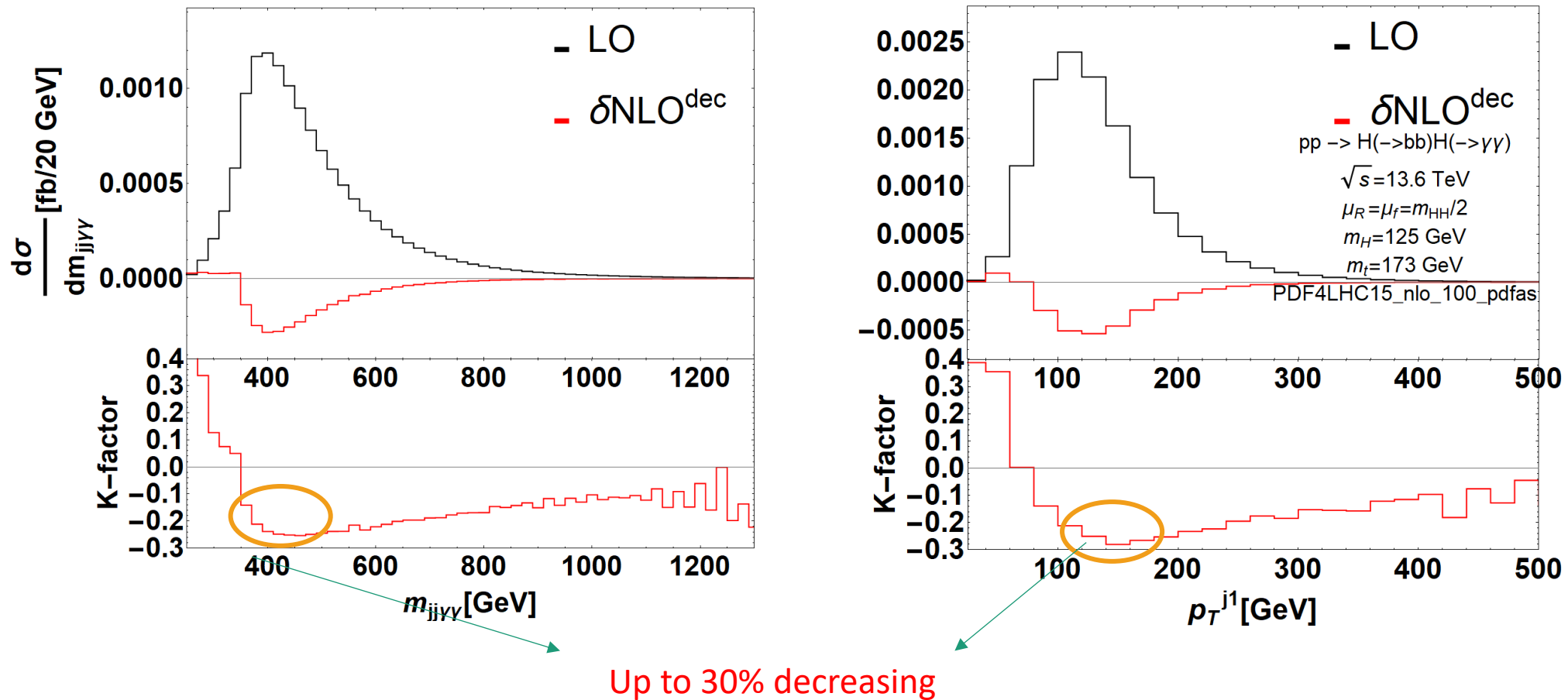
Numerical result

Reconstructed Higgs is different from intermediate Higgs



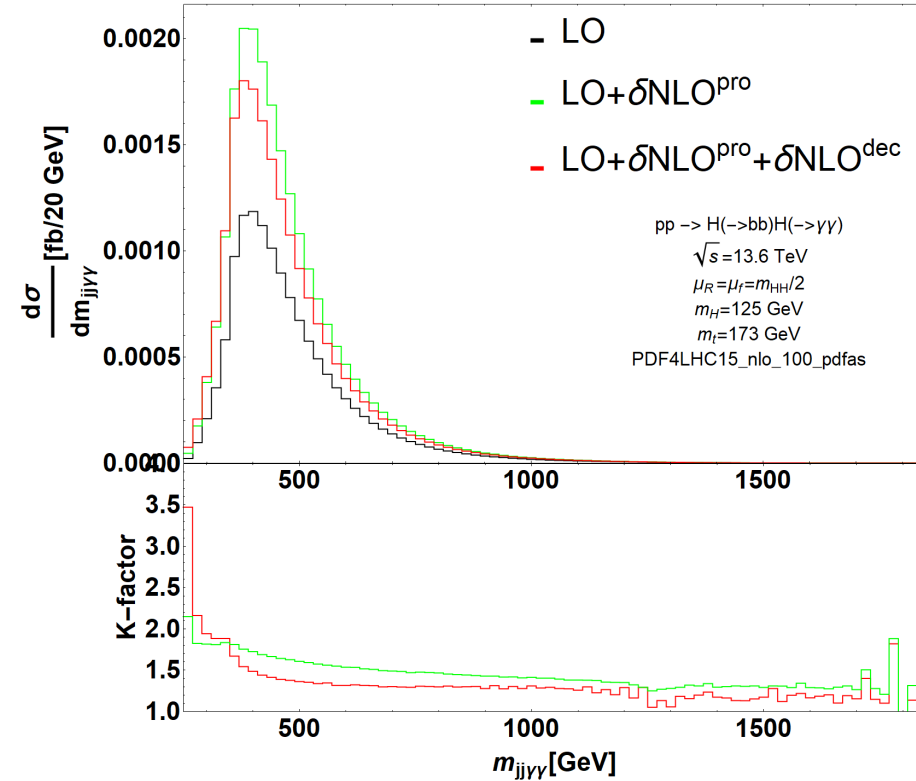
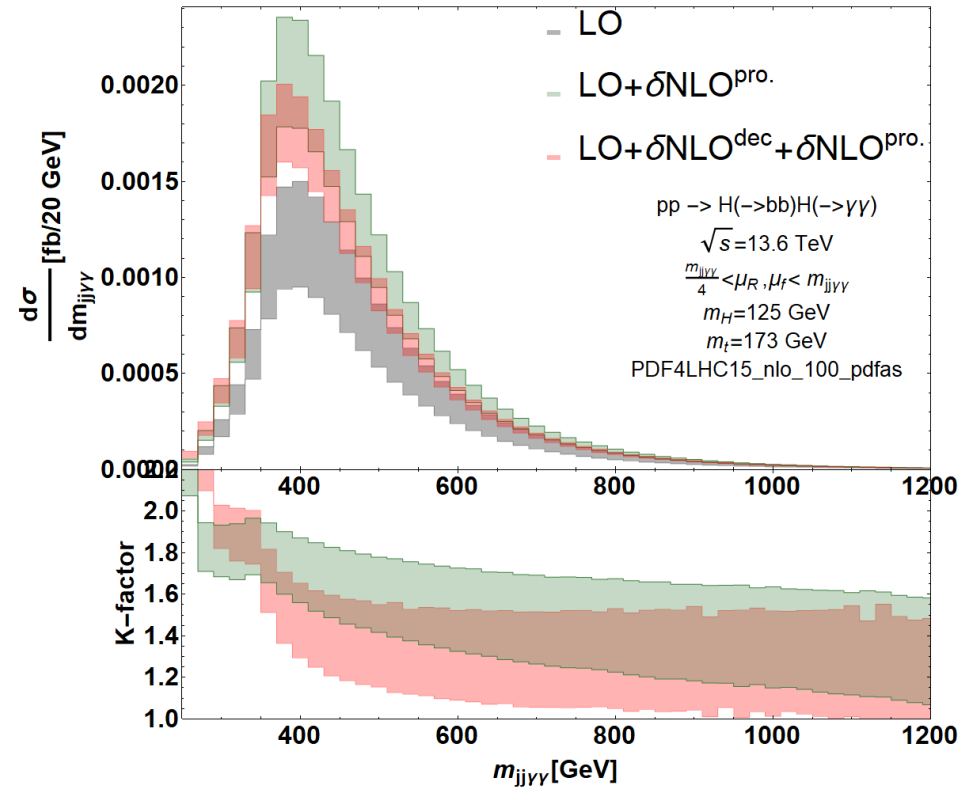
$$\begin{aligned}
 p_H^\nu &= (p_H^0, \vec{p}_T^H) & |\vec{p}_b^\nu| &= z |\vec{p}_T^H|, & |\vec{p}_{\bar{b}}^\nu| &= \bar{z} |\vec{p}_T^H| \\
 p_b^\nu &= (|\vec{p}_b^\nu|, \vec{p}_b^\nu) & & & & \\
 p_{\bar{b}}^\nu &= (|\vec{p}_{\bar{b}}^\nu|, \vec{p}_{\bar{b}}^\nu) & m_H^2 &= z\bar{z} \left(|\vec{p}_T^H| \right)^2 R^2
 \end{aligned}$$

Numerical result



In the peak region, the QCD correction is as large as about -30%.

Numerical result



- The NLO decay correction is negative but in some region is still positive.
- The K-factor is not constant.

Summary

- Higgs self-coupling λ_{HHH} is important for the SM and the best way to directly measure λ_{HHH} is from Higgs pair production.
- The experiment uncertainty will be reduced at the period of Run 3 although the constraint on κ_λ is loose up to now.
- The QCD correction in decay is significant especially for cross-sections after cut.
- $\delta\text{NLO}^{\text{dec}} \sim$ up to -30%(-19%) for distributions(total cross-section) with respect to LO

Thank you !