



Unveiling Time-Varying Signals of Ultralight Bosonic Dark Matter at Collider and Beam Dump Experiments

The 9th China LHC Physics Workshop (CLHCP2023)

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JHG, Yuxuan He , Jia Liu, Xiao-Ping Wang and Ke-Pan Xie, *Commun.Phys.* 6 (2023) 225

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- Time-varying Mass of Dark Photon
- Experiments Recast and Reanalysis
- Time-dependent Method Analysis
- Conclusion

Introduction



- Time varying physical constants Dirac, Nature 139 (1937) 323
- Ultralight dark matter

$$\text{EOM: } \ddot{\phi} + 3H\dot{\phi} + m_{\phi}^2\phi = 0$$

And misalignment mechanism

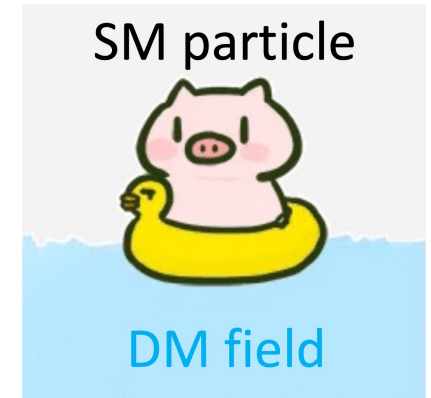
$$\phi(t) \approx \phi_0 \cos(m_{\phi}t) \quad \rho_{\text{DM}} = m_{\phi}^2\phi_0^2/2$$

- Ultralight dark matter and varying constants

Stadnik et al, 1412.7801, 1503.08540

$$\mathcal{L} \supset - \sum_{f=e,p,n} \frac{m_f}{\Lambda_f} \phi \bar{f} f + \frac{\phi}{4\Lambda_{\gamma}} F_{\mu\nu} F^{\mu\nu}$$

$$m_f \rightarrow m_f \left(1 + \frac{\phi_0 \cos(m_{\phi}t)}{\Lambda_f} \right), \quad \alpha \rightarrow \frac{\alpha}{1 - \phi_0 \cos(m_{\phi}t)/\Lambda_{\gamma}}$$

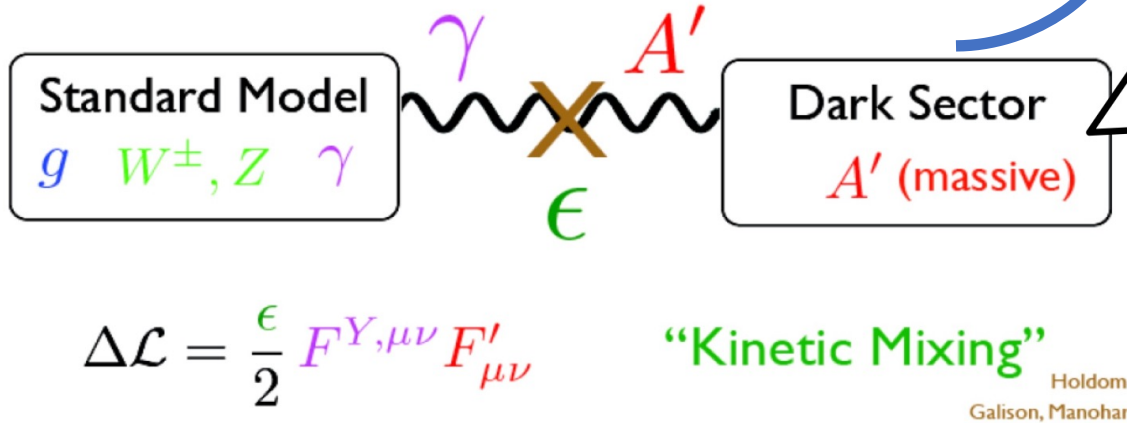


Introduction

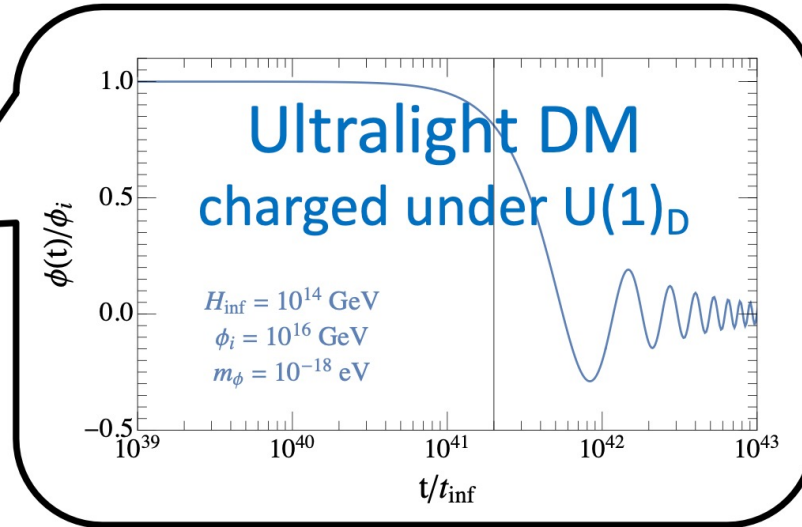


Varying dark particle mass:

The classic dark $U(1)_D$ model



Solve the $g - 2$ anomaly



$$\phi(t) \approx \phi_0 \cos(m_\phi t)$$

DM couples to A'

$$(D_\mu \phi)^* D^\mu \phi \supset (g' Q_\phi)^2 \phi^* \phi A'_\mu A'^\mu$$

Obtain a time varying mass:

$$m_{A'}(t) = \sqrt{m_0^2 + 2(g' Q_\phi)^2 |\phi_0|^2 \cos^2(m_\phi t)}$$

$$\equiv m_0 \sqrt{1 + \kappa \cos^2(m_\phi t)}$$

$$\kappa \equiv 2(g' Q_\phi)^2 \rho_{\text{DM}} / (m_\phi^2 m_0^2) = 10 \left(\frac{\rho_{\text{DM}}}{0.3 \text{ GeV/cm}^3} \right)$$

$$\times \left(\frac{g' Q_\phi}{1.5 \times 10^{-8}} \frac{10^{-19} \text{ eV}}{m_\phi} \frac{0.1 \text{ GeV}}{m_0} \right)^2,$$

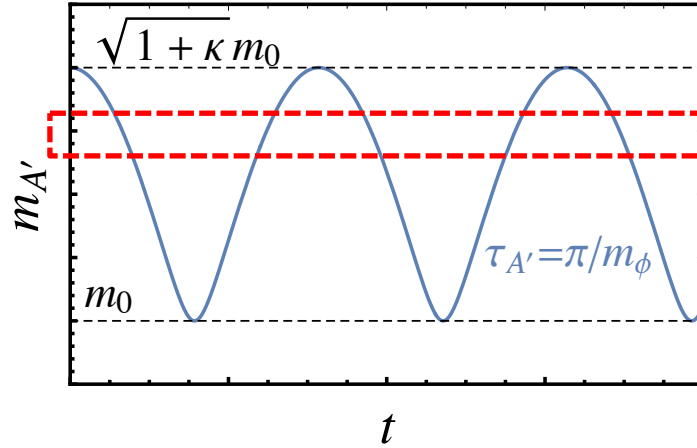
Time-varying Mass of Dark Photon

Varying dark particle mass:

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$$\equiv m_0 \sqrt{1 + \kappa \cos^2(m_\phi t)}$$

Time-varying resonant spectrum



$$m_{\text{res}}^2(t) = m_{\text{res}}^2(t + \tau)$$

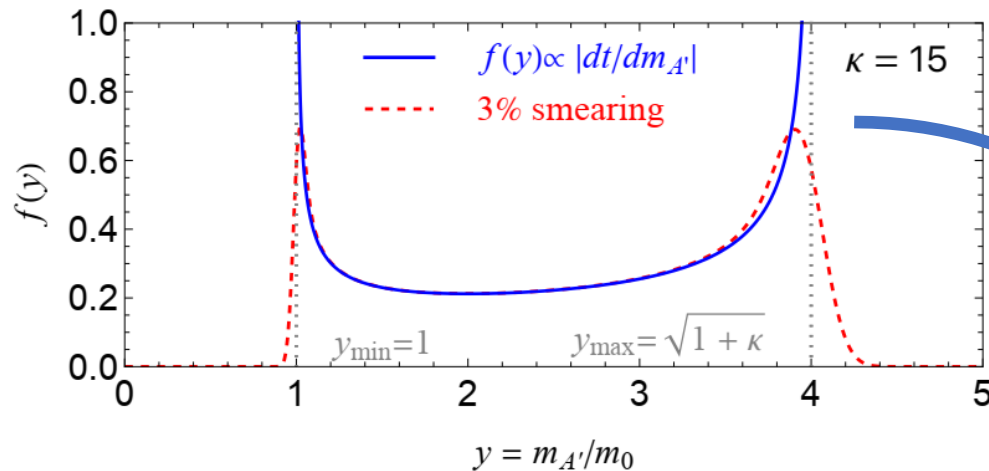
$$m_\phi \gtrsim 2 \times 10^{-20} \text{ eV}$$

$$t_{\text{exp}} \gg \tau$$

We focus on: $m_0 \sim \text{MeV}, \kappa \sim 10$

DP mass probability density function (pdf)

Number of events per bin:



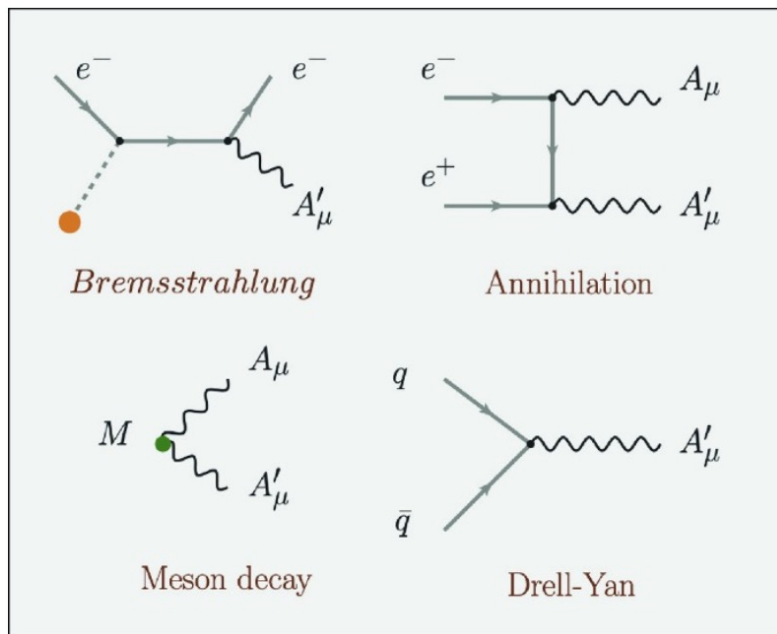
$$N_i = \sigma_{A'}^{(i)} \epsilon_i L \times \frac{1}{\tau_{A'}} \int_{m_i}^{m_{i+1}} \left| \frac{dt}{dm_{A'}} \right| dm_{A'}$$

$$\left| \frac{dt}{dm_{A'}} \right| = \frac{\tau}{m_0} f(y)$$

Double peak feature!!

Time-varying Mass of Dark Photon

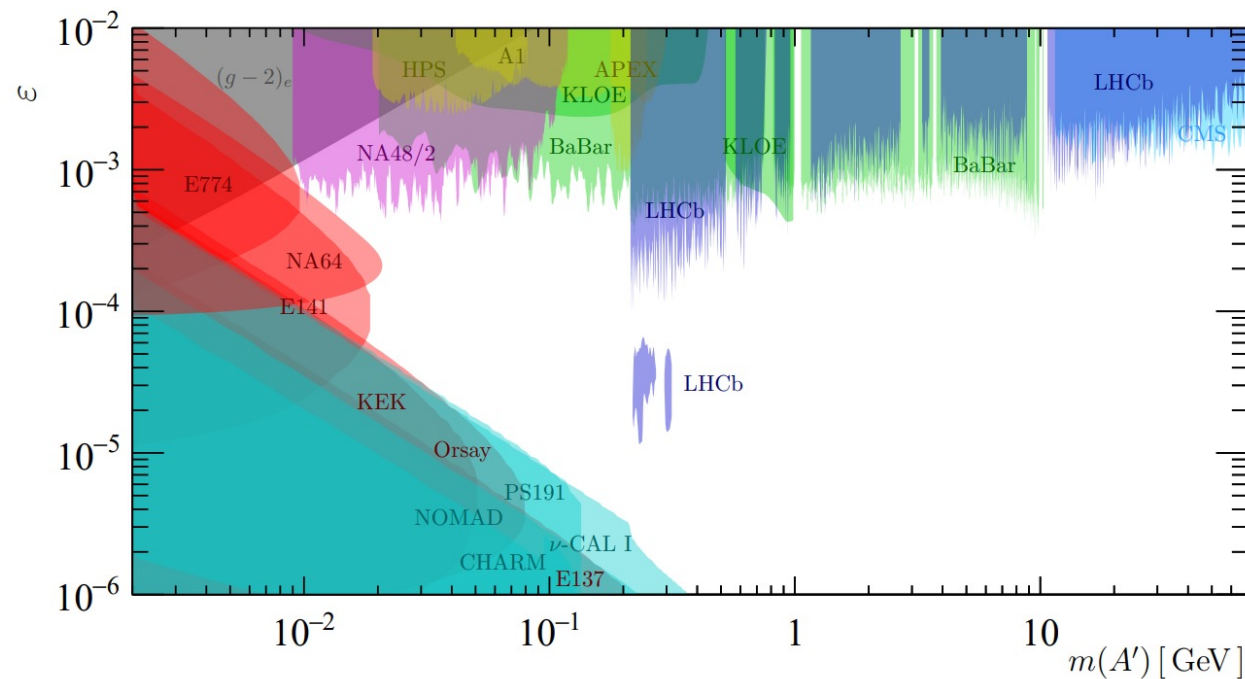
The existing constraints on dark photon mixing



$$\mathcal{L} \supset \epsilon e A'_\mu J_{em}^\mu$$

Constraints on mixing

$$\sigma_{\text{production}} \propto \epsilon^2$$



LHCb 1910.06926

Experiments Recast and Reanalysis



➤ Reanalyzing of existing experiments

Collaboration	Production mode	Experimental environment	Spectrum	Resolution σ_{re}	Fit window
BaBar [1406.2980]	$e^+e^- \rightarrow \gamma A'$	$\sqrt{s} \approx 10 \text{ GeV}, 514 \text{ fb}^{-1}$	$m_{ee}, m_{\mu\mu}$	[1.5, 8] MeV	$m_{A'} \pm 10 \sigma_{\text{re}}$
LHCb [1910.06926]	$pp \rightarrow A'$	$\sqrt{s} = 13 \text{ TeV}, \sim 5 \text{ fb}^{-1}$	$m_{\mu\mu}$	[0.12, 380] MeV	$m_{A'} \pm 12.5 \sigma_{\text{re}}$
A1 [1404.5502]	$e^-Z \rightarrow e^-ZA'$	$E_e \in [0.180, 0.855] \text{ GeV}$	m_{ee}	0.5 MeV	$m_{A'} \pm 3 \sigma_{\text{re}}$
NA48/2 [1504.00607]	$\pi^0 \rightarrow \gamma A'$	$1.69 \times 10^7 \pi^0 \rightarrow \gamma e^+e^-$ events	m_{ee}	[0.16, 1.33] MeV	single bin

➤ Log likelihood ratio (LLR)

$$-2 \log \left[\frac{\text{Max}_{\vec{a}'} \prod_i \mathcal{N}(B_i - B(m_i, \vec{a}') - S f_G(m_i) | B_i)}{\text{Max}_{\vec{a}} \prod_i \mathcal{N}(B_i - B(m_i, \vec{a}) | B_i)} \right]$$

- Background fit by $B(m_i, \vec{a}) = a_0 + a_1 m_i + a_2 m_i^2$
- Constraints obtained by $S \equiv -2 \ln(\mathcal{L}/\mathcal{L}_0) = 3.84$

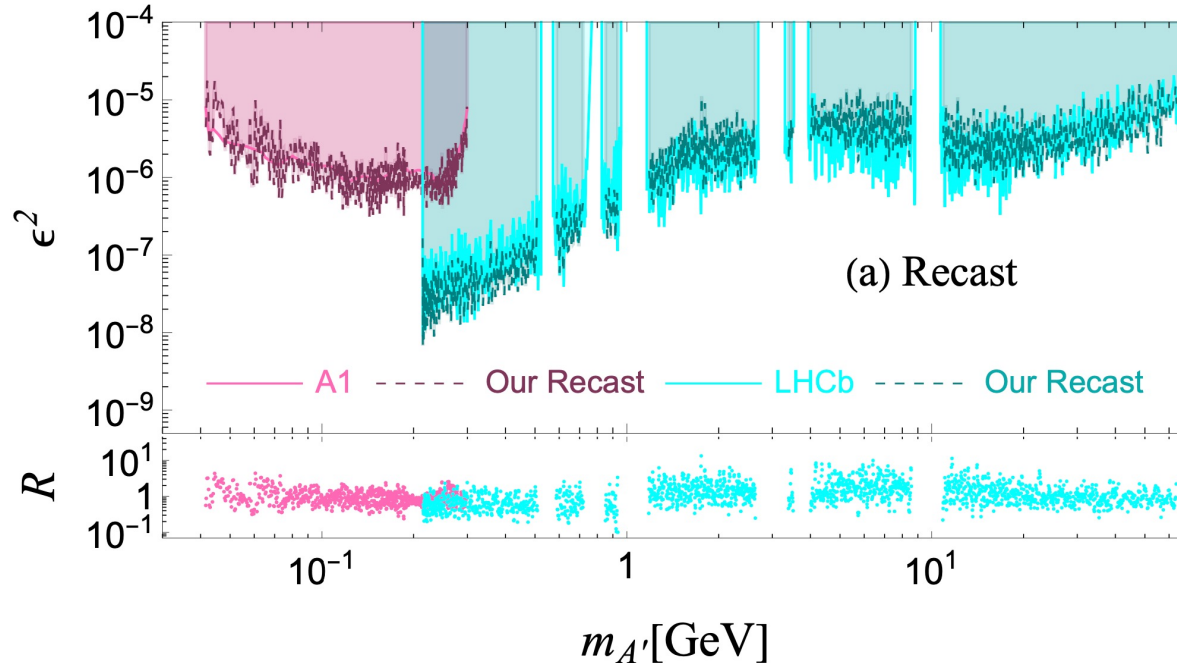
Experiments Recast and Reanalysis



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First, repeat the existing bounds with the single-peak signal



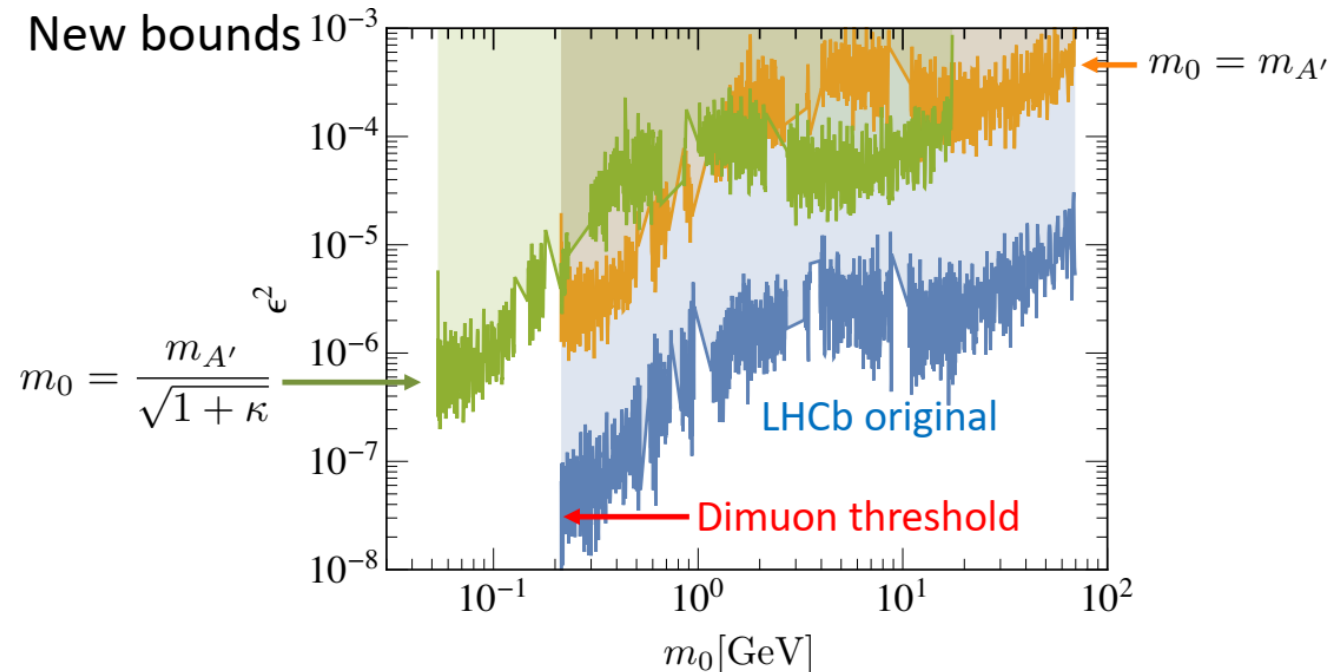
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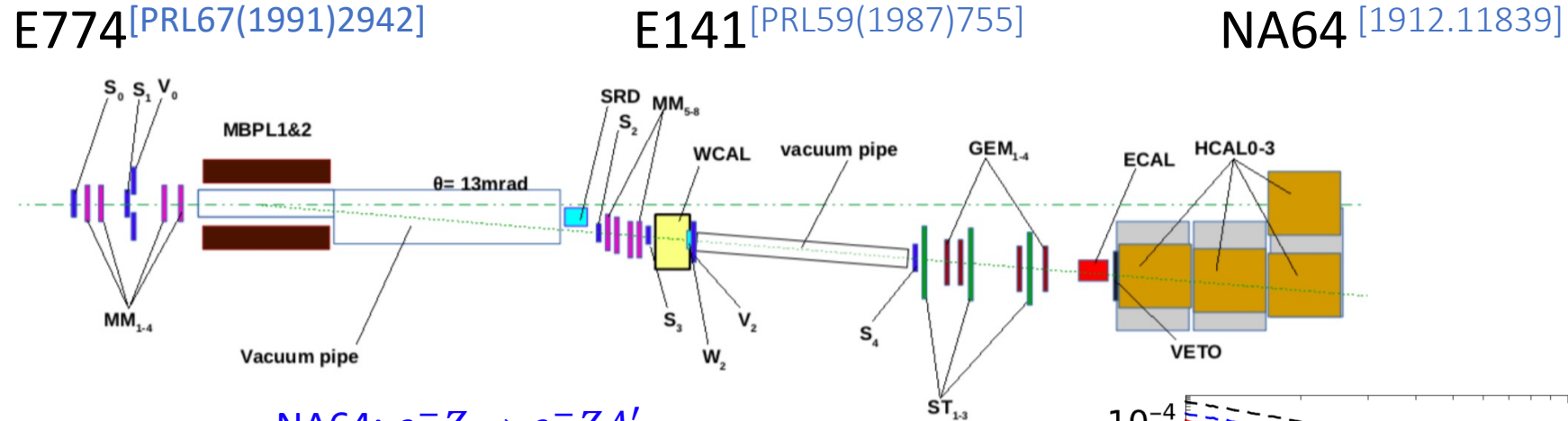
Then, obtain the new bounds:



Experiments Recast and Reanalysis



➤ Reanalyzing the beam dump experiments

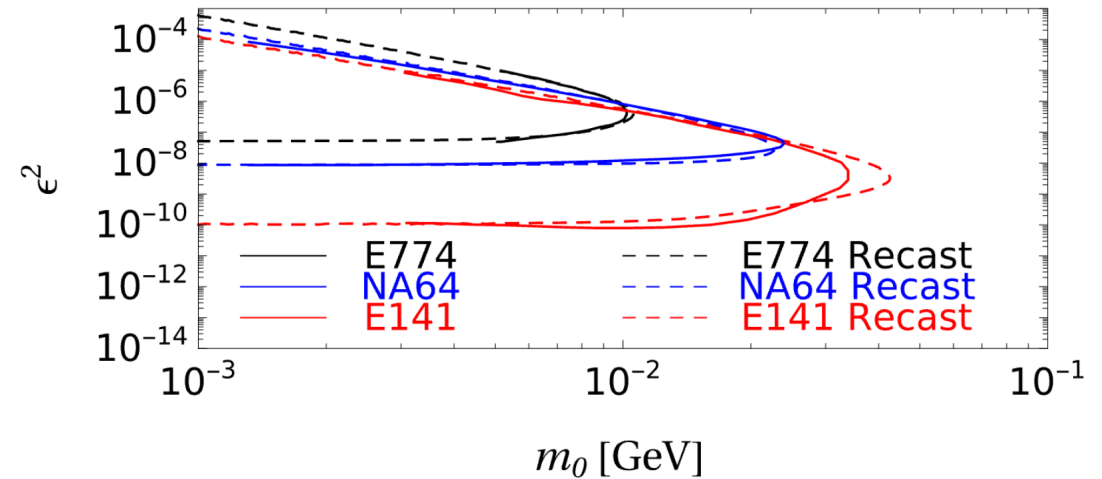


Event number:

$$N(\epsilon, m_{A'}) = N_e C' \epsilon^2 \frac{m_e^2}{m_{A'}^2} e^{-a_1 L_{sh} \Gamma_{A'}} \times (1 - e^{-a_2 L_{dec} \Gamma_{A'}}),$$

For our signals:

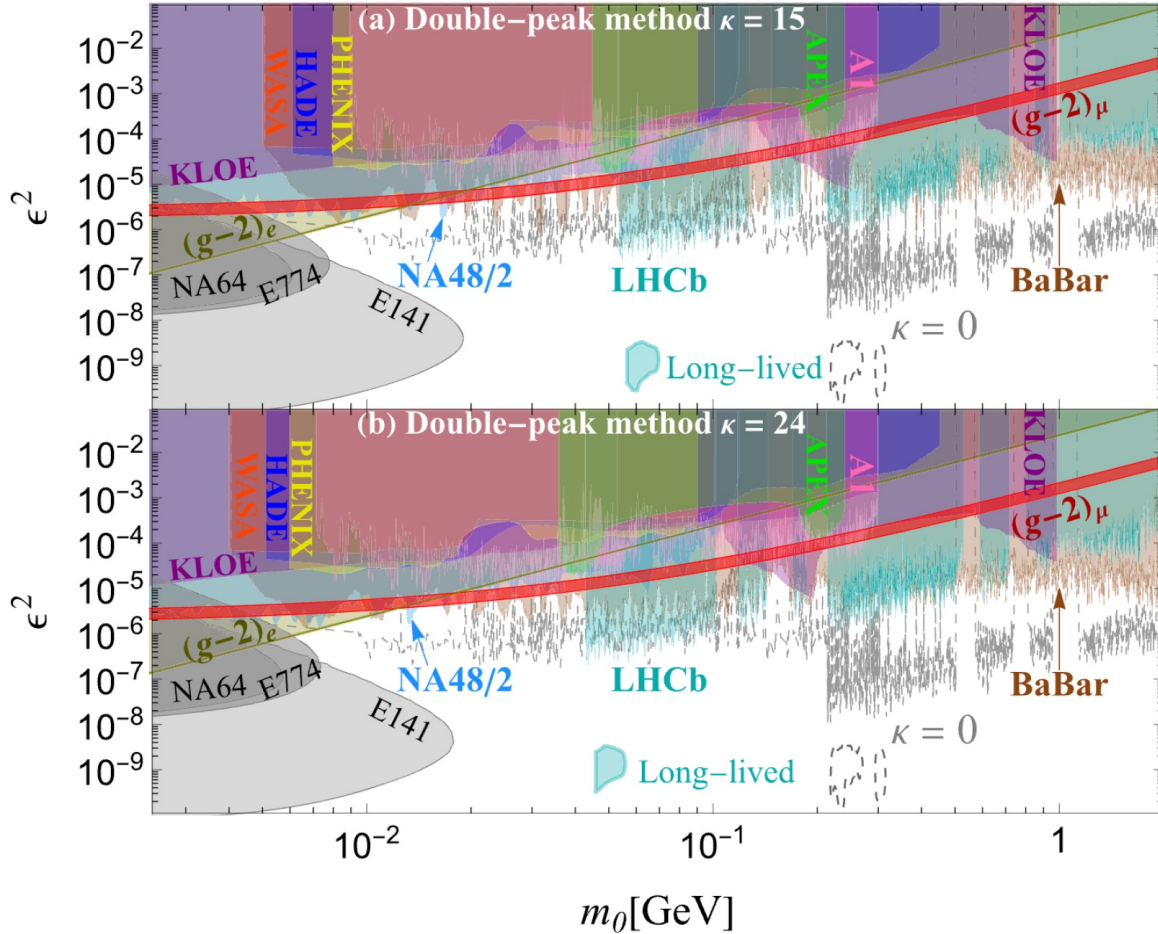
$$N(\epsilon, m_0, \kappa) = \frac{1}{\tau_{A'}} \int_{m_0}^{\sqrt{1+\kappa} m_0} N(\epsilon, m_{A'}) \left| \frac{dt}{dm_{A'}} \right| dm_{A'}.$$



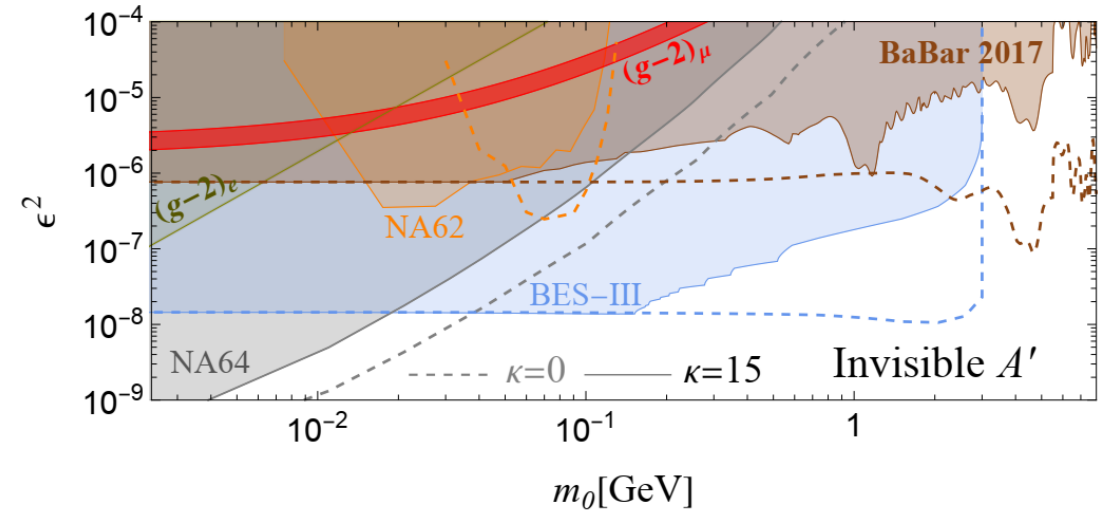
Experiments Recast and Reanalysis



➤ Obtain the new constraints:



➤ Invisible Dark Photon Searches ($A' \rightarrow \chi\bar{\chi}$)



- Monophoton search ($e^+e^- \rightarrow A'\gamma$):

$$E_\gamma = \frac{s - m_{A'}^2}{2\sqrt{s}}$$

- Time-varying photon energy spectrum:

$$\left| \frac{dt}{dE_\gamma} \right| = \frac{\tau}{\pi \sqrt{(E_\gamma - E_{\min})(E_{\max} - E_\gamma)}}$$

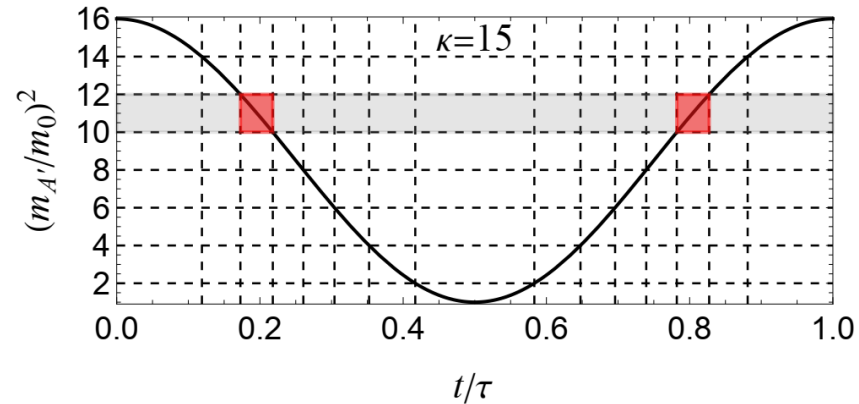
And the $g - 2$: $\frac{1}{\tau} \int_0^\tau dt \Delta a_\mu(m_{A'}(t))$

Time-dependent Method Analysis

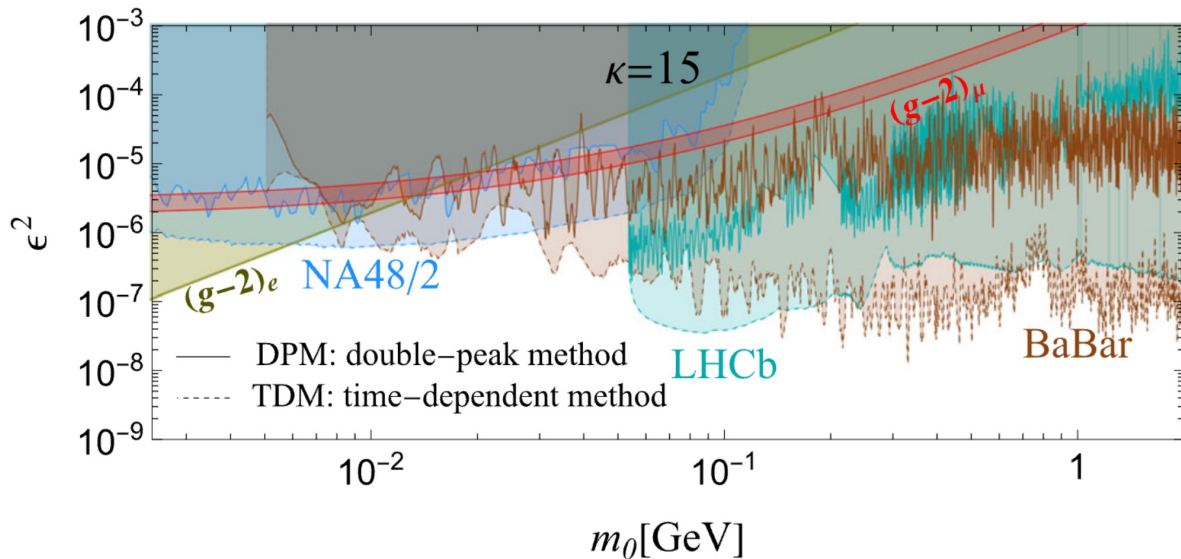
➤ Using the time information:

$$N_i^{\text{red}} = N_i \frac{1}{\tau} \int_{m_i}^{m_{i+1}} \left| \frac{dt}{dm_{A'}} \right| dm_{A'}$$

$$N_{i,j} = N_i \Delta t_j / \tau$$



- Signal not changed, while background suppressed!



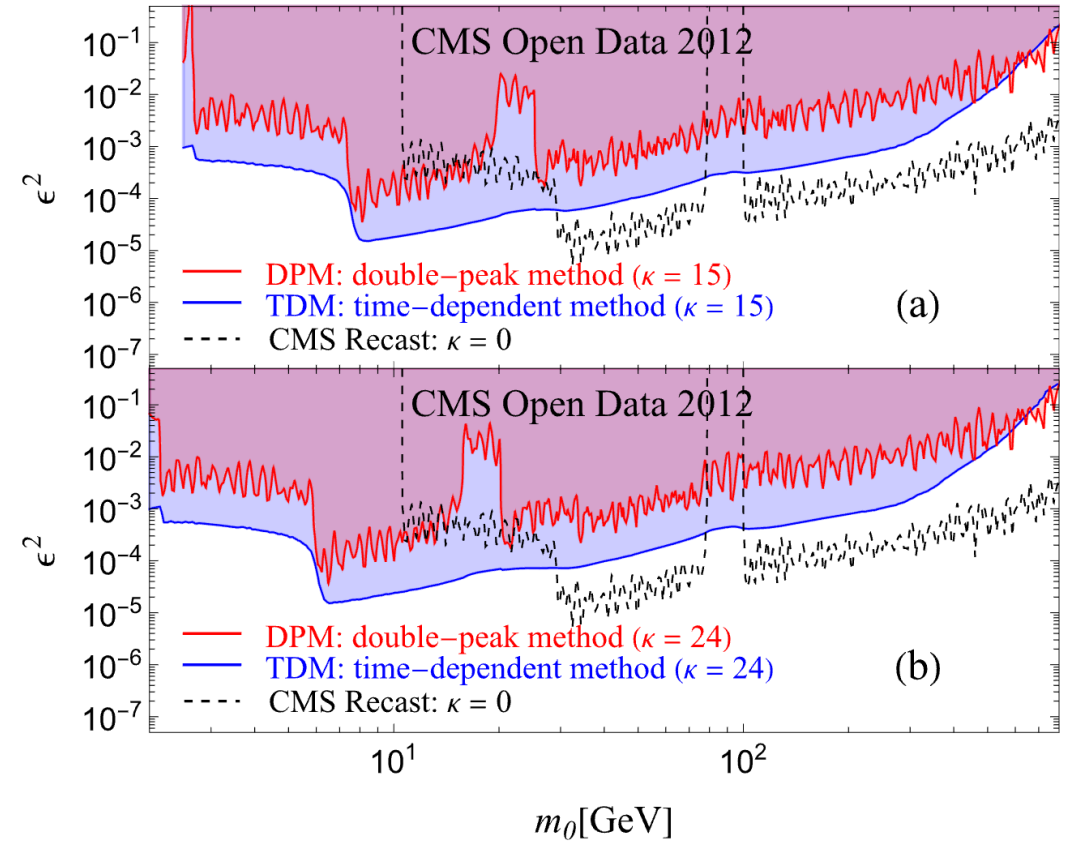
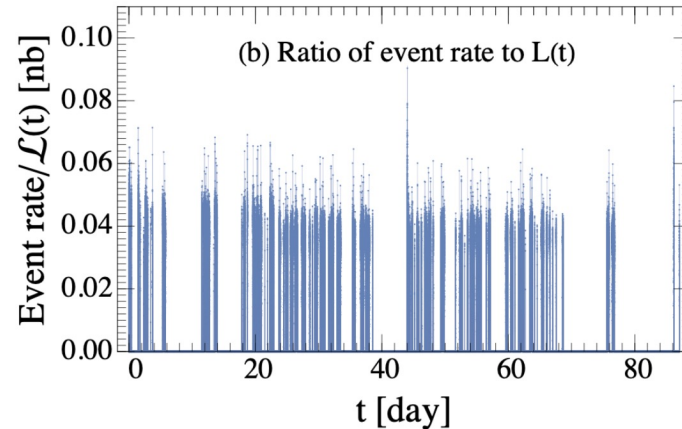
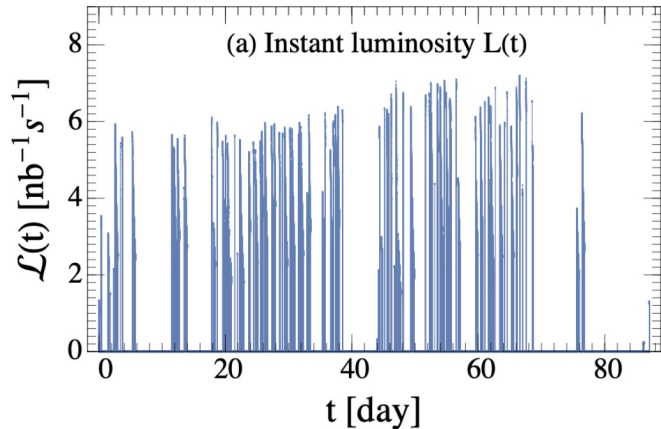
➔ TDM has better constraint than DPM

Time-dependent Method Analysis

➤ Analyzing CMS Open Data

We choose 2012 dimuon events from CMS Open Data to verify the two methods: DPM and TDM

$$\begin{aligned} \frac{dN_S}{dm_{\ell\ell}} &= \int_{t_1}^{t_2} dt \mathcal{L}(t) \frac{d\sigma_S}{dm_{\ell\ell}} \\ &= \epsilon_S(m_{\ell\ell}) \sigma_0(m_{\ell\ell}) \frac{f(y)}{m_0} \frac{\tau}{2} \left[\sum_i \mathcal{L}(t_i^+) + \mathcal{L}(t_i^-) \right], \\ \bar{f}(m_{\ell\ell}/m_0) &\equiv \frac{1}{\int \mathcal{L}(t) dt} \frac{dN_S}{d(m_{\ell\ell}/m_0)} \\ &= f(y) \frac{\tau}{2L} \left[\sum_i \mathcal{L}(t_i^+) + \mathcal{L}(t_i^-) \right]. \end{aligned}$$



Conclusion



Dark photon can have a time varying mass by coupling with ultralight scalar DM:

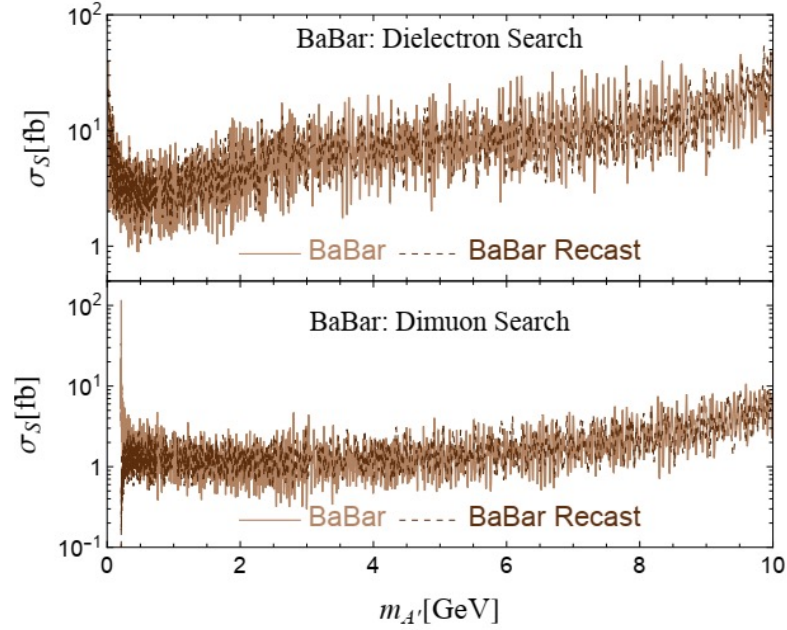
- The existing bounds are significantly weakened;
- The muon $g-2$ solution becomes viable again;
- Including time information of experiments can improve sensitivity;
- CMS Open Data shows the validity of these two methods;
- If we find this time-varying signature, it can be used to unveil the ultralight DM.

Thank You

Backup



➤ Repeating of existing constraints



➤ Other constraints

Varying SM fermion mass

$$\frac{\Delta m_f}{m_f} \simeq \frac{3 (e\epsilon Q_f)^2}{16\pi^2} \log \left(\frac{m_0^2 + 2(g' Q_\phi)^2 \phi^* \phi}{m_0^2} \right)$$

Early universe:

1. Thermalization: small coupling $g' Q_\phi \sim 10^{-6} - 10^{-10}$, $m_\phi \sim 10^{-20} - 10^{-17}$ eV, $\kappa \sim 10$; frozen scalar field gives a large DP mass;
2. Black hole super-radiance.