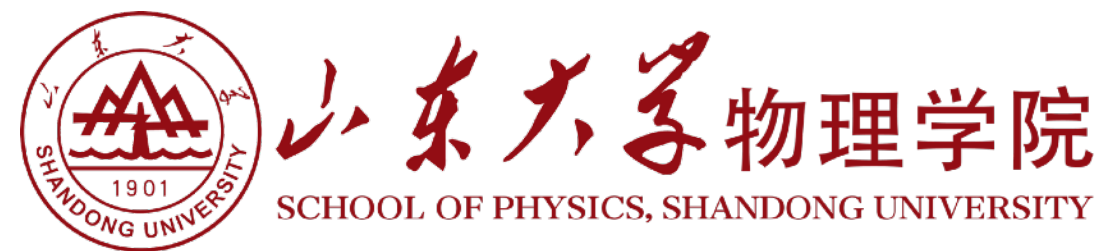


Theory overview: QCD

Haitao Li 李海涛



The 9th China LHC Physics Workshop (CLHCP2023)

2023-11-17

Asymptotic Freedom

The Nobel Prize in Physics 2004

50 years of QCD

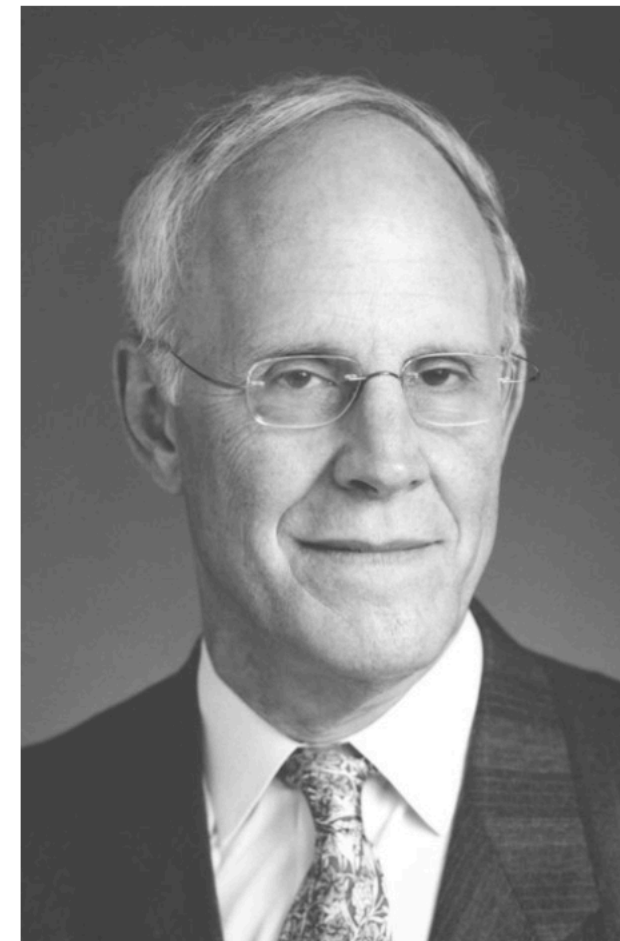


Photo from the Nobel Foundation archive.

David J. Gross

Prize share: 1/3



Photo from the Nobel Foundation archive.

H. David Politzer

Prize share: 1/3

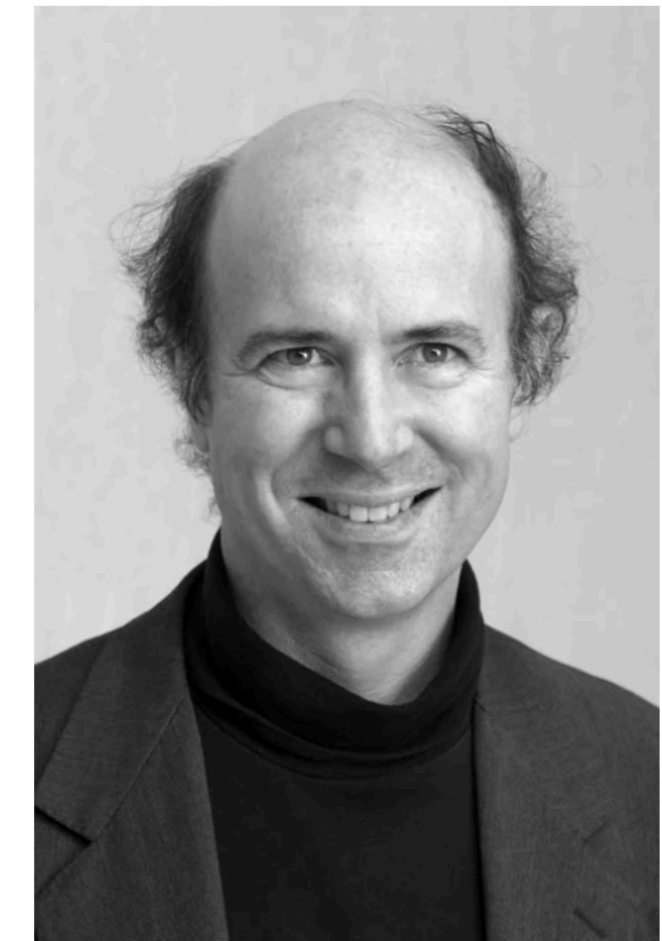


Photo from the Nobel Foundation archive.

Frank Wilczek

Prize share: 1/3

The Nobel Prize in Physics 2004 was awarded jointly to David J. Gross, H. David Politzer and Frank Wilczek "for the discovery of asymptotic freedom in the theory of the strong interaction"

Asymptotic Freedom

QCD beta function

$$\frac{d\alpha_s}{d\ln\mu^2} = \beta(\alpha_s) = -\alpha_s \left(\frac{\alpha_s}{\pi} \beta_0 + \left(\frac{\alpha_s}{\pi} \right)^2 \beta_1 + \left(\frac{\alpha_s}{\pi} \right)^3 \beta_2 + \left(\frac{\alpha_s}{\pi} \right)^4 \beta_3 + \left(\frac{\alpha_s}{\pi} \right)^5 \beta_4 + \dots \right)$$

$$\beta_0 = \frac{1}{4} \left\{ 11 - \frac{2}{3} n_f \right\}, \quad \beta_1 = \frac{1}{4^2} \left\{ 102 - \frac{38}{3} n_f \right\},$$

$$\beta_2 = \frac{1}{4^3} \left\{ \frac{2857}{2} - \frac{5033}{18} n_f + \frac{325}{54} n_f^2 \right\},$$

$$\beta_3 = \frac{1}{4^4} \left\{ \frac{149753}{6} + 3564 \zeta_3 - \left[\frac{1078361}{162} + \frac{6508}{27} \zeta_3 \right] n_f + \left[\frac{50065}{162} + \frac{6472}{81} \zeta_3 \right] n_f^2 + \frac{1093}{729} n_f^3 \right\},$$

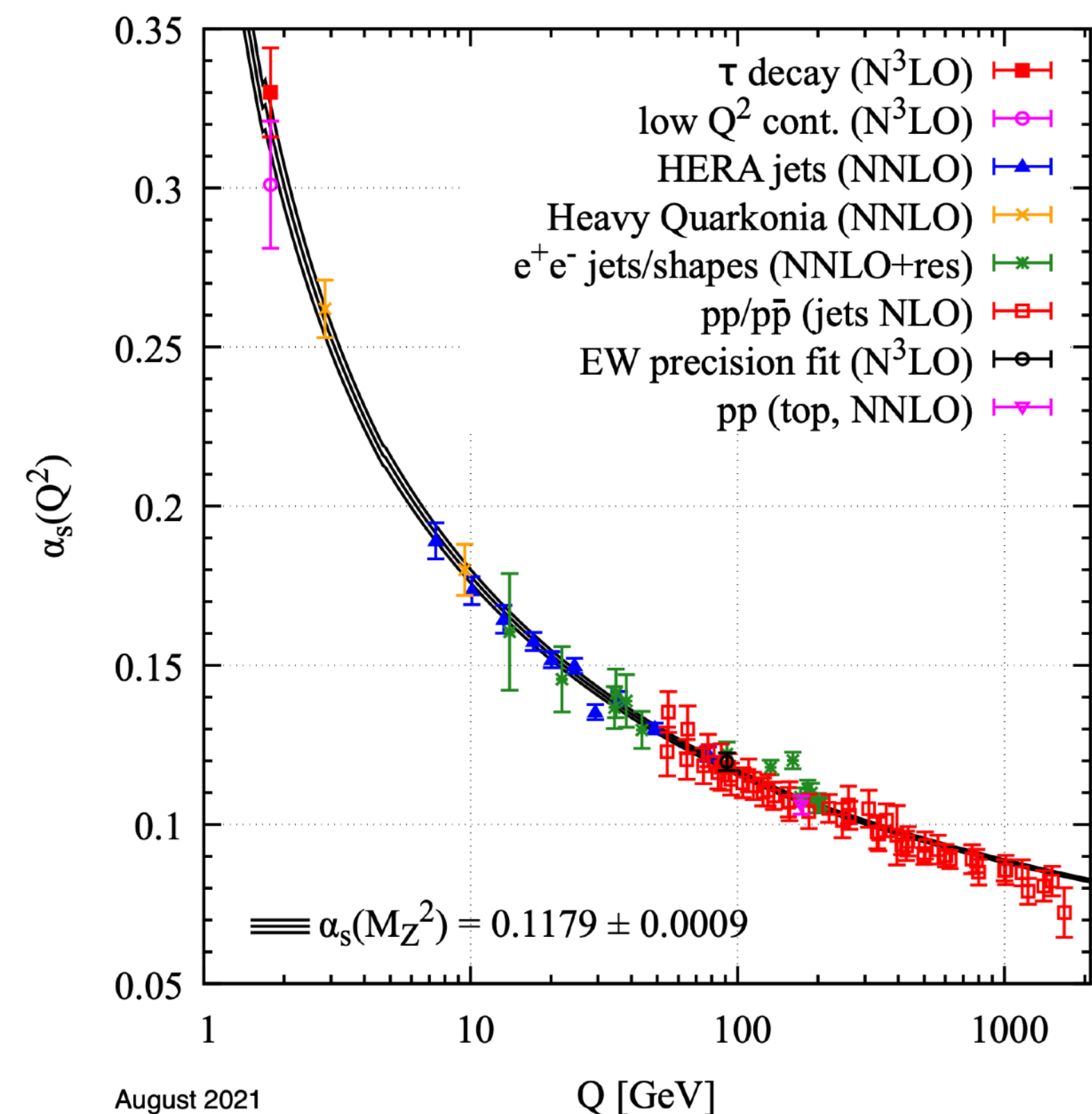
$$\begin{aligned} \beta_4 = & \frac{1}{4^5} \left\{ \frac{8157455}{16} + \frac{621885}{2} \zeta_3 - \frac{88209}{2} \zeta_4 - 288090 \zeta_5 \right. \\ & + n_f \left[-\frac{336460813}{1944} - \frac{4811164}{81} \zeta_3 \right. \\ & \quad \left. \left. + \frac{33935}{6} \zeta_4 + \frac{1358995}{27} \zeta_5 \right] \right. \\ & + n_f^2 \left[\frac{25960913}{1944} + \frac{698531}{81} \zeta_3 - \frac{10526}{9} \zeta_4 - \frac{381760}{81} \zeta_5 \right] \\ & + n_f^3 \left[-\frac{630559}{5832} - \frac{48722}{243} \zeta_3 + \frac{1618}{27} \zeta_4 + \frac{460}{9} \zeta_5 \right] \\ & \left. + n_f^4 \left[\frac{1205}{2916} - \frac{152}{81} \zeta_3 \right] \right\}, \end{aligned}$$

Baikov, Chetyrkin, Kuhn, 2016

'73 '74 '80 '97 2016

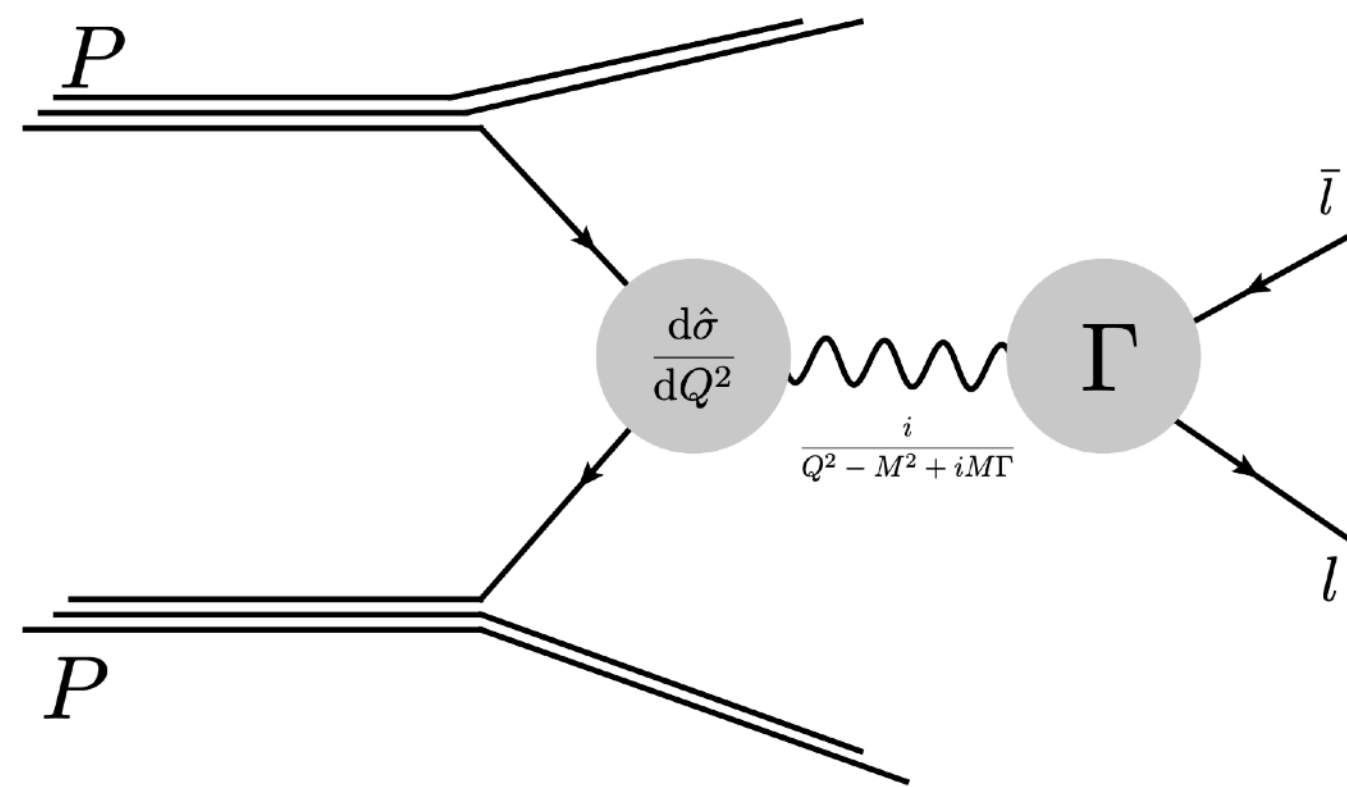
4-loop 5-loop

5000 Diagrams 1.5 Million Diagrams



QCD and LHCp

Factorize the physics associated with different scales



Hard process: where energy scale is large and perturbative theory works well.

Evolution: Resummation. A bridge.

Hadronization: Non-perturbative

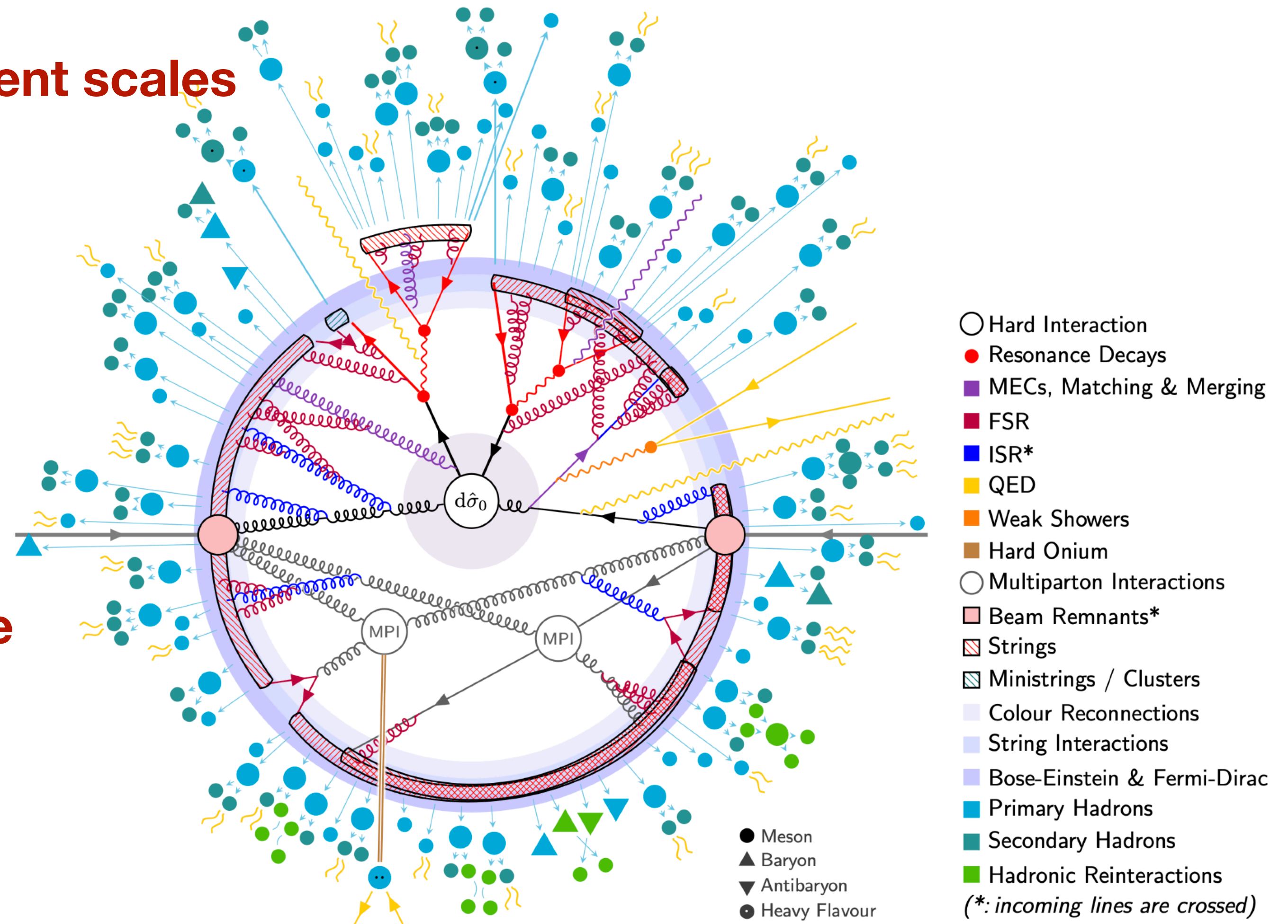


figure taken from PYTHIA 8.3 manual

Fixed order

$$\frac{d\sigma}{d\mathcal{O}} = \frac{1}{2s} \sum_{a,b} \left[\int \frac{dx_1}{x_1} f_a(x_1) \int \frac{dx_2}{x_2} f_b(x_2) \right] \left[\Pi_i \frac{d^3 \mathbf{p}_i}{(2\pi)^3 2E_i} (2\pi)^4 \delta(p_a + p_b - \sum_i p_i) \left| \mathcal{M}(a + b \rightarrow \{i\}) \right|^2 \delta(\mathcal{O} - \mathcal{O}(\{p_i\})) \right]$$

Theoretical Predictions rely on

PDFs and

Partonic Cross Section

Fixed order

$$\frac{d\sigma}{d\mathcal{O}} = \frac{1}{2s} \sum_{a,b} \left[\int \frac{dx_1}{x_1} f_a(x_1) \int \frac{dx_2}{x_2} f_b(x_2) \right] \left[\Pi_i \frac{d^3 \mathbf{p}_i}{(2\pi)^3 2E_i} (2\pi)^4 \delta(p_a + p_b - \sum_i p_i) \left| \mathcal{M}(a + b \rightarrow \{i\}) \right|^2 \delta(\mathcal{O} - \mathcal{O}(\{p_i\})) \right]$$

Theoretical Predictions rely on

PDFs and

Partonic Cross Section

PDFs

Non-Perturbative
Process Independent

CT, MSHT, NNPDF, PDF4LHC et al

reviews: Gao, Harland-Lang Rojo, Phys.Rep. 2018
Snowmass 2021 Whitepaper about PDF, arXiv:2203.13923

Fixed order

$$\frac{d\sigma}{d\mathcal{O}} = \frac{1}{2s} \sum_{a,b} \left[\int \frac{dx_1}{x_1} f_a(x_1) \int \frac{dx_2}{x_2} f_b(x_2) \right] \left[\Pi_i \frac{d^3 \mathbf{p}_i}{(2\pi)^3 2E_i} (2\pi)^4 \delta(p_a + p_b - \sum_i p_i) \left| \mathcal{M}(a + b \rightarrow \{i\}) \right|^2 \delta(\mathcal{O} - \mathcal{O}(\{p_i\})) \right]$$

Theoretical Predictions rely on

PDFs and

Partonic Cross Section

PDFs

**Non-Perturbative
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CT, MSHT, NNPDF, PDF4LHC et al

reviews: Gao, Harland-Lang Rojo, Phys.Rep. 2018
Snowmass 2021 Whitepaper about PDF, arXiv:2203.13923

NNLO PDFs require 3-loop splitting function

hep-ph/0403192, hep-ph/0404111

Fixed order

$$\frac{d\sigma}{d\mathcal{O}} = \frac{1}{2s} \sum_{a,b} \left[\int \frac{dx_1}{x_1} f_a(x_1) \int \frac{dx_2}{x_2} f_b(x_2) \right] \left[\Pi_i \frac{d^3 \mathbf{p}_i}{(2\pi)^3 2E_i} (2\pi)^4 \delta(p_a + p_b - \sum_i p_i) \left| \mathcal{M}(a + b \rightarrow \{i\}) \right|^2 \delta(\mathcal{O} - \mathcal{O}(\{p_i\})) \right]$$

Theoretical Predictions rely on PDFs and Partonic Cross Section

PDFs

**Non-Perturbative
Process Independent**
CT, MSHT, NNPDF, PDF4LHC et al

NNLO PDFs require 3-loop splitting function

**To get N3LO PDFs four loop QCD splitting
function required**

available only for some specific color structures

reviews: Gao, Harland-Lang Rojo, Phys.Rep. 2018
Snowmass 2021 Whitepaper about PDF, arXiv:2203.13923

hep-ph/0403192, hep-ph/0404111

arXiv:1610.07477, arXiv:1707.08315, arXiv:2307.04158,
arXiv:2308.07958, arXiv:2310.01245, arXiv:2310.05744,
arXiv:2310.12240

Fixed order

$$\frac{d\sigma}{d\mathcal{O}} = \frac{1}{2s} \sum_{a,b} \left[\int \frac{dx_1}{x_1} f_a(x_1) \int \frac{dx_2}{x_2} f_b(x_2) \right] \left[\Pi_i \frac{d^3 \mathbf{p}_i}{(2\pi)^3 2E_i} (2\pi)^4 \delta(p_a + p_b - \sum_i p_i) \left| \mathcal{M}(a + b \rightarrow \{i\}) \right|^2 \delta(\mathcal{O} - O(\{p_i\})) \right]$$

Theoretical Predictions rely on PDFs and Partonic Cross Section

PDFs	Non-Perturbative Process Independent <i>CT, MSHT, NNPDF, PDF4LHC et al</i>	reviews: Gao, Harland-Lang Rojo, Phys.Rep. 2018 Snowmass 2021 Whitepaper about PDF, arXiv:2203.13923
	NNLO PDFs require 3-loop splitting function	hep-ph/0403192, hep-ph/0404111
	To get N3LO PDFs four loop QCD splitting function required	arXiv:1610.07477, arXiv:1707.08315, arXiv:2307.04158, arXiv:2308.07958, arXiv:2310.01245, arXiv:2310.05744,
	available only for some specific color structures	arXiv:2310.12240
	Some Implementation about approximate N3LO PDFs	arXiv:2207.04739 (MSHT20aN ³ LO), arXiv:2306.15294 (NNPDF)

Fixed order

$$\frac{d\sigma}{d\mathcal{O}} = \frac{1}{2s} \sum_{a,b} \left[\int \frac{dx_1}{x_1} f_a(x_1) \int \frac{dx_2}{x_2} f_b(x_2) \right] \left[\Pi_i \frac{d^3\mathbf{p}_i}{(2\pi)^3 2E_i} (2\pi)^4 \delta(p_a + p_b - \sum_i p_i) \left| \mathcal{M}(a + b \rightarrow \{i\}) \right|^2 \delta(\mathcal{O} - \mathcal{O}(\{p_i\})) \right]$$

Theoretical Predictions rely on PDFs and Partonic Cross Section

Non-Perturbative

reviews: [Gao, Harland-Lano, Roio, Phys.Rev. 2018](#)

σ order		PDF order	$\sigma + \Delta\sigma_+ - \Delta\sigma_-$ (pb)	σ (pb) + $\Delta\sigma_+ - \Delta\sigma_-$ (%)
PDF uncertainties				
P	N ³ LO	aN ³ LO (no theory unc.)	4.1150 + 0.0638 − 0.0724	4.1150 + 1.55% − 1.76%
		aN ³ LO ($H_{ij} + K_{ij}$)	4.1150 + 0.0682 − 0.0755	4.1150 + 1.66% − 1.83%
		aN ³ LO (H'_{ij})	4.1150 + 0.0678 − 0.0742	4.1150 + 1.65% − 1.80%
		NNLO	3.9941 + 0.0558 − 0.0631	3.9941 + 1.40% − 1.58%
	NNLO	NNLO	3.9974 + 0.0557 − 0.0633	3.9974 + 1.39% − 1.58%

7.04158,
0.05744,

Some Implementation about approximate N3LO PDFs

[arXiv:2207.04739](#) (MSHT20aN³LO),
[arXiv:2306.15294](#) (NNPDF)

Fixed order

$$\frac{d\sigma}{d\mathcal{O}} = \frac{1}{2s} \sum_{a,b} \left[\int \frac{dx_1}{x_1} f_a(x_1) \int \frac{dx_2}{x_2} f_b(x_2) \right] \left[\Pi_i \frac{d^3 \mathbf{p}_i}{(2\pi)^3 2E_i} (2\pi)^4 \delta(p_a + p_b - \sum_i p_i) \left| \mathcal{M}(a + b \rightarrow \{i\}) \right|^2 \delta(\mathcal{O} - \mathcal{O}(\{p_i\})) \right]$$

Theoretical Predictions rely on

PDFs and
Universal

Partonic Cross Section

**Perturbative but
Process Dependent**

The hadronic cross section

$$\sigma = \sigma_0 \times \left(1 + \frac{\alpha_s}{4\pi} \kappa_1 + \left(\frac{\alpha_s}{4\pi} \right)^2 \kappa_2 + \left(\frac{\alpha_s}{4\pi} \right)^3 \kappa_3 + \dots \right)$$

LO

NLO

NNLO

N3LO

NLO QCD corrections:

well studied method
with automatic tools

[arXiv:1405.0301](https://arxiv.org/abs/1405.0301)

However,

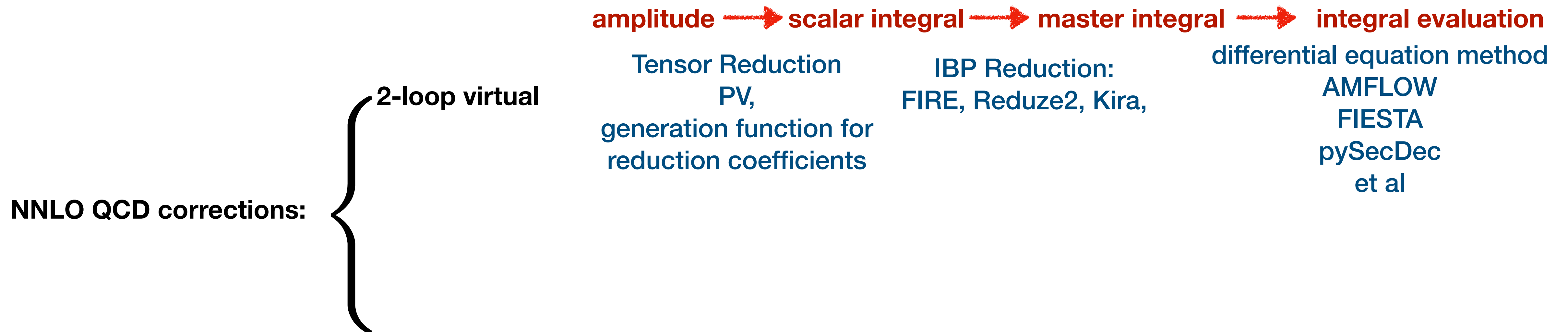
1. QCD corrections to loop induced process (virtual correction corresponds to 2loop)
2. NLO EW corrections
3. NLO QCD corrections for BSM
- ...

see talks by Dan Zhao and Xiao Zhang

Fixed order

$$\sigma = \sigma_0 \times \left(1 + \frac{\alpha_s}{4\pi} \kappa_1 + \left(\frac{\alpha_s}{4\pi} \right)^2 \kappa_2 + \left(\frac{\alpha_s}{4\pi} \right)^3 \kappa_3 + \dots \right)$$

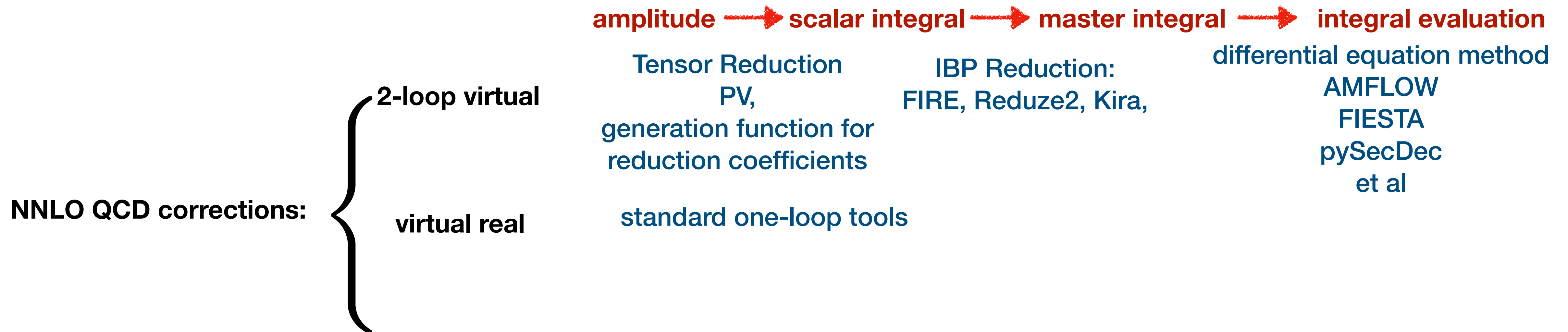
LO NLO NNLO N3LO



Fixed order

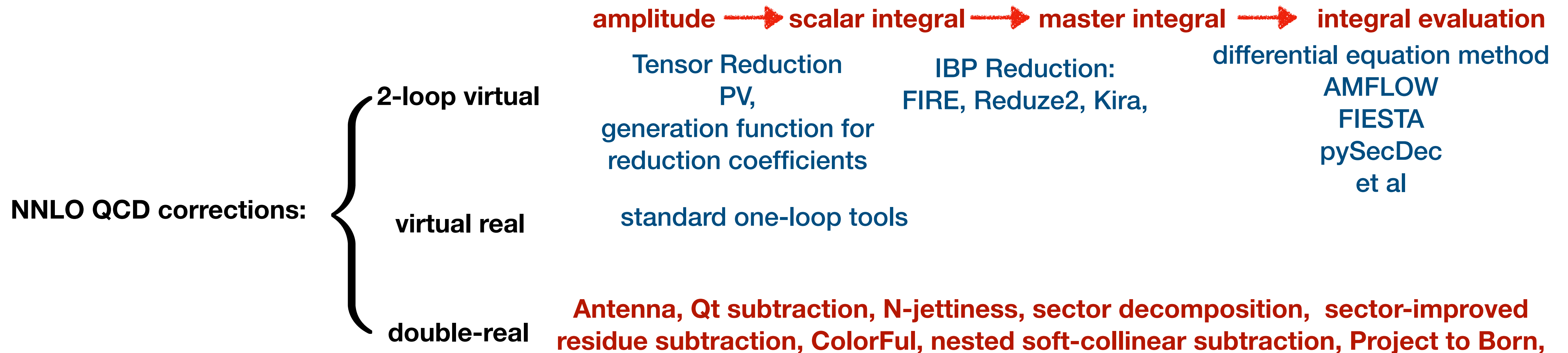
$$\sigma = \sigma_0 \times \left(1 + \frac{\alpha_s}{4\pi} \kappa_1 + \left(\frac{\alpha_s}{4\pi} \right)^2 \kappa_2 + \left(\frac{\alpha_s}{4\pi} \right)^3 \kappa_3 + \dots \right)$$

LO NLO NNLO N3LO



Fixed order

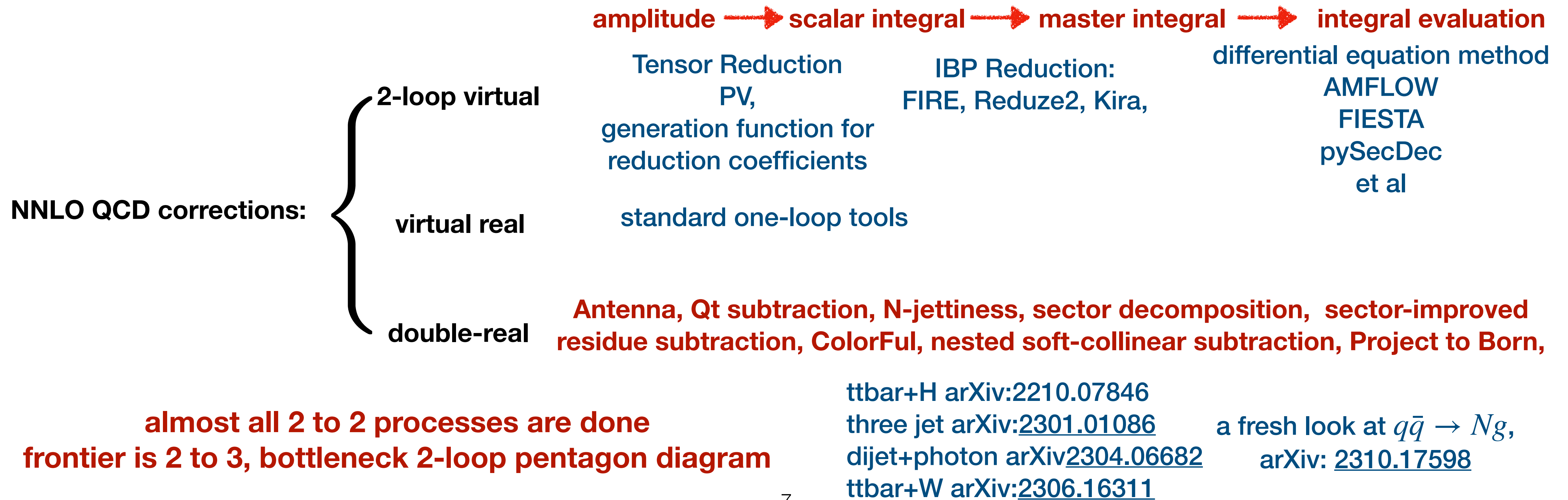
$$\sigma = \sigma_0 \times \left(1 + \underbrace{\frac{\alpha_s}{4\pi} \kappa_1}_{\text{LO}} + \underbrace{\left(\frac{\alpha_s}{4\pi} \right)^2 \kappa_2}_{\text{NLO}} + \underbrace{\left(\frac{\alpha_s}{4\pi} \right)^3 \kappa_3}_{\text{NNLO}} + \dots \right)$$



Fixed order

$$\sigma = \sigma_0 \times \left(1 + \frac{\alpha_s}{4\pi} \kappa_1 + \left(\frac{\alpha_s}{4\pi} \right)^2 \kappa_2 + \left(\frac{\alpha_s}{4\pi} \right)^3 \kappa_3 + \dots \right)$$

LO NLO NNLO N3LO



Fixed order

Hard interactions

$$\sigma = \sigma_0 \times \left(\underbrace{1}_{\text{LO}} + \underbrace{\frac{\alpha_s}{4\pi} \kappa_1}_{\text{NLO}} + \underbrace{\left(\frac{\alpha_s}{4\pi} \right)^2 \kappa_2}_{\text{NNLO}} + \underbrace{\left(\frac{\alpha_s}{4\pi} \right)^3 \kappa_3}_{\text{N3LO}} + \dots \right)$$

gg2H

inclusive '15, differential '21

ff2H

inclusive '19, differential '23

DY

inclusive '20, differential 22

WH

Baglio, Duhr, Mistlberger, Szafron, 2022

NNLO QCD corrections:

HH

Chen, HTL, et al, 2019

Hjj (VBF)

Dreyer, Karlberg, 2016

HHjj(VBF)

Dreyer, Karlberg, 2018

$e^+e^- \rightarrow t\bar{t}$

Chen, Guan, He, Liu, Ma

...

...

Fixed order



Chen, HTL, Shao, Wang, PLB and JHEP, 2020

Developments in Effective theory

● **LO**

● **NLO**

arXiv: hep-ph/9805244

● **NNLO**

arXiv:1305.5206., arXiv:1309.6594,
arXiv: 1408.2422, arXiv:1606.09519

● **NNNLO**

Our works: arXiv:1909.06808, arXiv:1912.13001

● **NNLL+(N)NLO**

D.Y.Shao, C.S.Li, HTL, J.Wang, arXiv:1301.1245

arXiv:1505.07122, arXiv: 1807.03704

Developments in Full theory

● **LO**

PLB 197 (1987) 269,
NPB 309 (1988) 282,
arXiv:hep-ph/9603205

● **NLO**

arXiv: 1604.06447, arXiv:1608.04798,
arXiv:1811.05692, arXiv:1907.06408

● **PS+NLO**

arXiv:1703.09252, arXiv:1711.03319,
arXiv:1903.08137

● **NLL+NLO q_T** arXiv:1609.01691

●

$\sqrt{s}$	13 TeV	14 TeV	27 TeV	100 TeV
LO	$13.80^{+31\%}_{-22\%}$	$17.06^{+31\%}_{-22\%}$	$98.22^{+26\%}_{-19\%}$	$2015^{+19\%}_{-15\%}$
NLO	$25.81^{+18\%}_{-15\%}$	$31.89^{+18\%}_{-15\%}$	$183.0^{+16\%}_{-14\%}$	$3724^{+13\%}_{-11\%}$
NNLO	$30.41^{+5.3\%}_{-7.8\%}$	$37.55^{+5.2\%}_{-7.6\%}$	$214.2^{+4.8\%}_{-6.7\%}$	$4322^{+4.2\%}_{-5.3\%}$
N ³ LO	$31.31^{+0.66\%}_{-2.8\%}$	$38.65^{+0.65\%}_{-2.7\%}$	$220.2^{+0.53\%}_{-2.4\%}$	$4439^{+0.51\%}_{-1.8\%}$

❑ the QCD corrections are very prominent

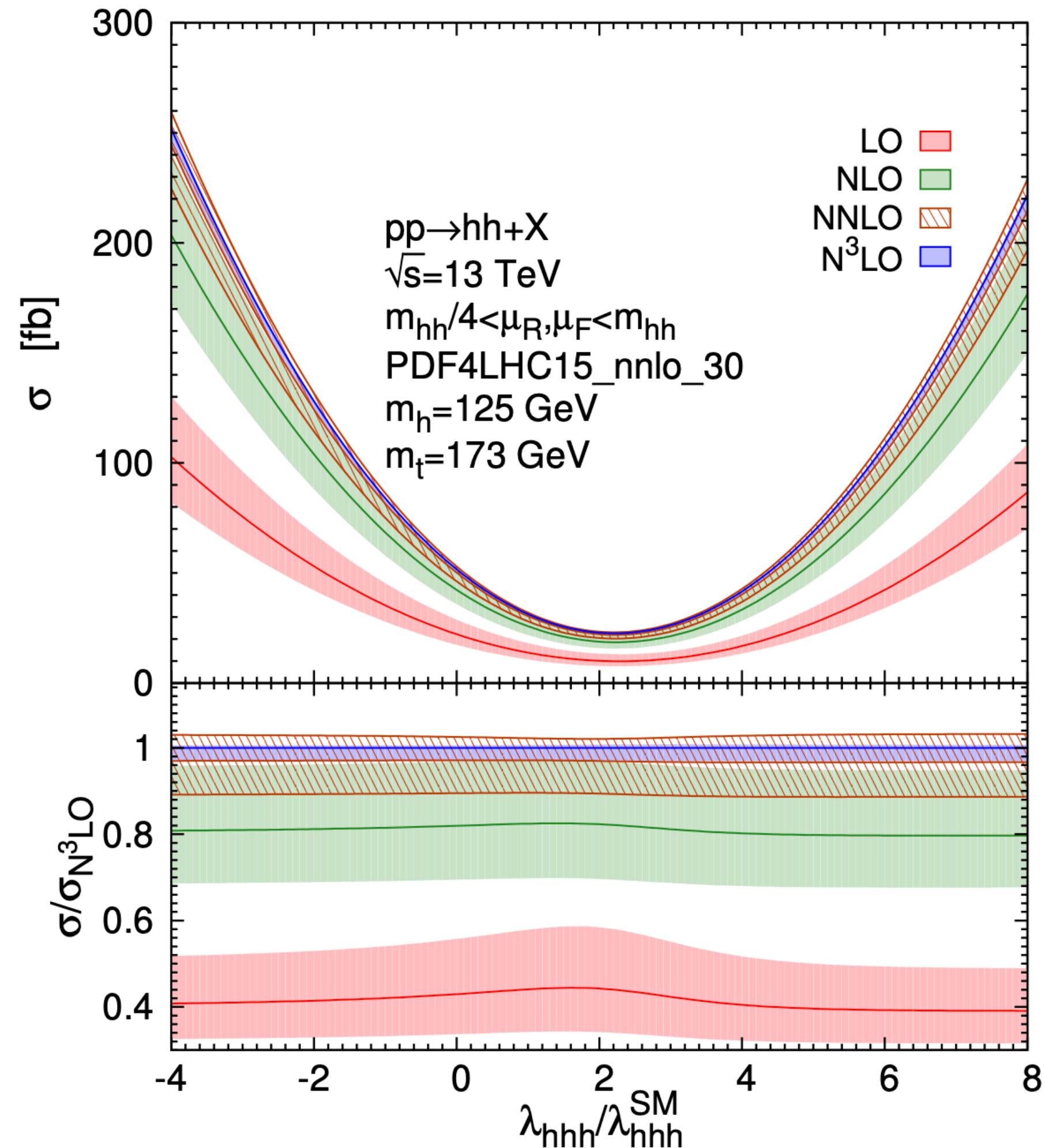
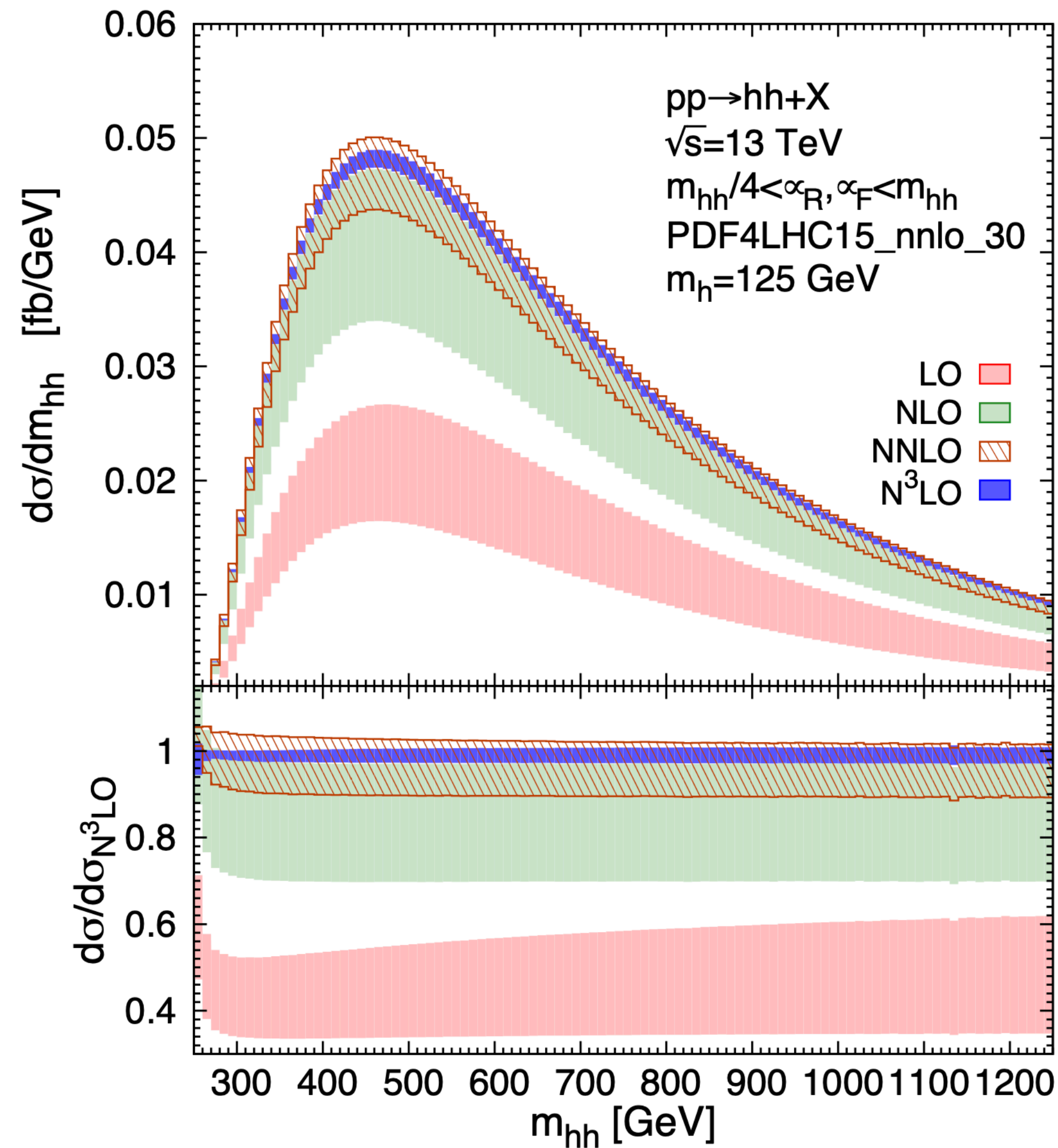
❑ NLO QCD corrections are around 87%

❑ NNLO QCD corrections are around 18%

❑ N³LO QCD corrections are around 3%

❑ N³LO scale uncertainties are less than 3%

Fixed order



Full theory

269,
282,
103205

Xiv:1608.04798,
iv:1907.06408

arXiv:1711.03319,

Xiv:1609.01691

dominant

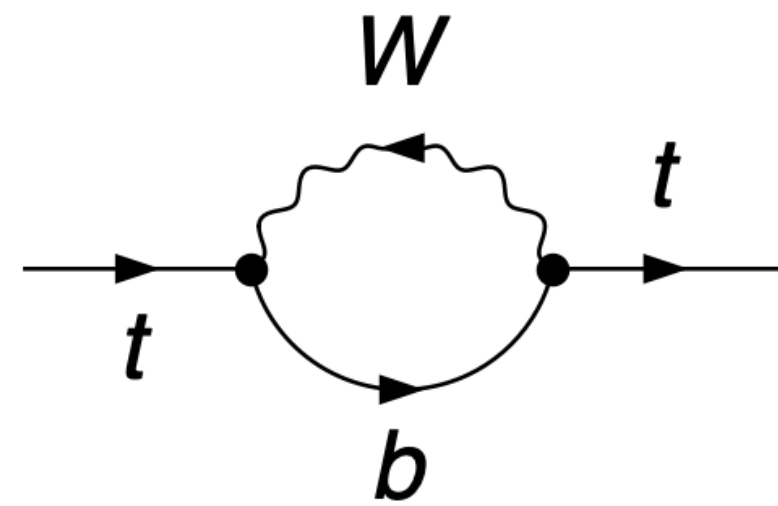
187%

and 18%

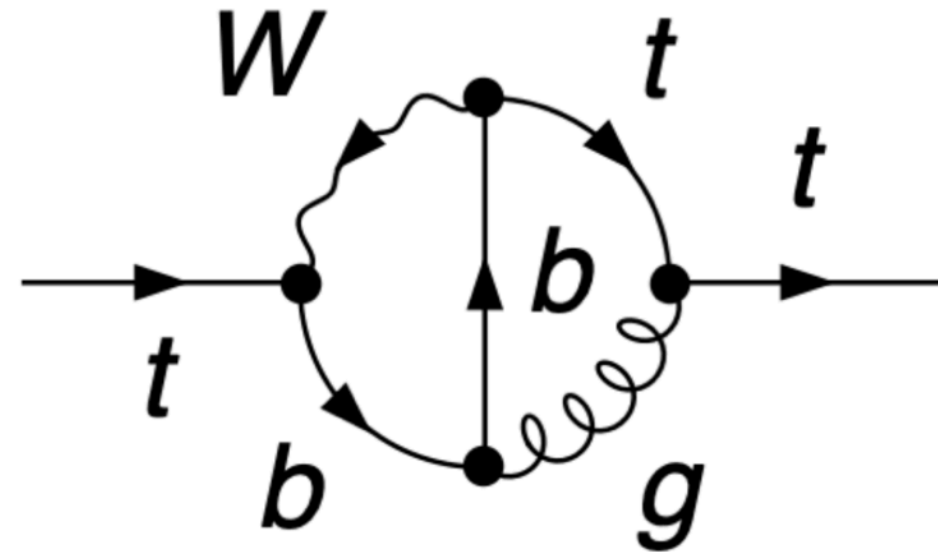
and 3%

is than 3%

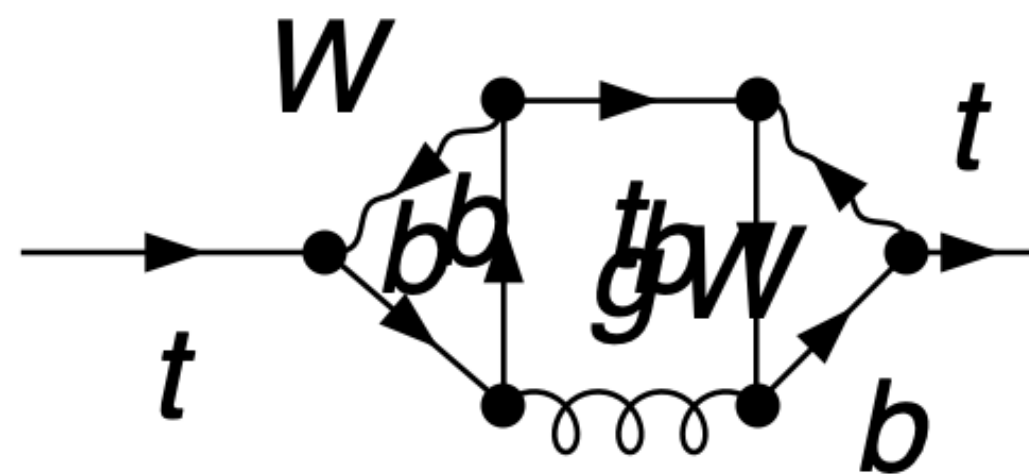
Fixed order



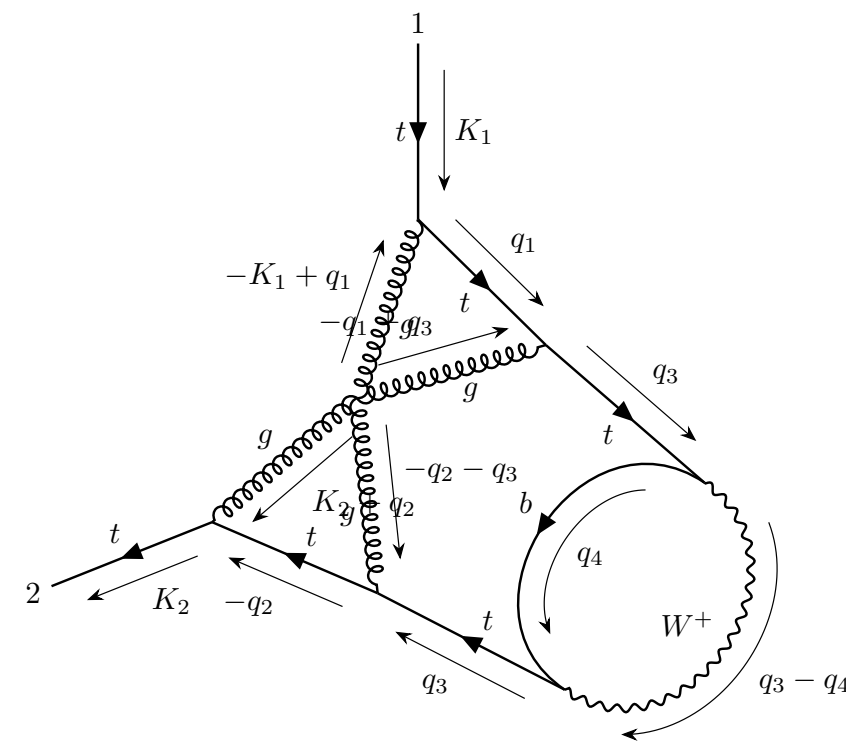
1-loop, LO



2-loop, NLO

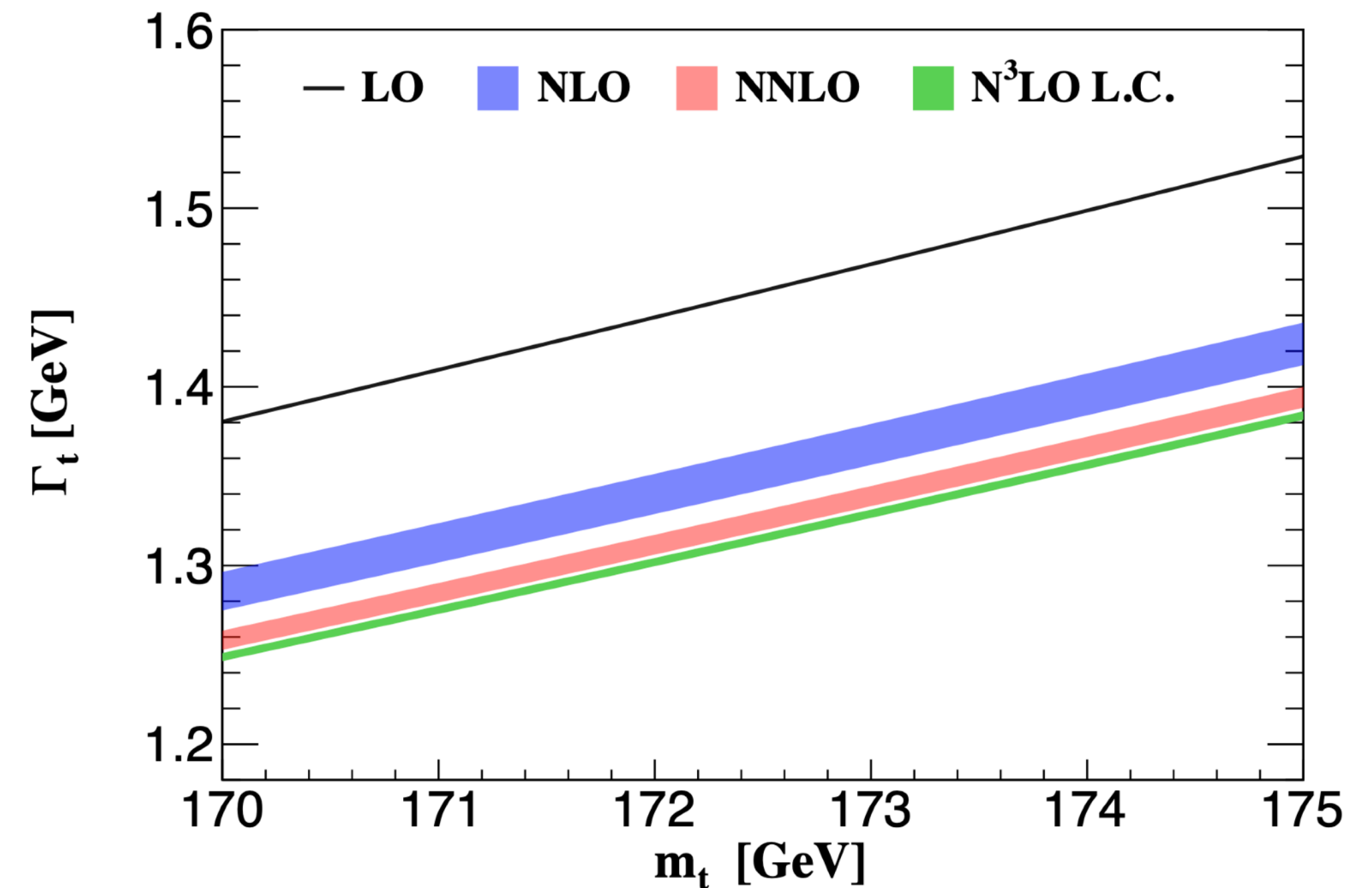


3-loop, NNLO



4-loop, NNNLO

$$\Gamma_t = \Gamma_0 \left[X_0 + \frac{\alpha_s}{\pi} X_1 + \left(\frac{\alpha_s}{\pi} \right)^2 X_2 + \left(\frac{\alpha_s}{\pi} \right)^3 X_3 \right]$$



NLO: Bigi, Dokshitzer, Khoze, Kuhn, Zerwas, 1986, Jezabek, Kuhn 1989, Czarnecki 1990, Li, Oakes, Yuan 1991

NNLO: Gao, Li, Zhu 2013, Brucherseifer, Caola, Melnikov 2013

Chen, HTL, Wang, Wang, PRD, arXiv:2212.06341
Chen, HTL, Li, Wang, Wang, Wu, arXiv 2309.00762

Chen, Chen, Guan, Ma, arXiv:2309.0193
see the talks by Long Chen and YeFan Wang

Resummation

At leading power, generally the cross section can be factorized into

$$d\sigma = H \otimes (B_1 \otimes B_2) \otimes (J_1 \otimes J_2) \otimes S$$

hard function
happens at high scale

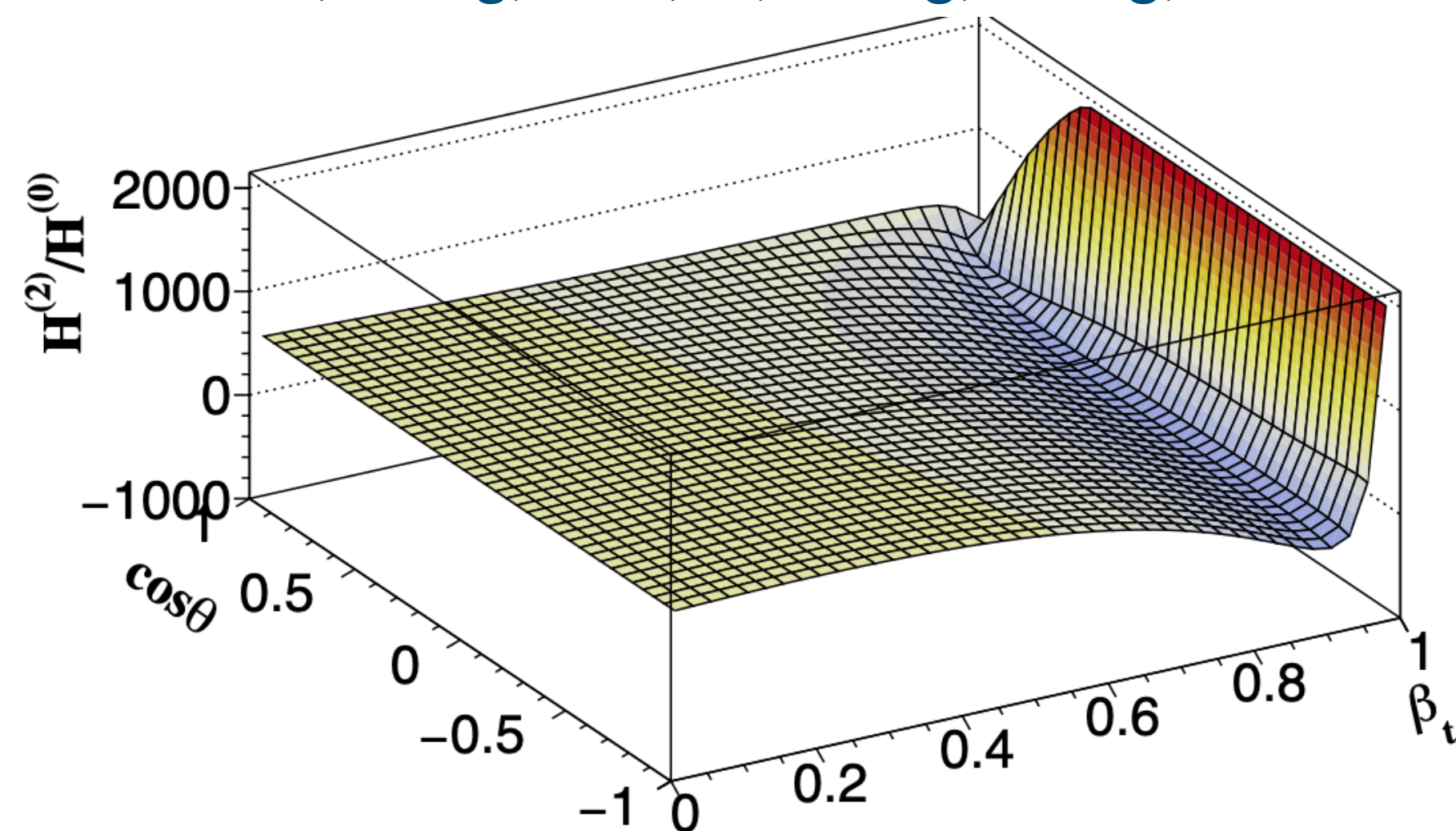
beam function
initial state collinear
radiation

Jet function
final state collinear
radiation

soft radiation

Hard function

2-loop hard function for tW,
Chen, Dong, HTL, Li, Wang, Wang, JHEP 2022



3-loop and 4-loop calculation about the form factors

3-loop anomalous dimensions

$$\begin{aligned} \Gamma(\{p\}, \{\underline{m}\}, \mu) = & \sum_{(i,j)} \frac{\mathbf{T}_i \cdot \mathbf{T}_j}{2} \gamma_{\text{cusp}}(\alpha_s) \ln \frac{\mu^2}{-s_{ij}} + \sum_{I,j} \mathbf{T}_I \cdot \mathbf{T}_j \gamma_{\text{cusp}}(\alpha_s) \ln \frac{m_I \mu}{-s_{Ij}} - \sum_{(I,J)} \frac{\mathbf{T}_I \cdot \mathbf{T}_J}{2} \gamma_{\text{cusp}}(\beta_{IJ}, \alpha_s) \\ & + \sum_i \gamma^i(\alpha_s) \mathbf{1} + \sum_I \gamma^I(\alpha_s) \mathbf{1} + f(\alpha_s) \sum_{(i,j,k)} \mathcal{T}_{ijk} + \sum_{(i,j,k,l)} \mathcal{T}_{ijkl} F_4(\beta_{ijkl}, \beta_{ijkl} - 2\beta_{ilkj}, \alpha_s) \\ & + \sum_I \sum_{(i,j)} \mathcal{T}_{ijII} F_{h2}(r_{ijI}, \alpha_s) + \sum_I \sum_{(i,j,k)} \mathcal{T}_{ijkI} F_{h3}(r_{ijI}, r_{ikI}, r_{jkI}, \alpha_s) \\ & + [\text{non-dipole contributions involving two or more massive partons}] + \mathcal{O}(\alpha_s^4) \end{aligned}$$

One Massive Parton, Liu, Schalch, PRL, 2022

Resummation

At leading power, generally the cross section can be factorized into

$$d\sigma = \boxed{H} \otimes \boxed{B_1 \otimes B_2} \otimes \boxed{J_1 \otimes J_2} \otimes \boxed{S}$$

hard function

happens at high scale

beam function

initial state collinear radiation

Jet function

final state collinear radiation

soft radiation

Beam, soft and jet functions are process independent

3-loop TMD beam function (TMDPDFs) Luo, Yang et al PRL, 2020
Elbert, Mistlberger, Vita, JHEP, 2019

3-loop n-jettiness beam function
Elbert, Mistlberger, Vita, JHEP, 2020
Baranowski et al, JHEP, 2023

4-loop Rapidity Anomalous Dimension, Moult, Zhu, Zhu, JHEP, 2020, Duhr, Mistlberger, Vita, JHEP, 2020

2-loop soft function for Jet veto, Abreu, Gaunt, Monni, Szafron, JHEP, 2022

2-loop beam for Jet veto, Bell, Brune, Das, Wald, JHEP, 2023, Abreu, Gaunt, Monni, Szafron, JHEP 2023

Resummation

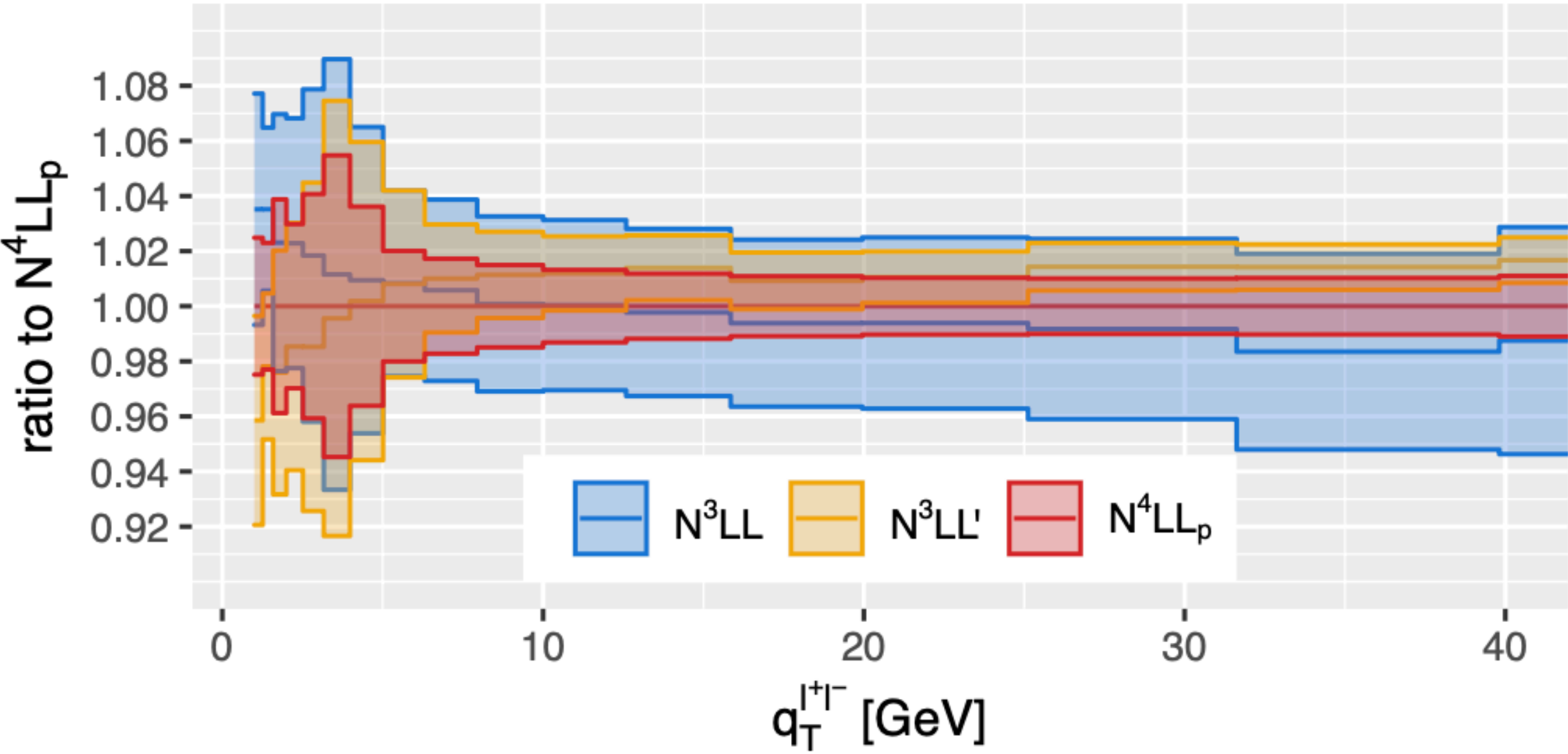
$$d\sigma = H \otimes B_1 \otimes B_2 \otimes J_1 \otimes J_2 \otimes S$$

Large logs can be resummed by solving the RG equations of different functions

Drell-Yan q_T resummation at N4LL or N3LL

Camarda, Cieri, Ferrera, 2023
Neumann, Campbell, PRD 2023,

Order	Boundary conditions	Anomalous dimensions		FO matching (nonsingular)
		γ_i (noncusp)	$\Gamma_{\text{cusp}}, \beta$	
LL	1	-	1-loop	-
NLL	1	1-loop	2-loop	-
NLL' (+NLO)	α_s	1-loop	2-loop	α_s
NNLL (+NLO)	α_s	2-loop	3-loop	α_s
NNLL' (+NNLO)	α_s^2	2-loop	3-loop	α_s^2
N ³ LL (+NNLO)	α_s^2	3-loop	4-loop	α_s^2
N ³ LL' (+N ³ LO)	α_s^3	3-loop	4-loop	α_s^3
N ⁴ LL (+N ³ LO)	α_s^3	4-loop	5-loop	α_s^3

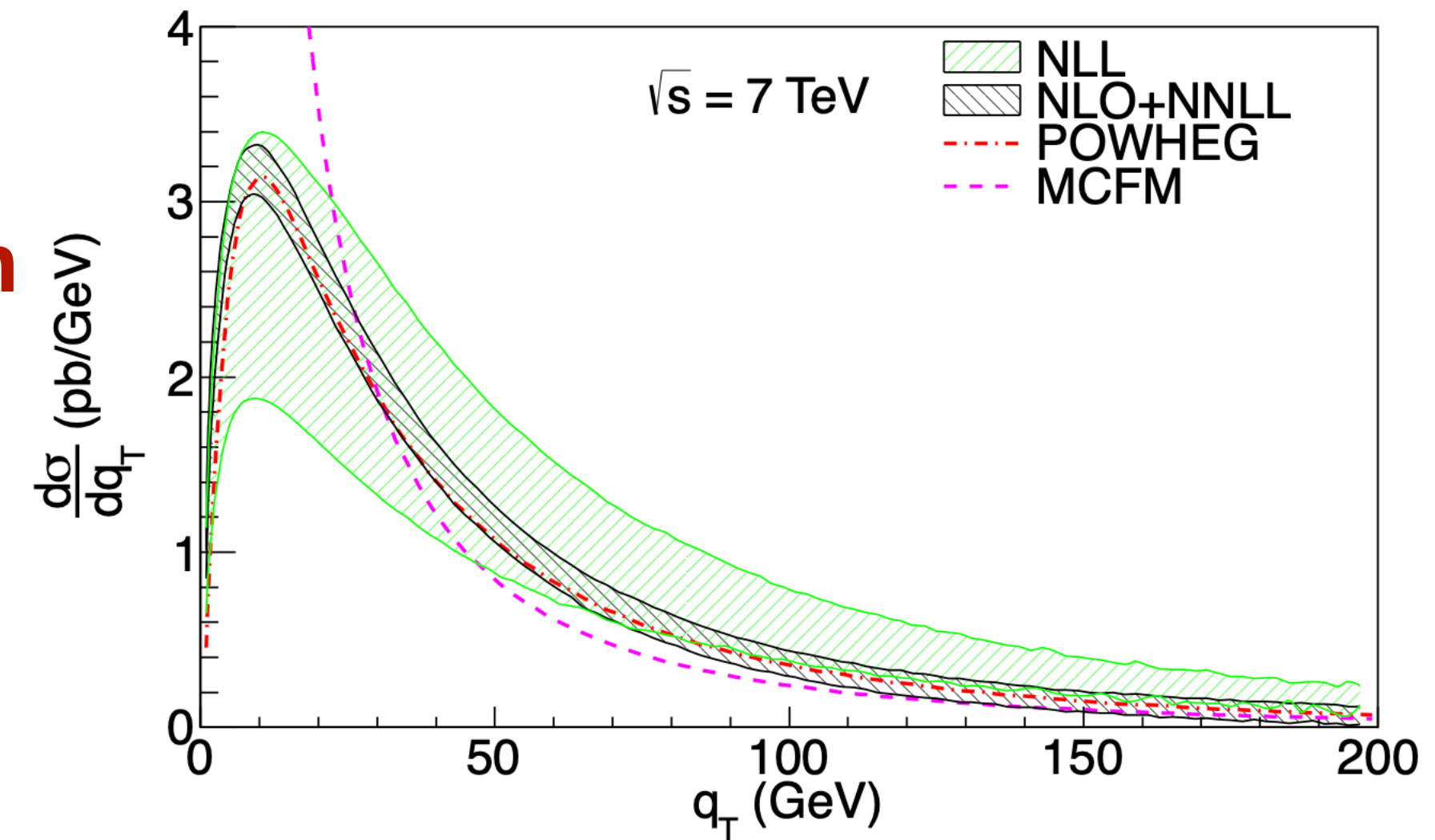


See Y. Fu’s talk about this topic

Resummation

TMD resummation for top quark pair production

Zhu, Li, HTL, Shao, Yang, PRL and PRD, 2013



TEEC for dijet production at hadron collide.

$$\text{TEEC} = \sum_{a,b} \int d\sigma_{pp \rightarrow a+b+X} \frac{2E_{T,a}E_{T,b}}{|\sum_i E_{T,i}|^2} \delta(\cos \phi_{ab} - \cos \phi)$$

Gao, HTL, Moult, Zhu, PRL, 2019

For TEEC/EEC observable to DIS process

HTL, Vitev, Zhu, JHEP, 2020

$$\text{TEEC} = \sum_a \int d\sigma_{lp \rightarrow l+a+X} \frac{E_{T,l}E_{T,a}}{E_{T,l} \sum_i E_{T,i}} \delta(\cos \phi_{la} - \cos \phi)$$

HTL, Makris, Vitev, PRD, 2021

$$\text{EEC} = \sum_a \int d\sigma_{lp \rightarrow l+a+X} \left(\frac{p \cdot p_a}{\sum_i p \cdot p_i} \right) \delta(\cos \chi - \cos \theta_{ap})$$

Resummation with Parton Shower

- Parton shower accuracy is critical at the LHC
- the interpretation of data
 - data analysis with machine-learning

Showers resum NLL resummation or beyond, and spin correlations

Many PanScales papers, [arXiv:2305.08645](#), PanScales for DIS

Ravasio, Hamilton et al PRL, 2023, double soft emission can be included by veto algorithms

$$P_{\text{accept}} = \frac{|M_{\text{DS}}|^2}{\sum_h |M_{\text{shower},h}|^2}$$

Nagy, Soper, PRD, 2021, Dedector, how parton shower sums large logarithms

Herren, Hoche et al, JHEP, 2023, Alaric, new approach for parton shower

Assi, Hoche, 2023, Alaric for massive partons

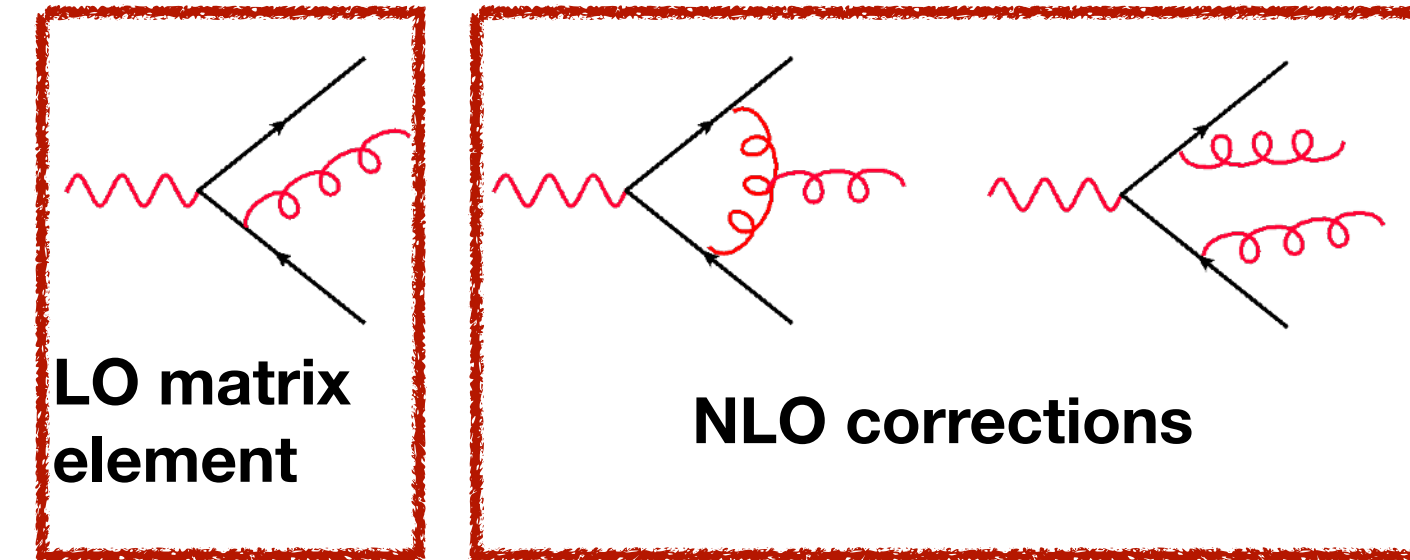
Matching with fixed order calculations combine resummation with hard emission

^{Mi}Method developed: UNLOPS, NNLOPS, Geneva, MiNNLOPS

NNLO+PS done for most of 2to2 process, but no light jet final state yet

Resummation with Parton Shower

In order to fully cover the phase space and reproduce the singularity at NNLO level, we need to have $2 \rightarrow 3$ and $2 \rightarrow 4$ showers



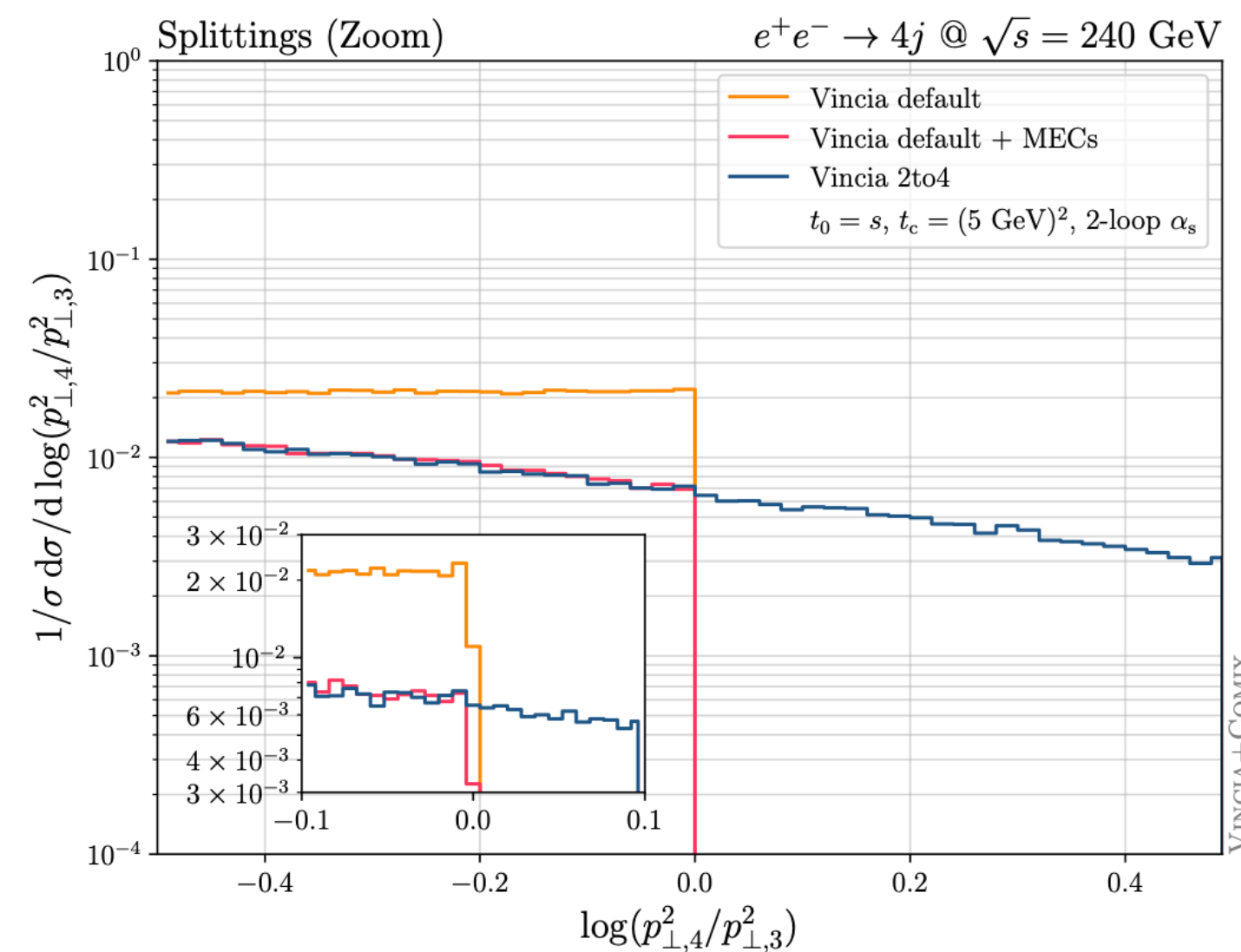
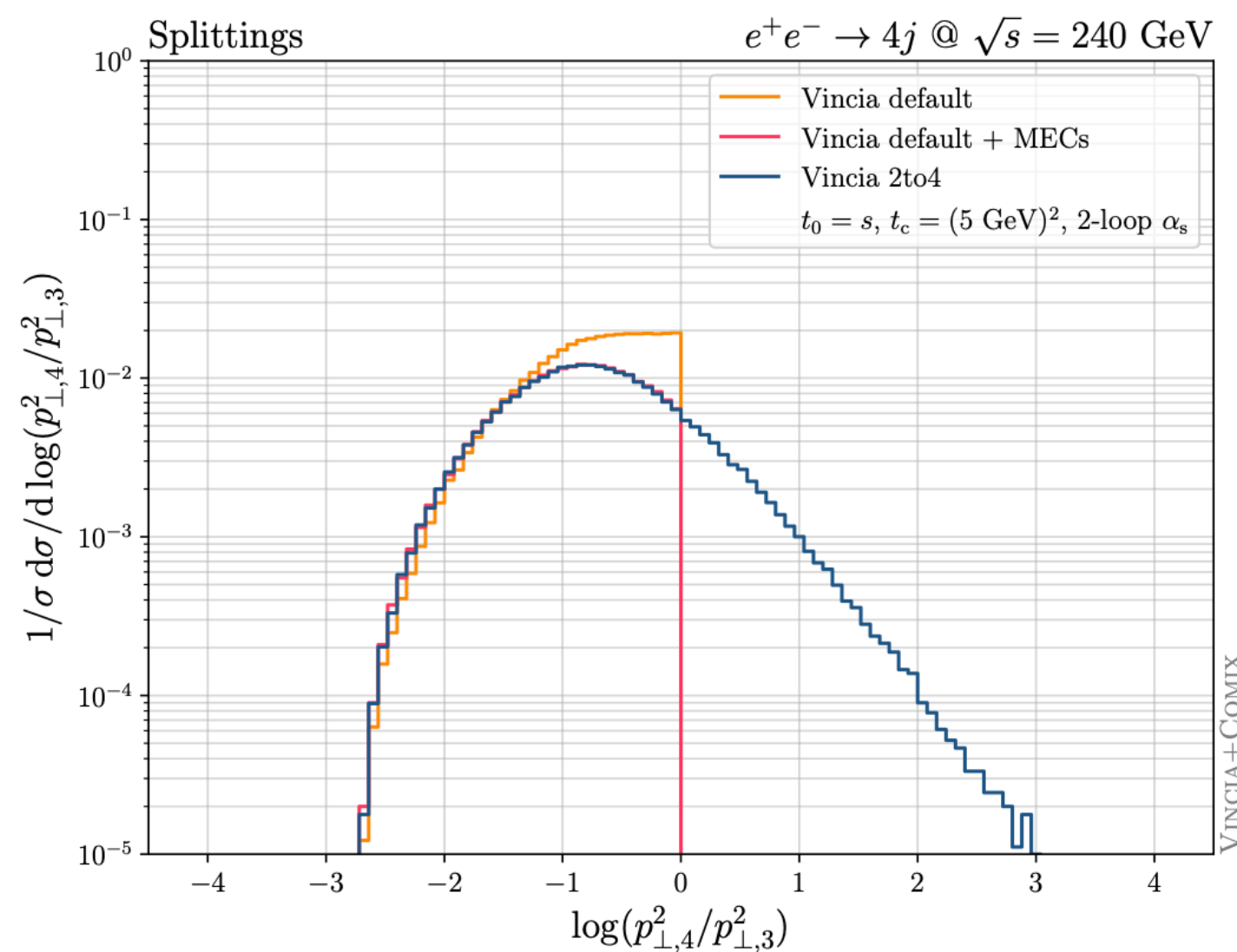
3 parton final state

3+4 parton final state

HTL, Skands, PLB, 2017

match final-state parton showers fully-differentially to NNLO calculations

Compbel, Hoech, HTL, Preuss, Skands, PLB, 2023



Jet

Jet Physics and it's applications

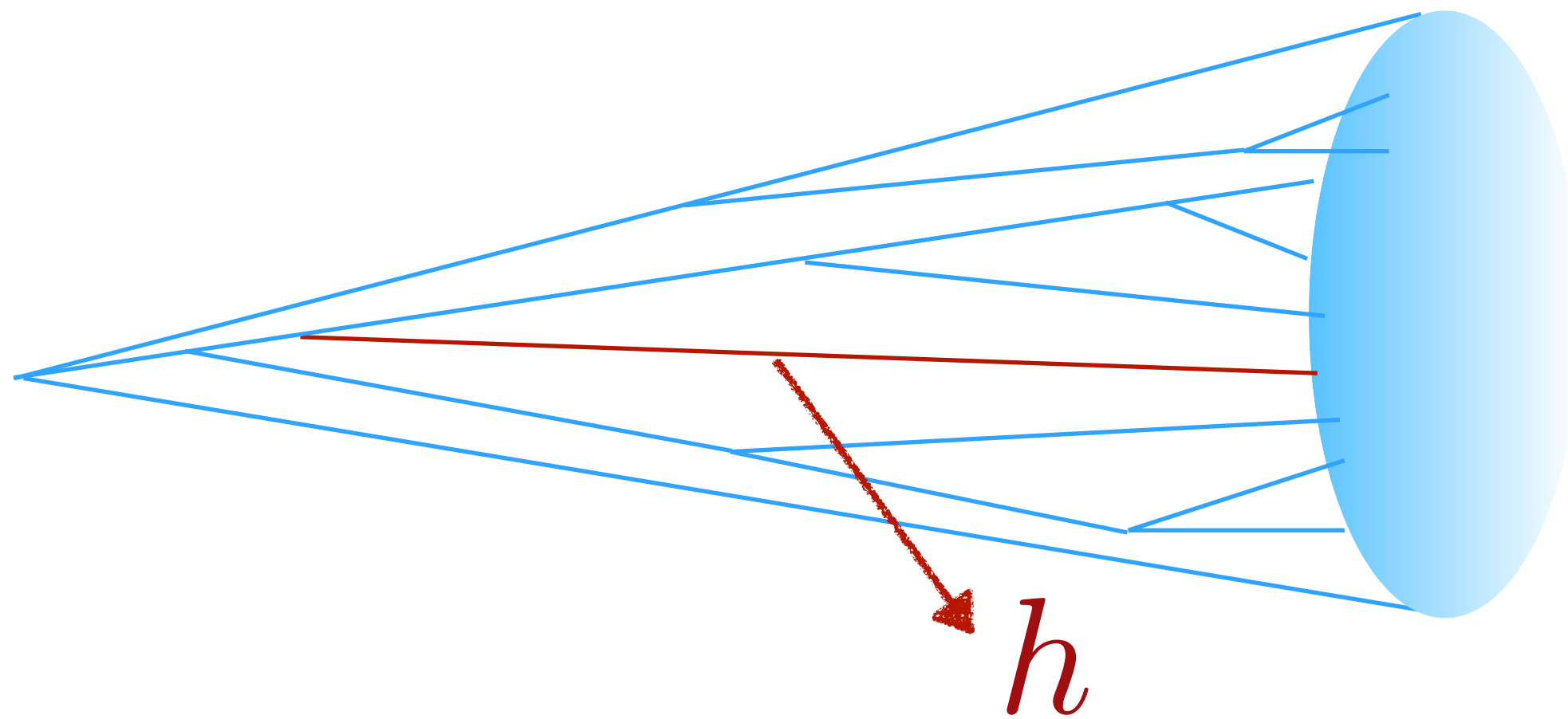
Object to study perturbative and nonperturbative QCD

Production of jets: automatical NLO generator, [arXiv:2305.14620](https://arxiv.org/abs/2305.14620)

talk by C.Y. Liu

Resummation of Non-Global logs

talk by D.Y. Shao



Jet Property: groomed angle or momentum fraction, mass, angularity, jet fragmentation function, et al

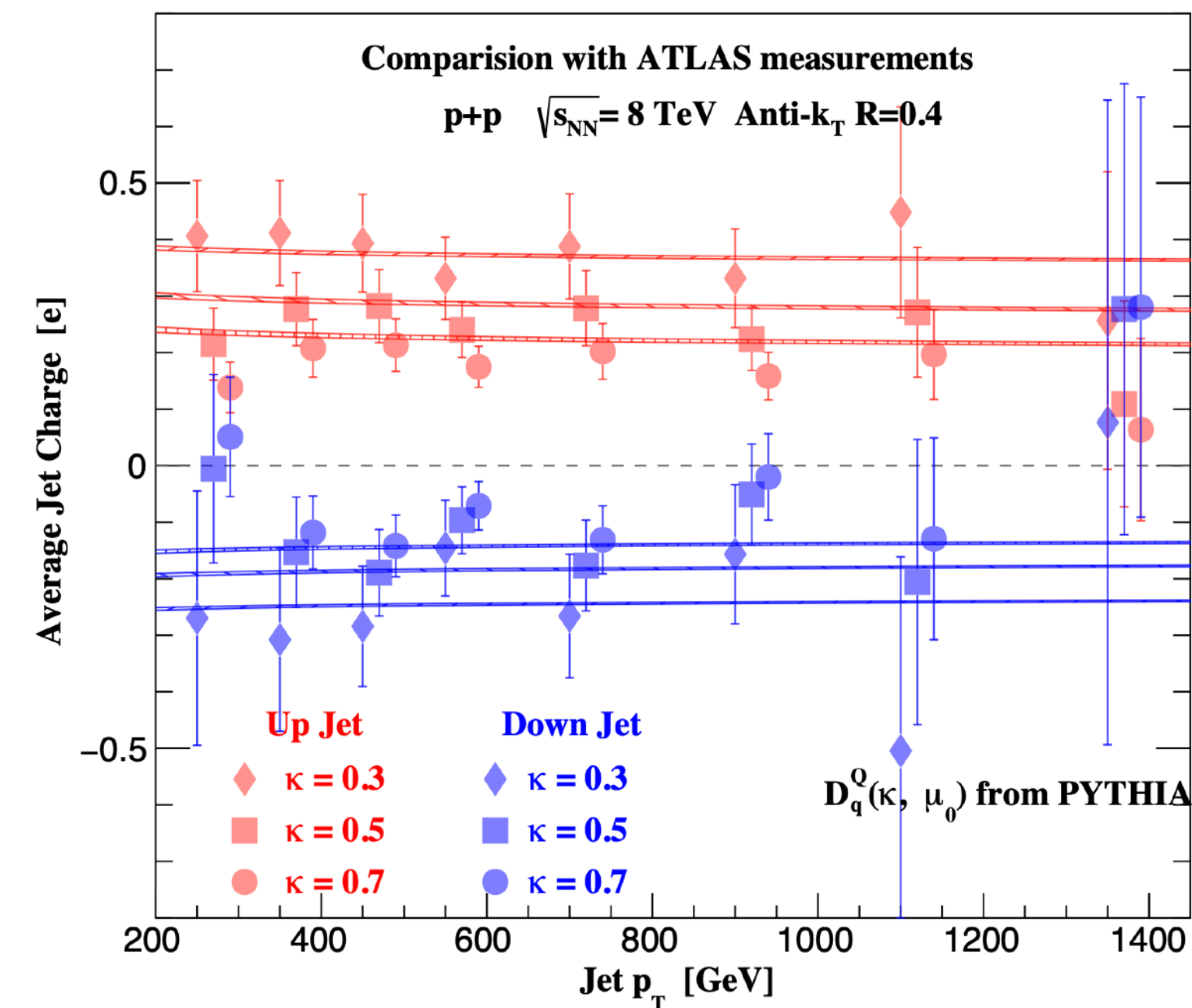
about jet charge and its applications

HTL, Vitev, PRD, 2020

HTL, Vitev, PRL, 2021

HTL, Yan, Yuan, PLB, 2022

HTL, Yan, Yuan, PRL, 2023



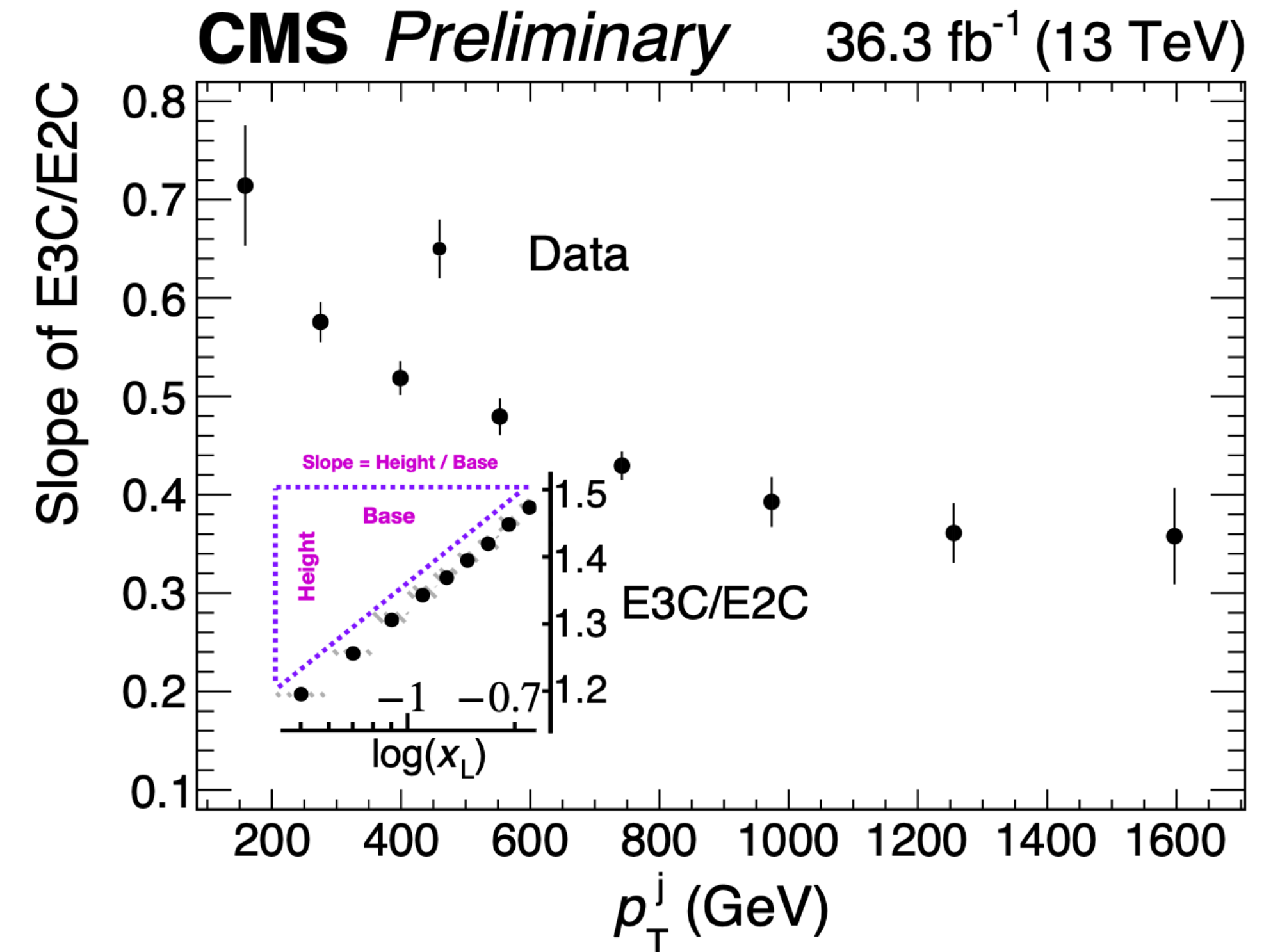
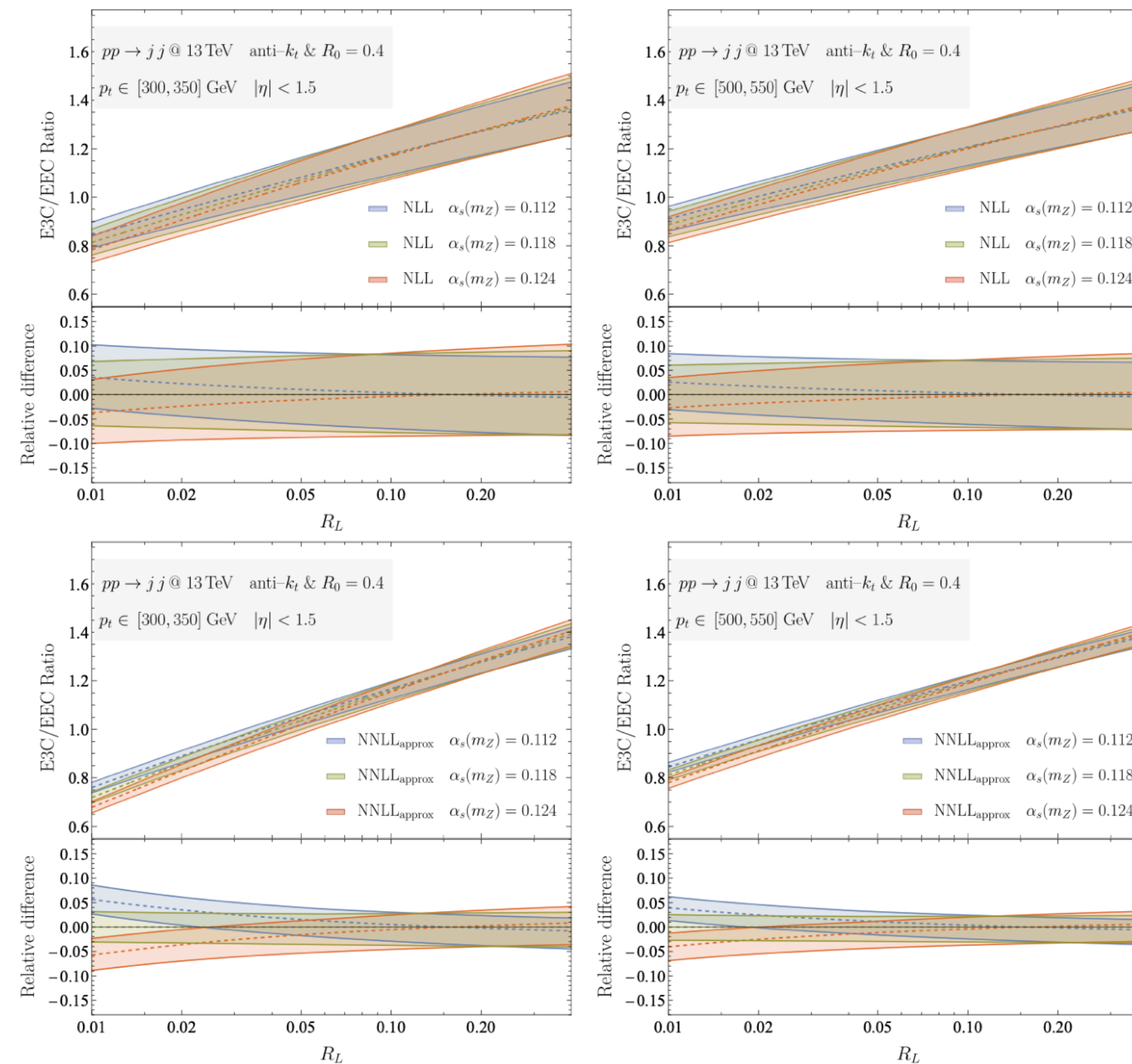
Jet

Jet Physics and it's applications Energy Correlators

$$\frac{d\sigma^{[N]}}{dx_L} \equiv \sum_n \sum_{1 \leq i_1, \dots, i_N \leq n} \int d\sigma \frac{\prod_{a=1}^N E_{i_a}}{Q^N} \delta \left(x_L - \max \left\{ x_{i_1, i_2}, x_{i_1, i_3}, \dots, x_{i_{N-1}, i_N} \right\} \right)$$

[arXiv:2307.07510](#)
[arXiv:2307.15739](#)
[arXiv:2303.01530](#)
[arXiv:2301.01788](#)
[arXiv:2210.10061](#)
[arXiv:2209.02080](#)

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Status: February 2022



Summary

Compared to High Luminosity LHC (goal 4000 fb^{-1}), roughly 5% of data collected

Small deviations, if there are, will indicate possible signal for new physics

Precision predictions or tools are essential

**Many other interesting results not covered by this talk.
Apology for it.**

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Thank You !