



北京航空航天大学
BEIHANG UNIVERSITY

Freeze-in of WIMP dark matter

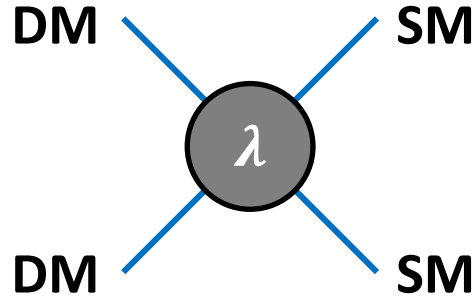
Ke-Pan Xie (谢柯盼)

Beihang University

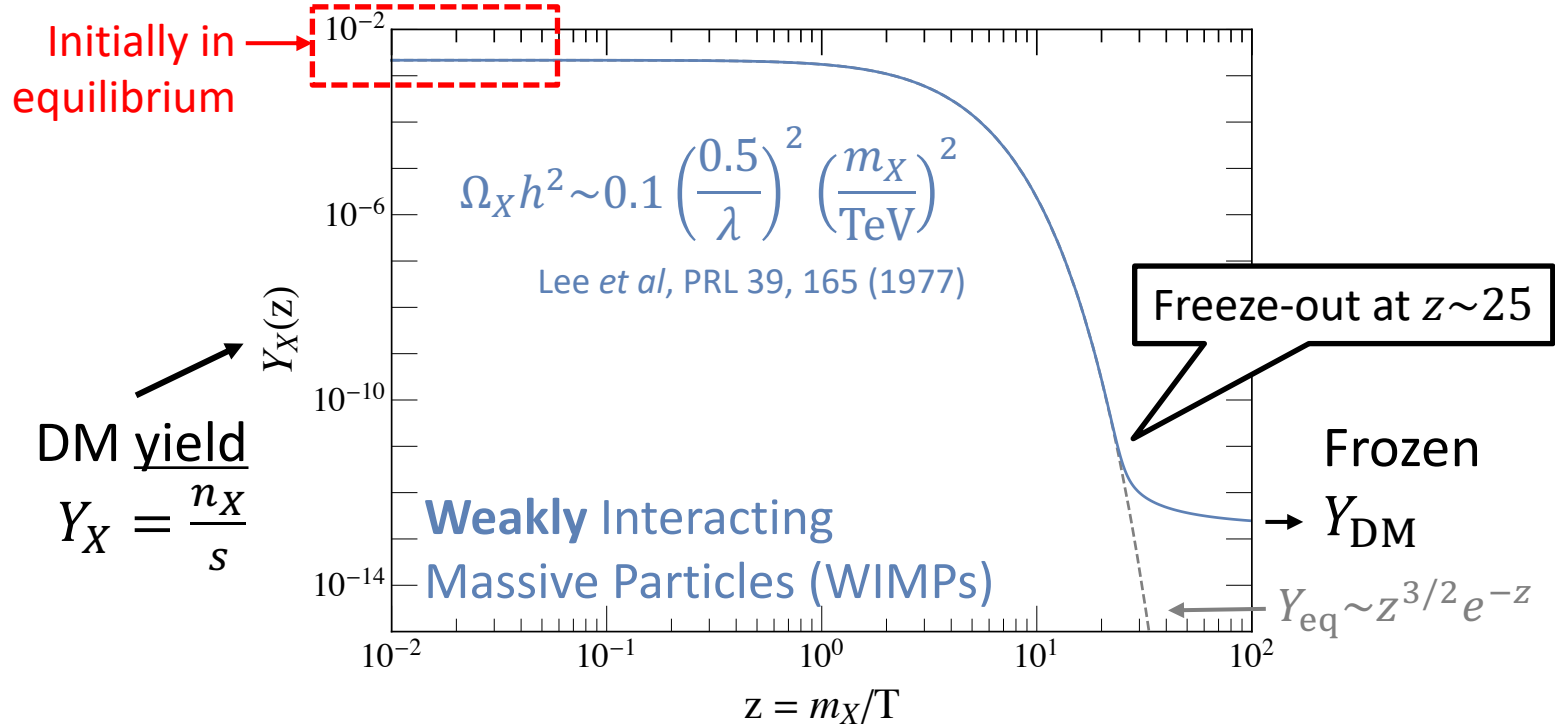
2023.6.3 @TOPAC2023, TDLI

With Xiaorui Wong, 2304.00908

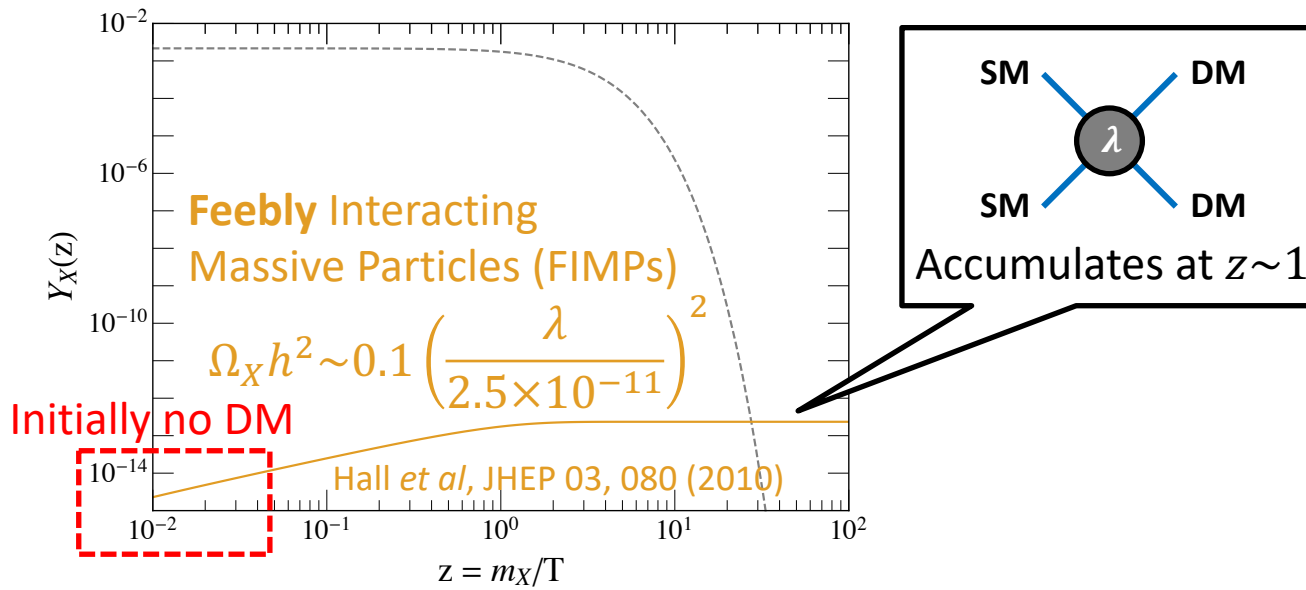
The most popular explanation for particle DM



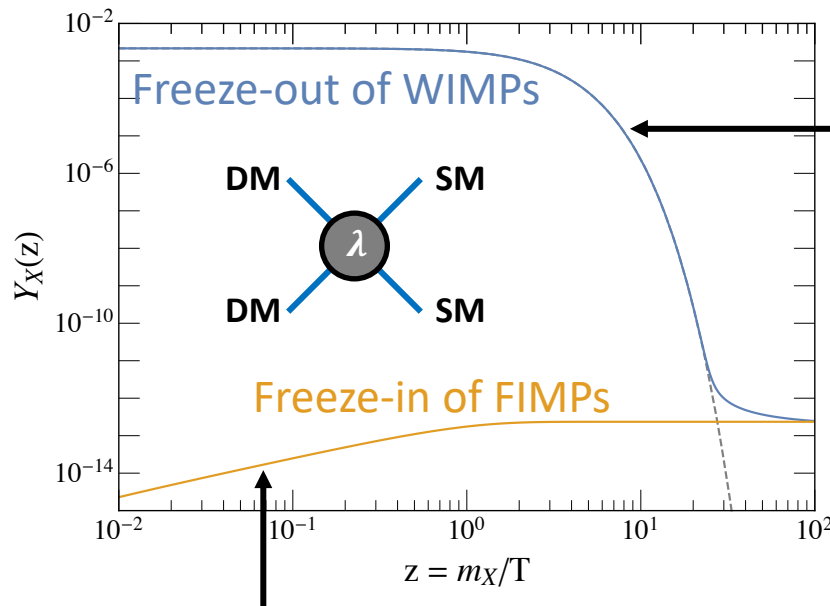
$$\langle \sigma v_{\text{rel}} \rangle \sim \frac{\lambda^2}{32\pi m_X^2}$$



WIMP freeze-out and FIMP freeze-in



WIMP freeze-out and FIMP freeze-in

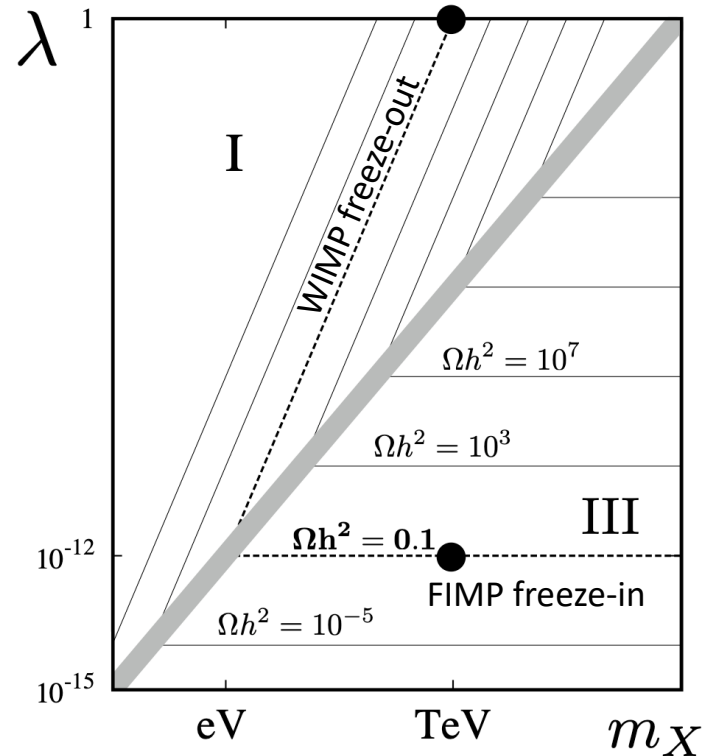


WIMPs: initially in equilibrium

DM scenario $\left(\frac{0.5}{\lambda}\right)^2 \left(\frac{m_X}{\text{TeV}}\right)^2 \sim 1$

FIMPs: initially no abundance

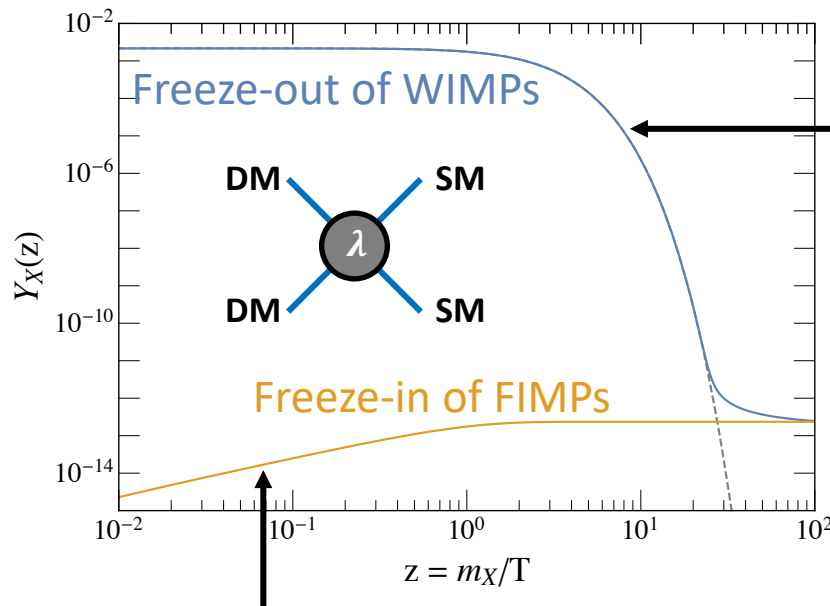
DM scenario $\left(\frac{\lambda}{2.5 \times 10^{-11}}\right)^2 \sim 1$



“Phase diagram”

Hall et al, JHEP 03, 080 (2010)

WIMP freeze-out and FIMP freeze-in



WIMPs: initially in equilibrium

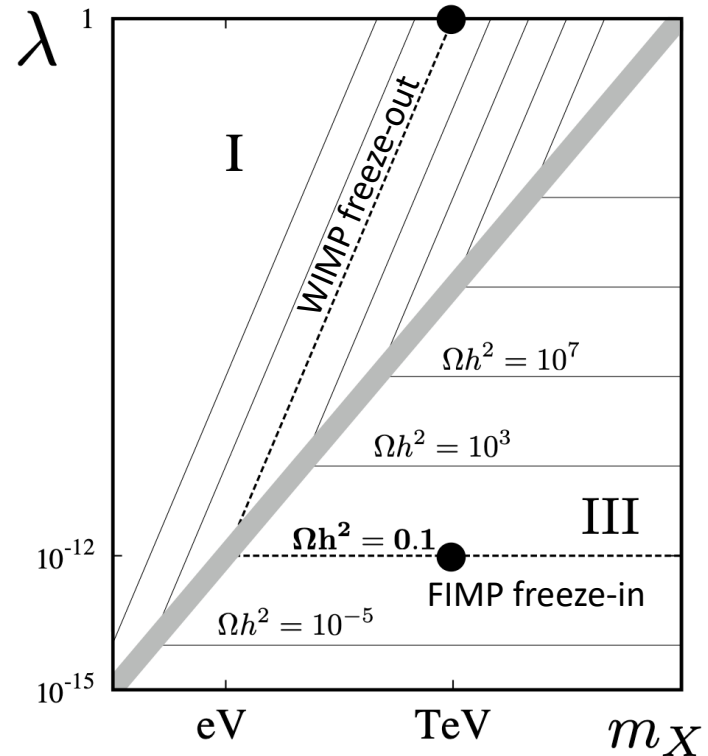
DM scenario $\left(\frac{0.5}{\lambda}\right)^2 \left(\frac{m_X}{\text{TeV}}\right)^2 \sim 1$

FIMPs: initially no abundance

DM scenario $\left(\frac{\lambda}{2.5 \times 10^{-11}}\right)^2 \sim 1$

This talk: WIMP can also freeze-in!

- The third scenario based on the $2 \rightarrow 2$ annihilation



“Phase diagram”

Hall et al, JHEP 03, 080 (2010)

How to realize freeze-in for WIMPs?

Will be **in equilibrium (very abundant)** since λ is NOT feeble



Incompatible!



Needs **out-of-equilibrium** dynamics,
especially **no DM** initially

How to realize freeze-in for WIMPs?

Will be **in equilibrium (very abundant)** since λ is NOT feeble



WIMP can **freeze-in**
if the early Universe
is boiling



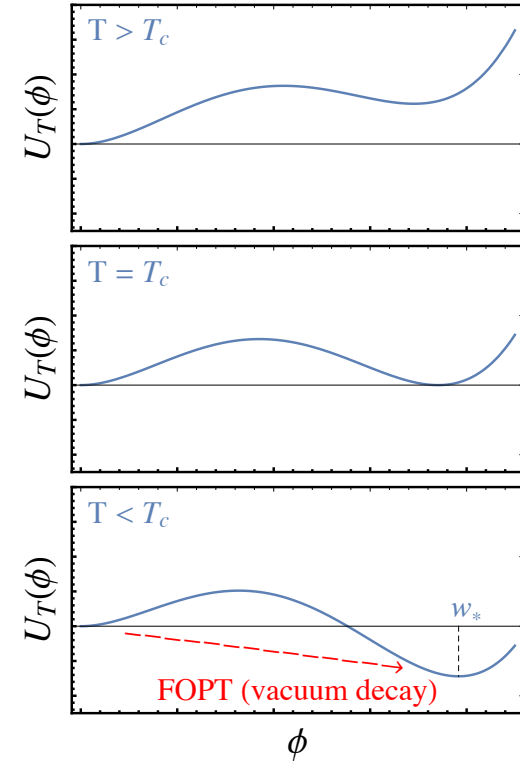
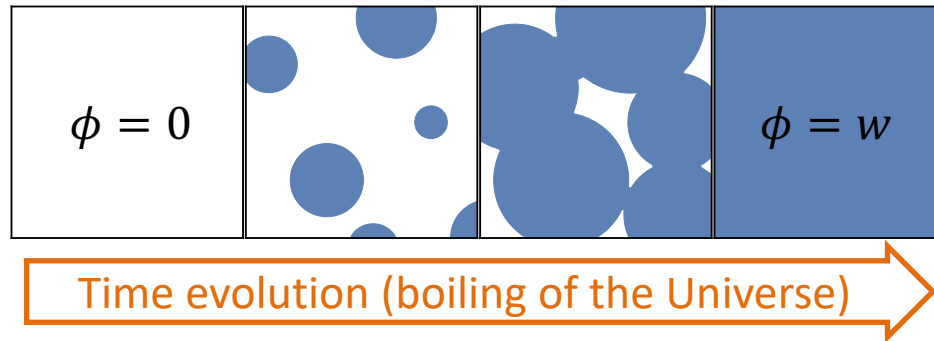
Needs **out-of-equilibrium** dynamics,
especially **no DM** initially

First-order phase transitions

FOPT: decay of the vacuum

$$\mathcal{L} \supset \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - U(\phi)$$

Early Universe $\Rightarrow U_T(\phi, T)$

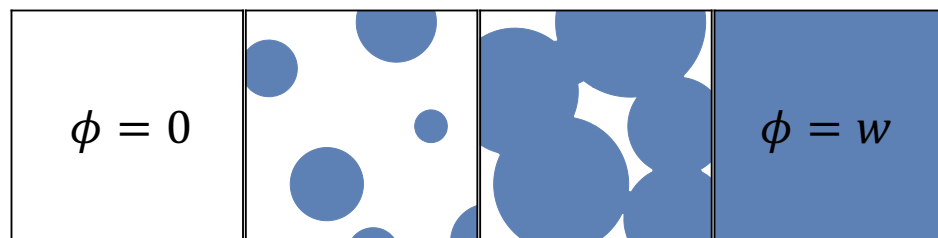


First-order phase transitions

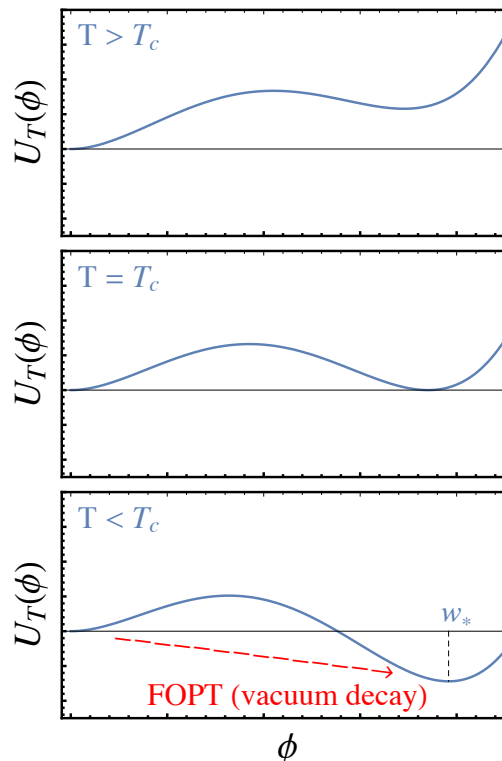
FOPT: decay of the vacuum

$$\mathcal{L} \supset \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - U(\phi)$$

Early Universe $\Rightarrow U_T(\phi, T)$



Time evolution (boiling of the Universe)



Vacuum energy ΔU released

- Cosmic temperature $T_1 \rightarrow T_2$
- Entropy injection $s_1 \rightarrow s_1 \left(\frac{T_2}{T_1}\right)^3$

Dilution of DM $Y_X \rightarrow Y_X \left(\frac{T_1}{T_2}\right)^3$

Assume $\mathcal{L} \supset -\frac{\lambda_{\phi X}}{2} \phi^2 |X|^2$

- Before transition $m_X = 0$
- After transition $m_X = \sqrt{\frac{\lambda_{\phi X}}{2}} w$

Jump of $z = \frac{m_X}{T} \Rightarrow 0 \rightarrow \mathcal{O}\left(\frac{w}{T_2}\right)$

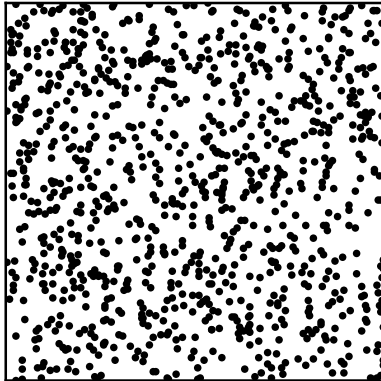
Realizing freeze-in with WIMPs in a strong FOPT

Freeze-in: needs zero DM abundance as initial condition

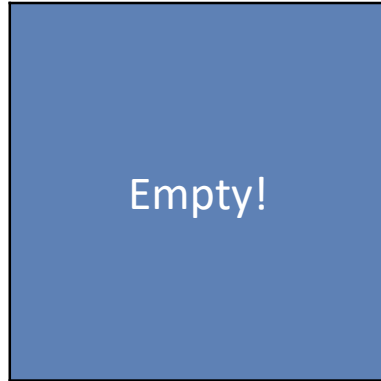
✗ Conventional WIMP: initially in equilibrium, **very abundant**

✓ Our scenario: preexisting DM is **diluted** to **zero**

Before FOPT



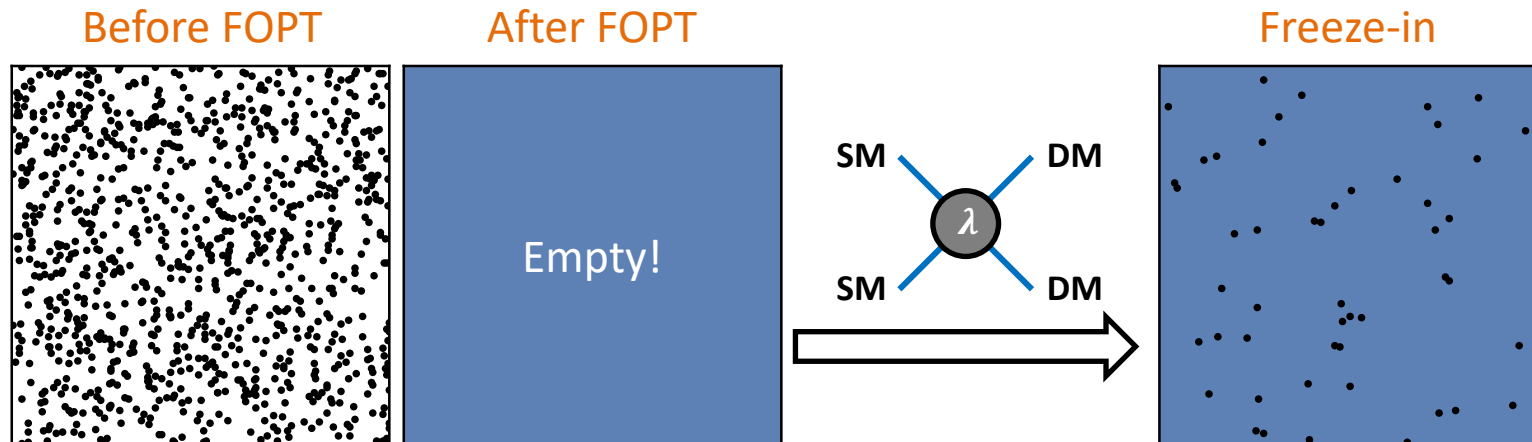
After FOPT



Realizing freeze-in with WIMPs in a strong FOPT

Freeze-in: needs zero DM abundance as initial condition

- ✗ Conventional WIMP: initially in equilibrium, **very abundant**
- ✓ Our scenario: preexisting DM is **diluted** to **zero**

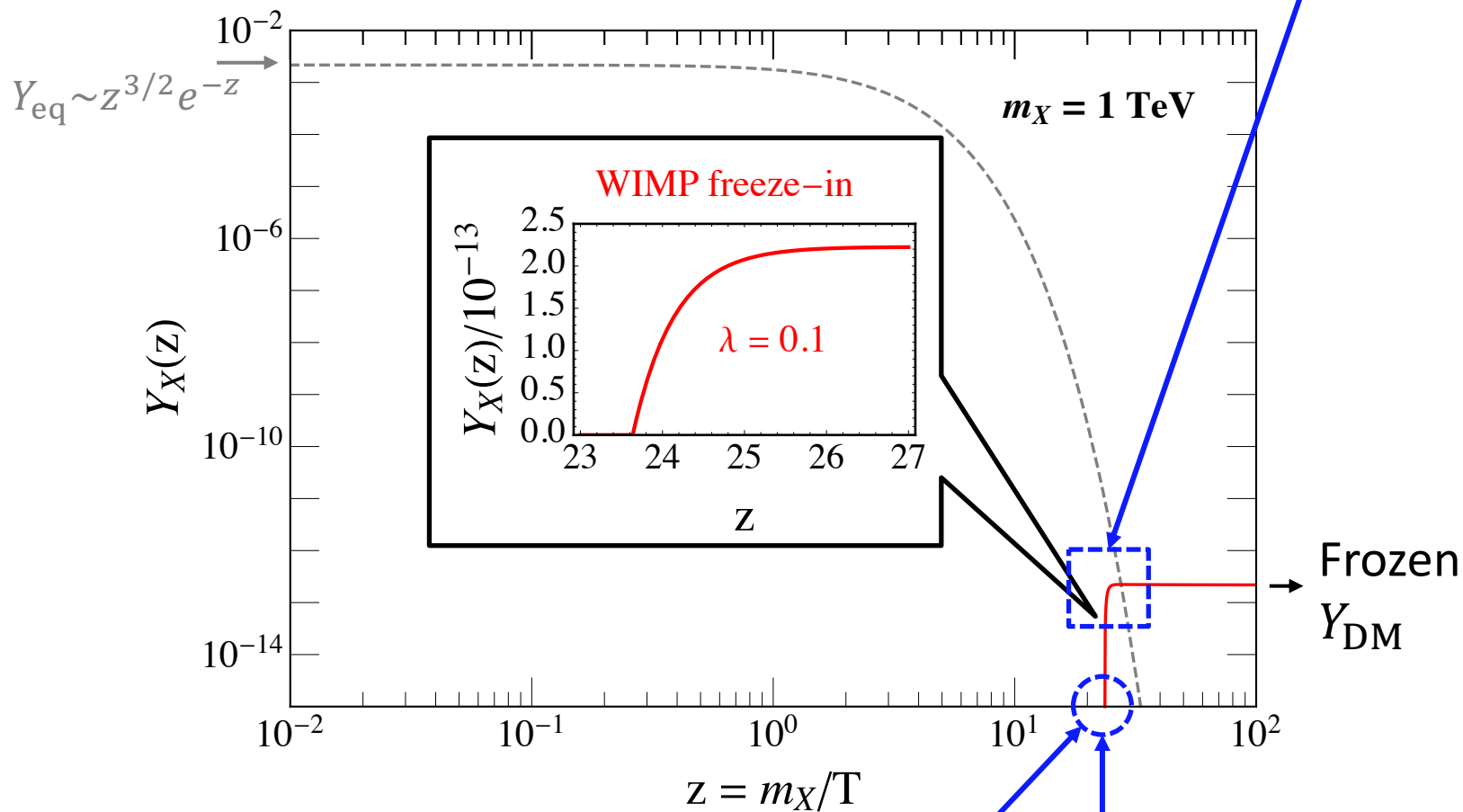


Freeze-in: during evolution, DM is never in equilibrium

- ✗ Conventional WIMP: be back to equilibrium **rapidly**, since λ is NOT feeble,
- ✓ Our scenario: cannot get back to equilibrium because of the **Boltzmann suppression** $e^{-\frac{m_X}{T_2}} \equiv e^{-z_2}$ (if $z_2 \sim \frac{w}{T_2} \gg 1$)

Freeze-in of WIMPs in a boiling Universe

2. WIMP cannot go back to equilibrium due to $e^{-z_2} \ll 1$

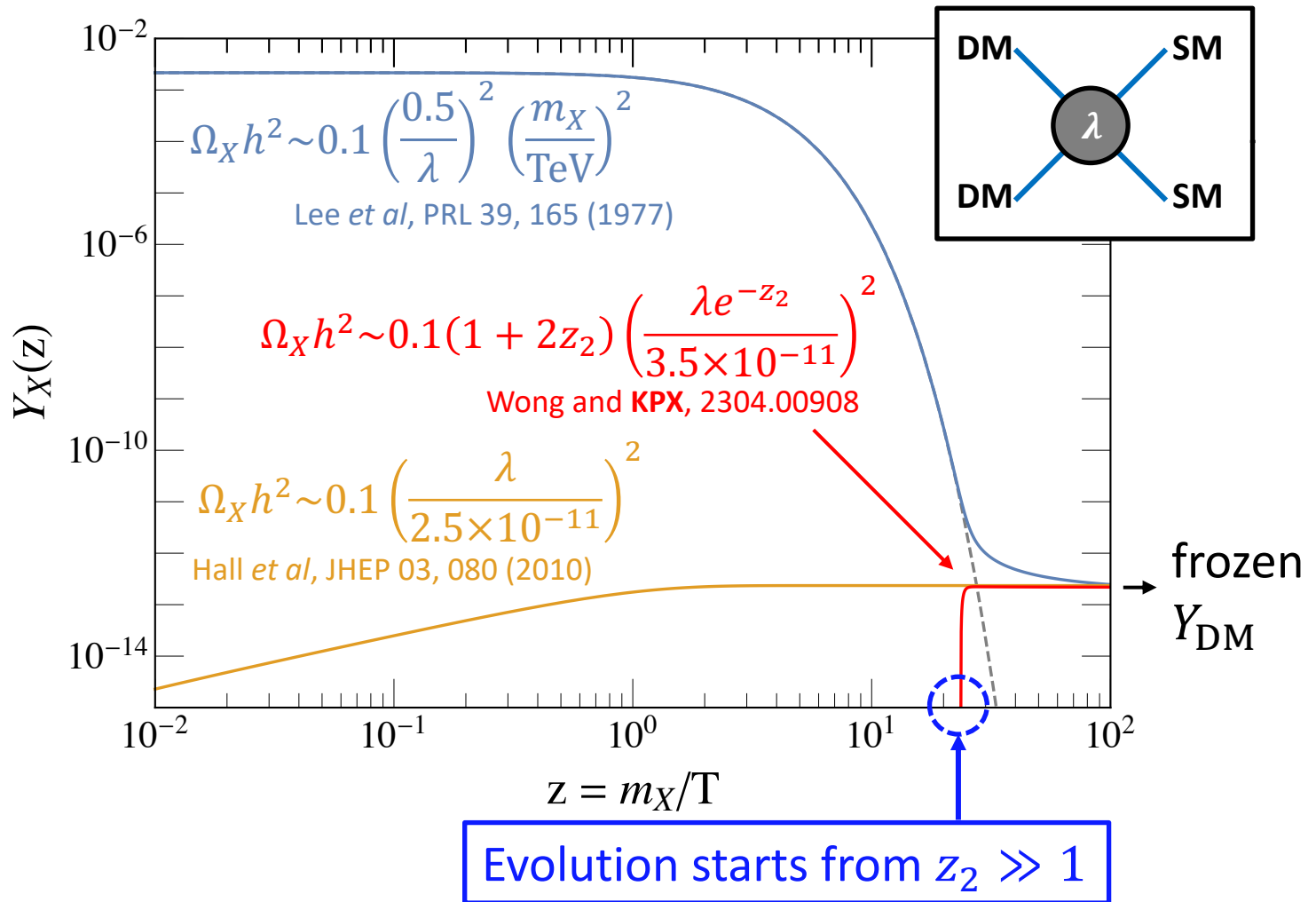


1. WIMP is diluted to zero after the FOPT

z_2 (set by FOPT)

Comparison of three scenarios

A typical freeze-in scenario, but happens for WIMPs



A realistic model with Higgs portal WIMP

SM + two singlet scalars [Kawana, PRD 105, 103515 (2022)]

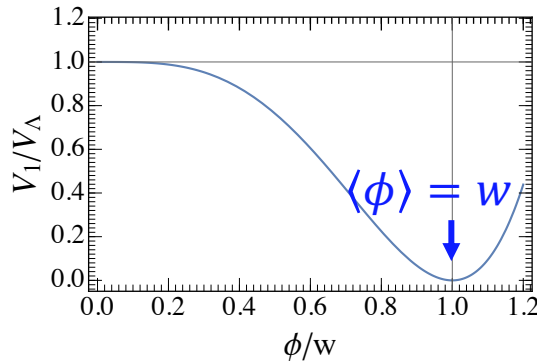
$$V = \lambda_h |H|^4 + \frac{\lambda_\phi}{4} \phi^4 + \lambda_x |X|^4 + \underbrace{\frac{\lambda_{h\phi}}{2} \phi^2 |H|^2 + \frac{\lambda_{\phi x}}{2} \phi^2 |X|^2 + \lambda_{hx} |X|^2 |H|^2}_{\text{Portal couplings}}$$

↑ SM Higgs
↑ FOPT scalar
↑ dark matter candidate

Coleman-Weinberg potential [Coleman et al, PRD 7, 1888 (1973)]

$$V_1 \approx V_\Lambda + \frac{\lambda_{\phi s}^2 - 2y_\psi^4}{64\pi^2} \phi^4 \left(\log \frac{\phi}{w} - \frac{1}{4} \right);$$

Also contribution from fermion $\mathcal{L}_I \supset -\frac{y_\psi}{\sqrt{2}} \phi \bar{\psi} \psi - y_\nu \bar{\ell}_L \tilde{H} \psi$



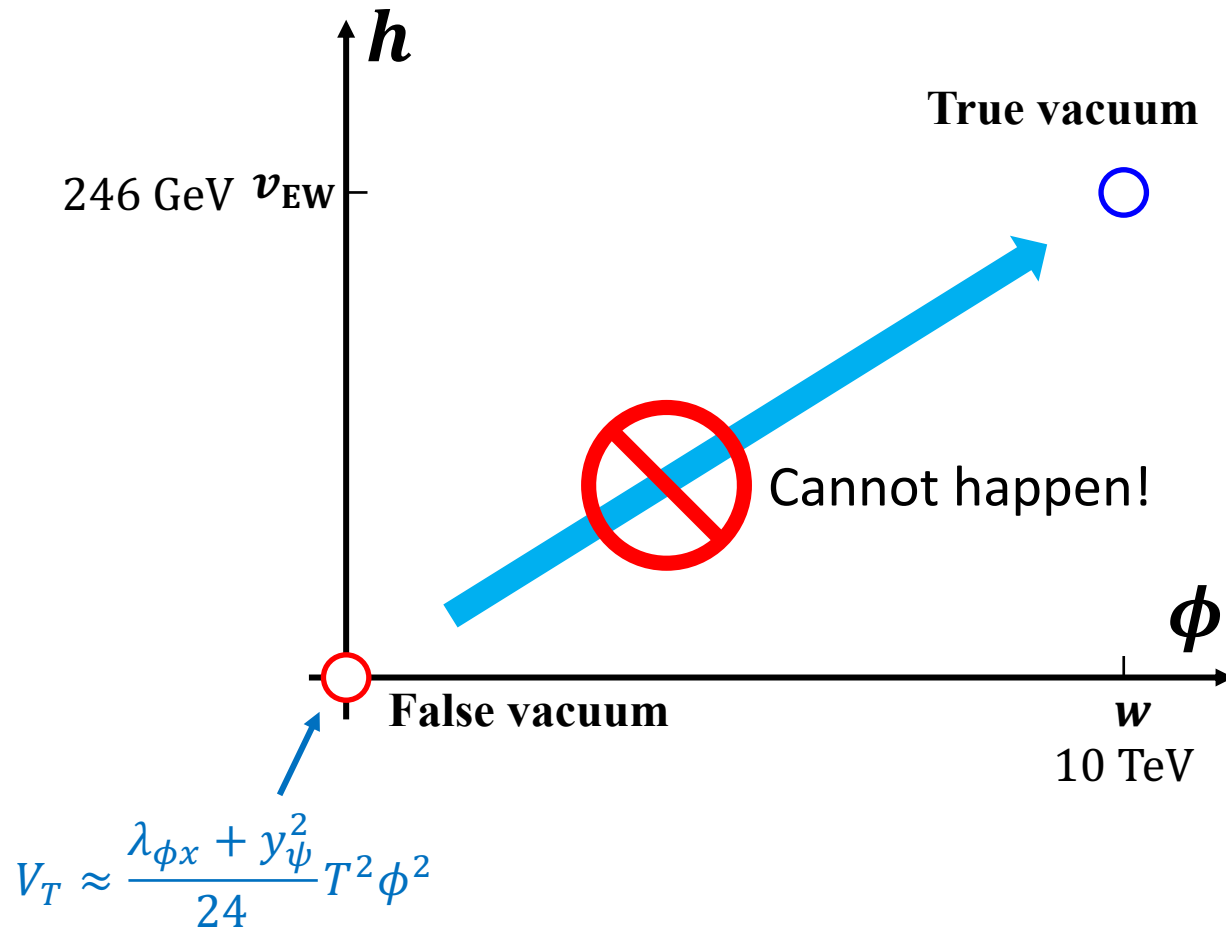
Triggers EW symmetry breaking: $\lambda_{h\phi} = -\frac{m_h^2}{w^2}$

- $V \rightarrow -\frac{m_h^2}{2} |H|^2 + \lambda_h |H|^4$
- $\langle h \rangle = v_{EW} = 246 \text{ GeV}; m_h = 125 \text{ GeV}$
- DM mass $m_X = \sqrt{\frac{\lambda_{\phi x} w^2 + \lambda_{hx} v_{EW}^2}{2}}$

Thermal history (1)

Thermal barrier forbids the direct tunneling

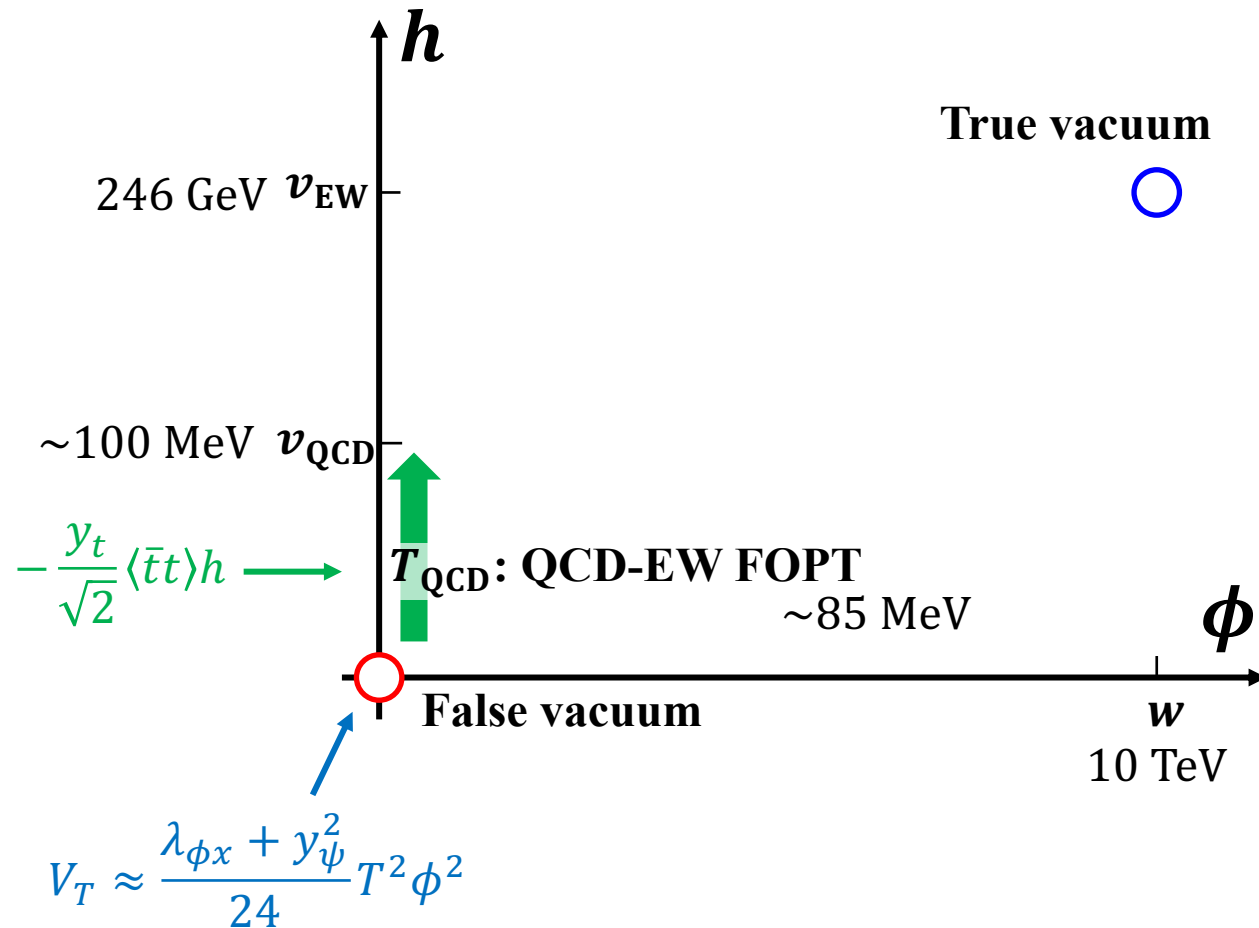
- The Universe is supercooled



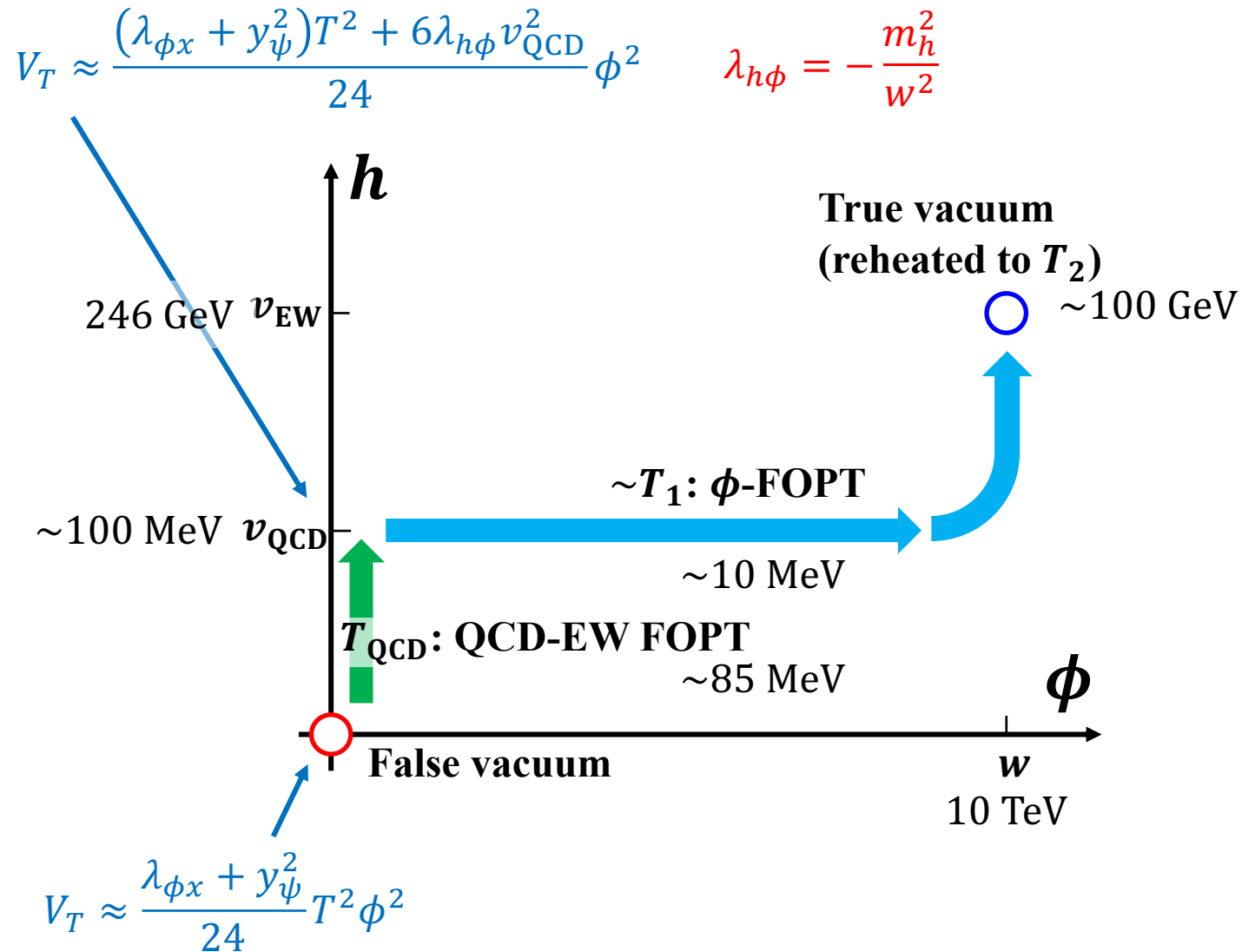
Thermal history (2)

QCD FOPT with 6-flavor massless quarks [\[Pisarski et al, PRD 29, 338 \(1984\)\]](#)

- QCD phase transition triggers EW phase transition



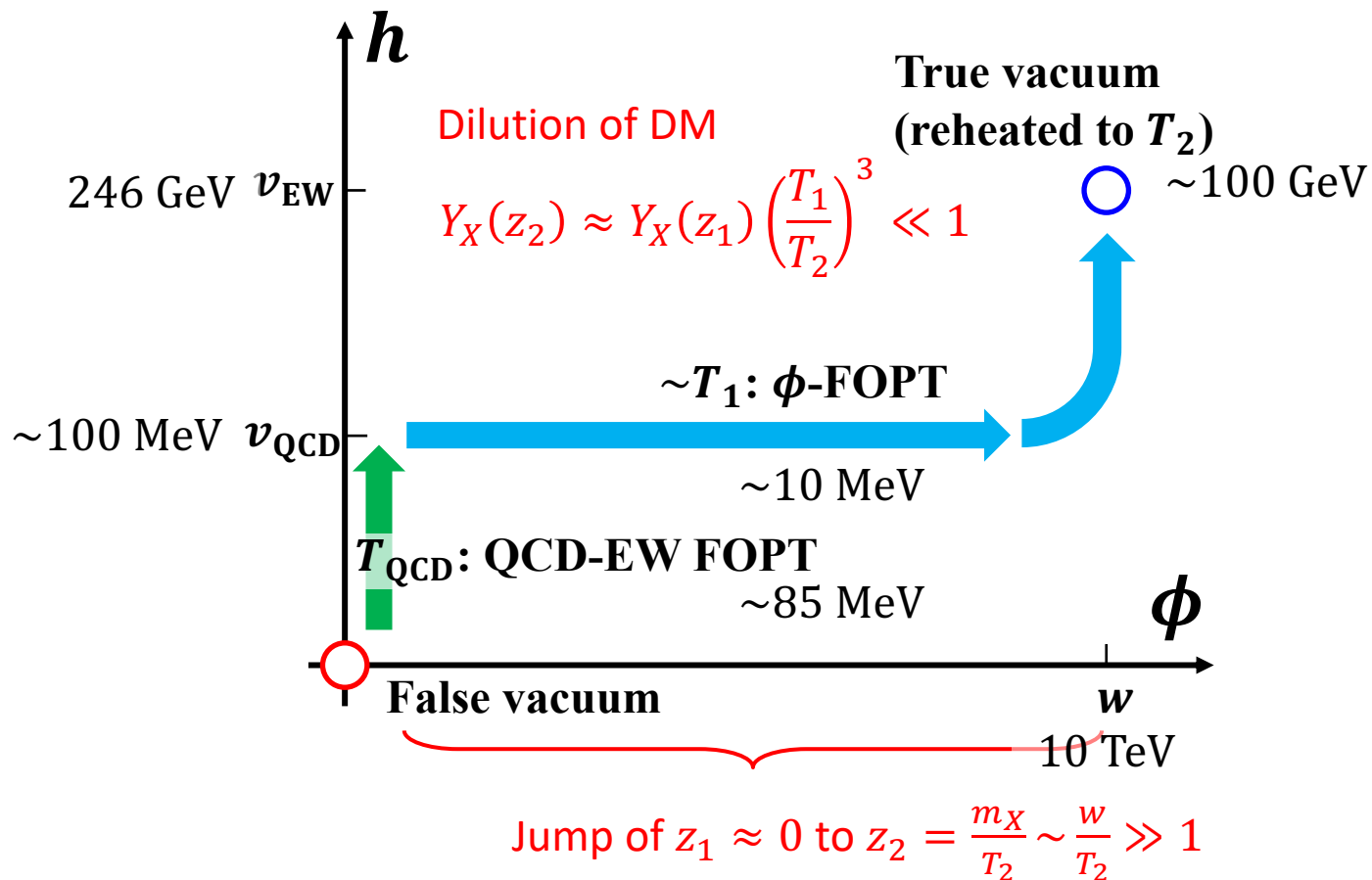
Thermal history (3)



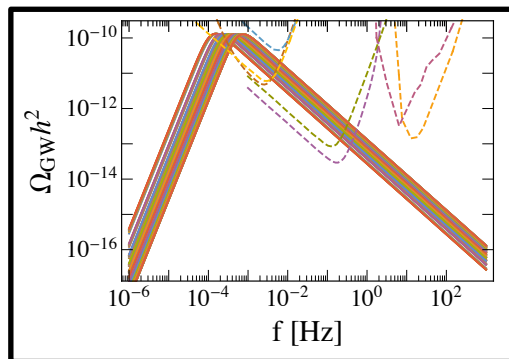
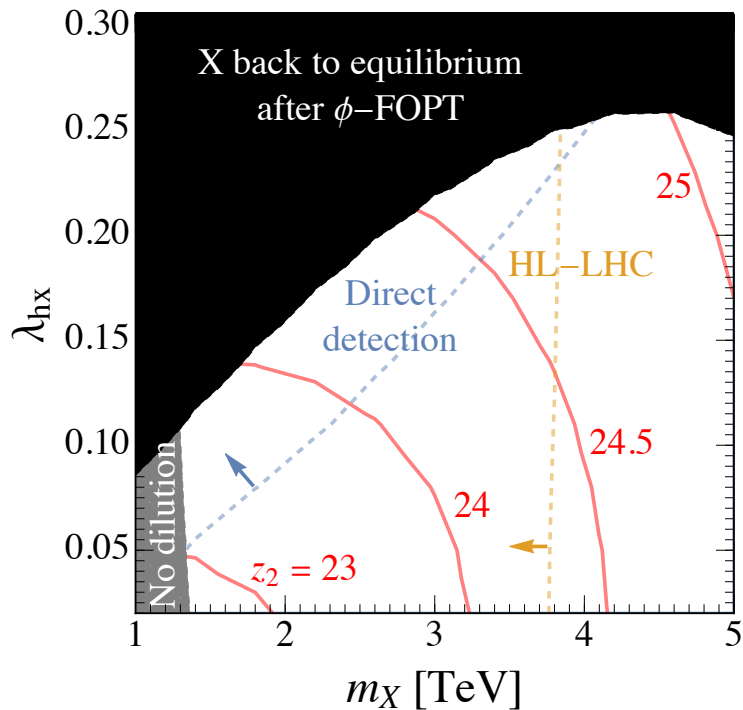
Satisfying conditions of WIMP freeze-in

Satisfying the two conditions of WIMP freeze-in

- Dilution & large z_2

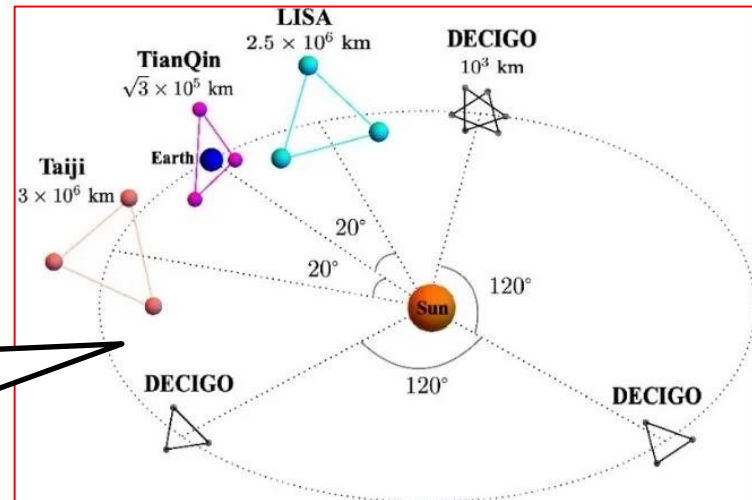


Viable parameter space



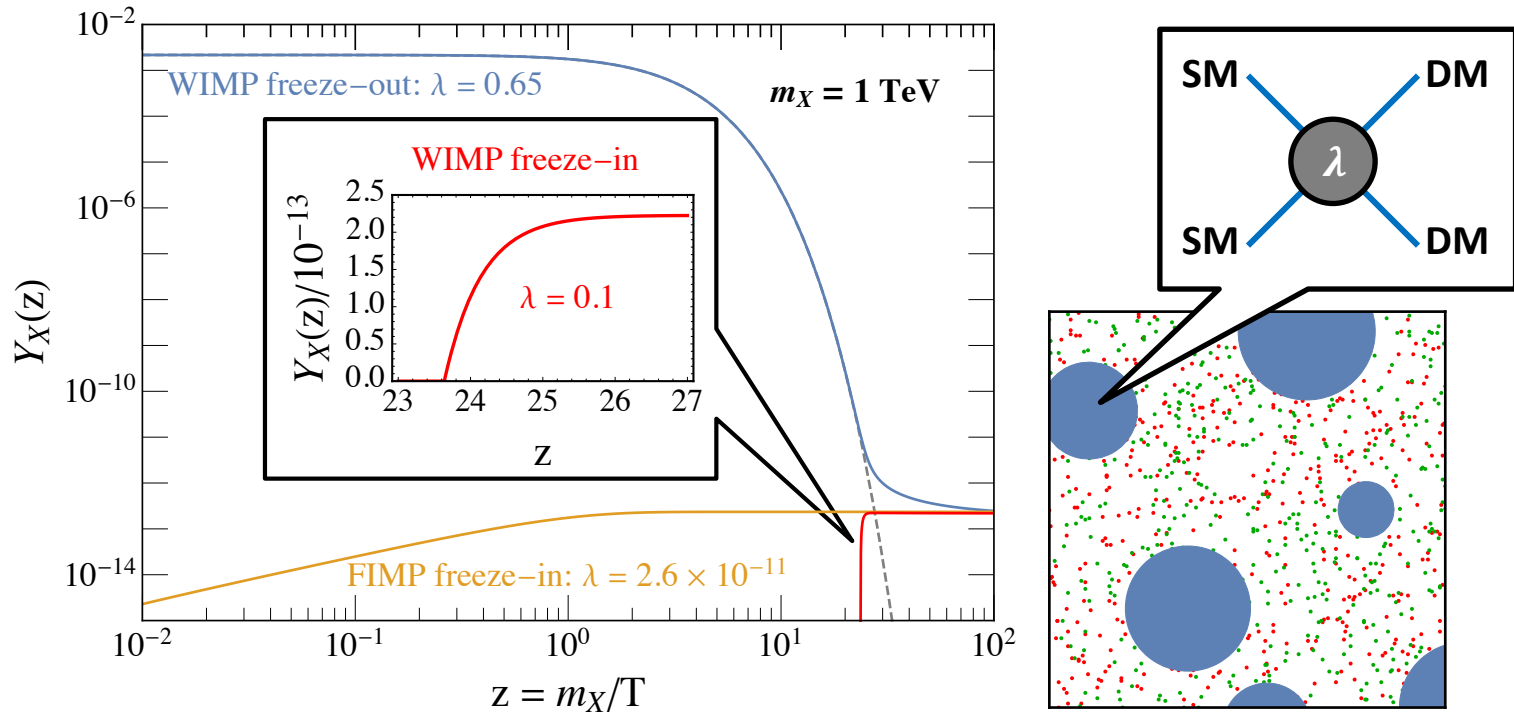
Higgs portal WIMP X , but freeze-in: $hh \rightarrow XX^\dagger$; $\phi\phi \rightarrow XX^\dagger$

- Direct detection $\sigma_{SI} \sim 10^{-48} \text{ cm}^2$ (smaller than standard WIMP)
- $\text{Br}(h \rightarrow \phi\phi) \sim 4\% - 9\%$;
- Gravitational waves $f \sim 10^{-3} \text{ Hz}$, LISA, TianQin, Taiji, ...



Conclusion

A novel dark matter scenario based on the $2 \rightarrow 2$ process



- Generally applied to a lot of new physics models;
- Phenomenology: Correlation between **WIMP searches** & **gravitational waves**

Thank you!

Backup: guide for model building (1)

Dark matter

$$\mathcal{L} \supset -\lambda B^\dagger B X^\dagger X - \frac{\lambda_{\phi X}}{2} \phi^2 |X|^2;$$

Thermal bath particle FOPT scalar

Minimal setup:

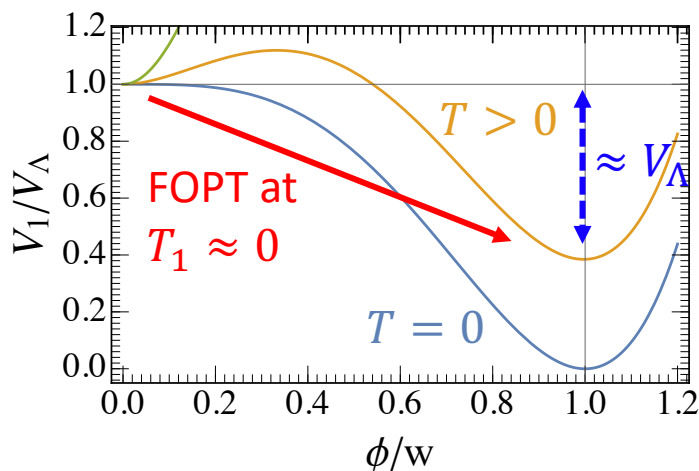
$$B = \frac{\phi}{\sqrt{2}} \text{ and hence } \lambda_{\phi X} = \lambda$$

Classically conformal principle $V_{\text{tree}} = \lambda_\phi \phi^4 / 4$

One-loop level: Coleman-Weinberg potential

$$V_1(\phi) = V_\Lambda + \frac{\lambda_B^2}{64\pi^2} \phi^4 \left(\log \frac{\phi}{w} - \frac{1}{4} \right)$$

$\Omega_X h^2$ determined by λ and λ_B



- $m_X^2 = \frac{\lambda}{2} w^2$
- $V_\Lambda = \frac{\lambda_B^2}{256\pi^2} w^4$
- $\frac{\pi^2}{30} g_* T_2^4 \approx V_\Lambda$

$$z_2 \equiv \frac{m_X}{T_2}$$

- $\Omega_X h^2 \approx 0.1 (1 + 2z_2) \left(\frac{\lambda e^{-z_2}}{3.5 \times 10^{-11}} \right)^2$

Backup: guide for model building (2)

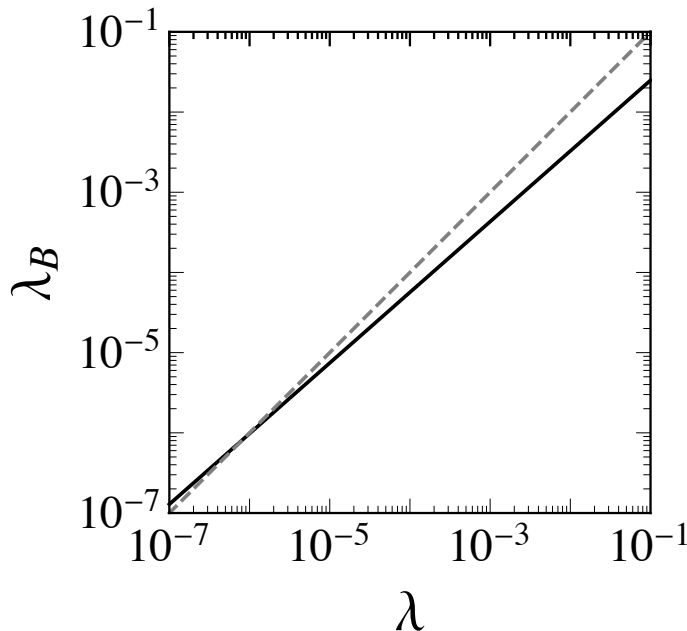
Dark matter

$$\mathcal{L} \supset -\frac{\lambda}{2} \phi^2 X^\dagger X - \frac{\lambda}{2} \phi^2 |X|^2; \quad V_1(\phi) = V_\Lambda + \frac{\lambda_B^2}{64\pi^2} \phi^4 \left(\log \frac{\phi}{w} - \frac{1}{4} \right).$$

↑ Thermal bath particle
 ↑

Classically conformal principle

$\Omega_X h^2$ determined by λ and λ_B



$$\Omega_{\text{DM}} h^2 = 0.12--$$

$$\text{Numerical: } \lambda_B \approx 0.189 \lambda^{0.881}$$

Minimal setup: $\lambda_B \approx \lambda$

$$\Rightarrow \lambda \sim 10^{-6};$$

If $\lambda \sim 0.1$ (WIMP)

$$\Rightarrow \lambda_B \approx 0.024;$$

\Rightarrow **additional fermions included**