



I FOUND THE HUGS BISON.

95 GeV Excess Theory Insight

Sven Heinemeyer, IFT (CSIC, Madrid)

Shanghai/zoom, 09/2023

1. Evidence for a light Higgs boson
2. Possible model interpretation
3. SUSY realizations
4. Physics opportunities at future e^+e^- colliders
5. Conclusions

1. Evidence for a light Higgs boson



Measurement of Higgs boson production and search for new resonances in final states with photons and Z bosons

by Chiara Arcangeletti (INFN e Laboratori Nazionali di Frascati (IT))

📅 Tuesday 6 Jun 2023, 11:00 → 12:00 Europe/Zurich

📍 500/1-001 - Main Auditorium (CERN)

New ATLAS result on the low-mass Higgs search in $pp \rightarrow \phi \rightarrow \gamma\gamma$

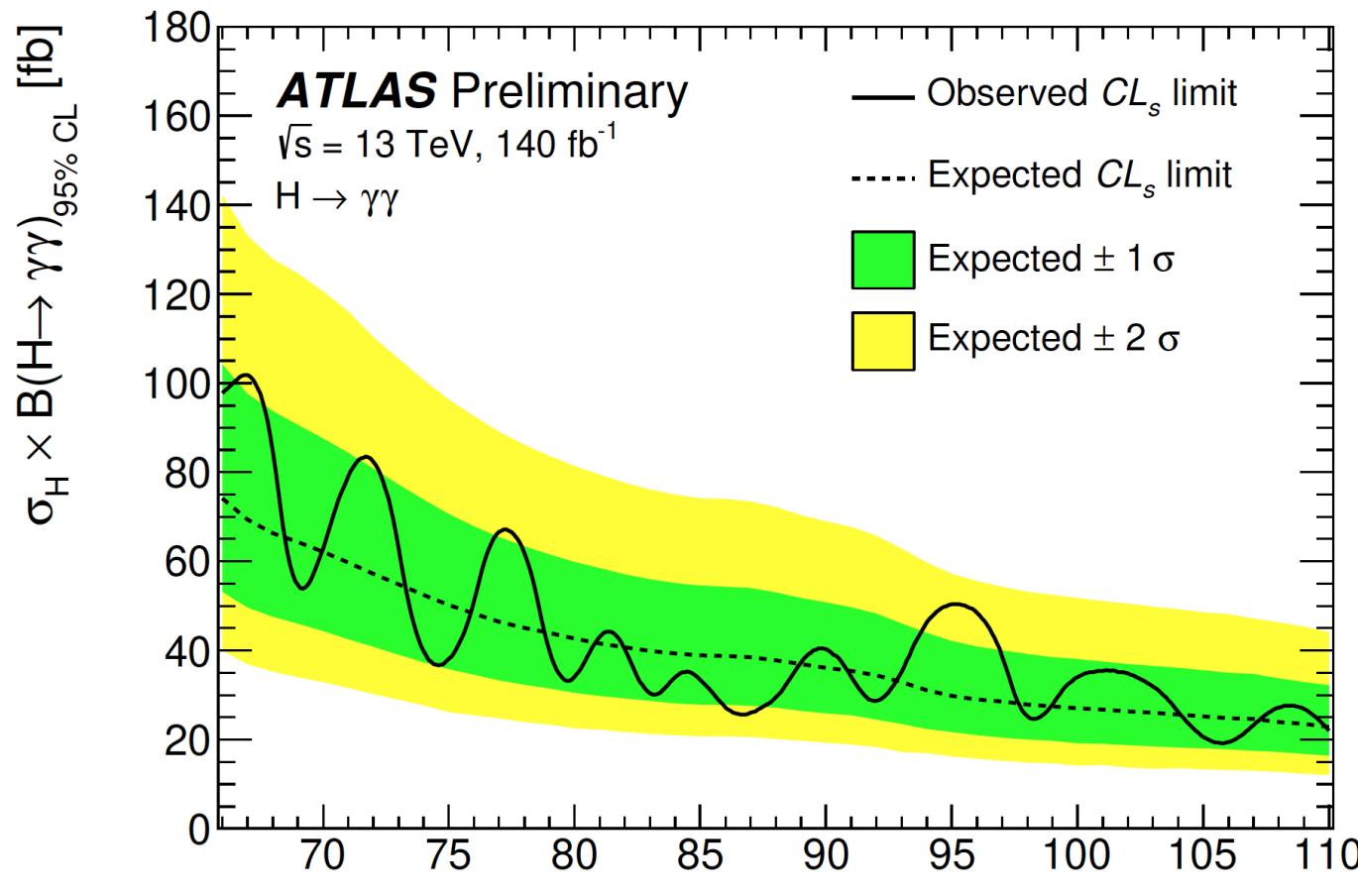
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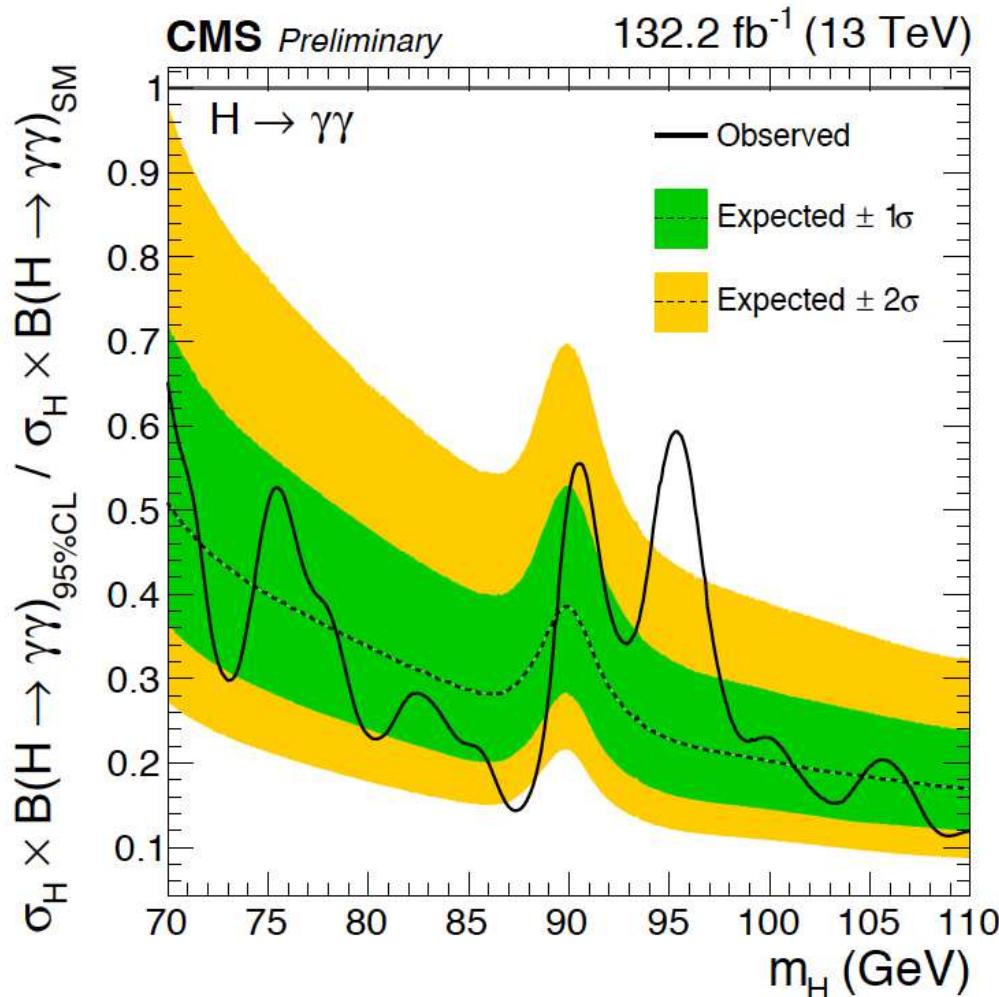
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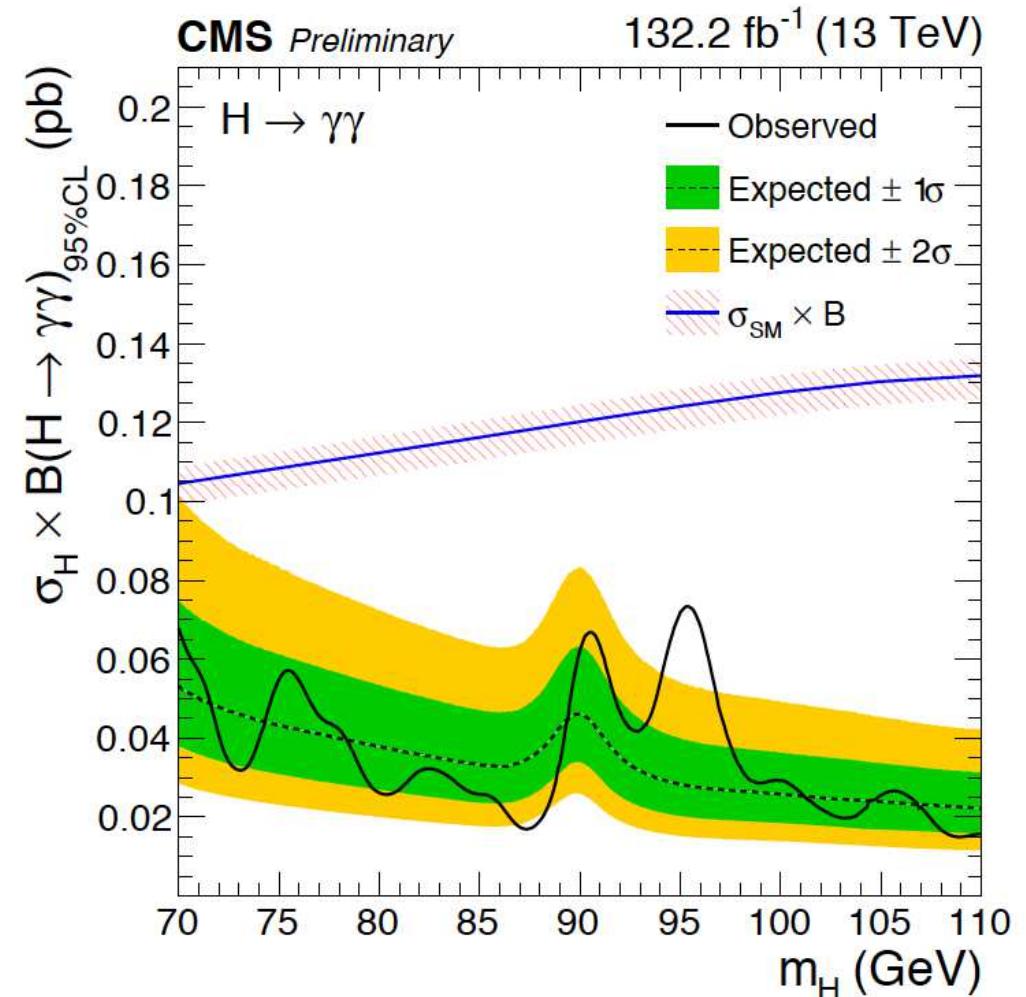


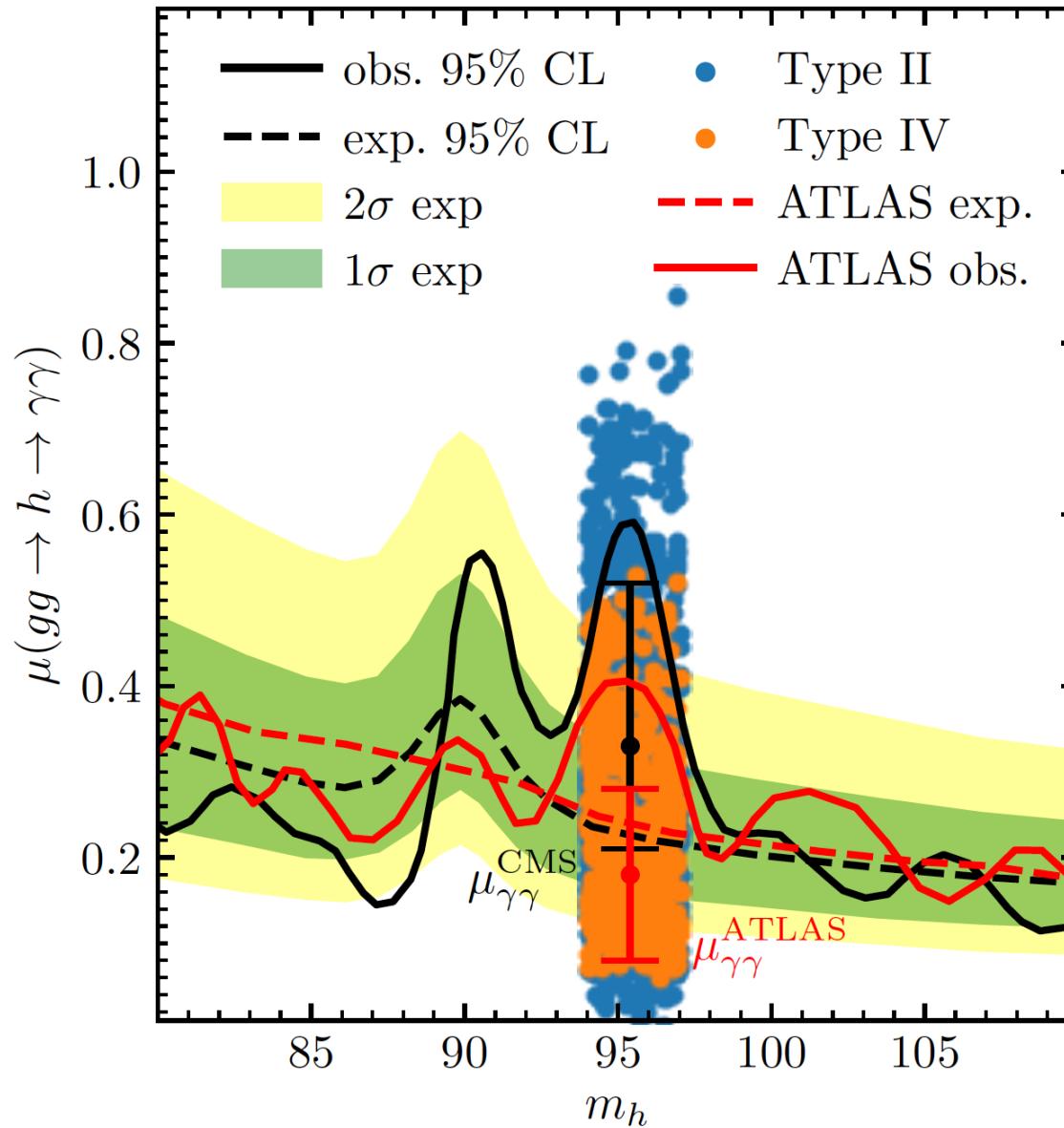
⇒ highest “excess” at 95.4 GeV



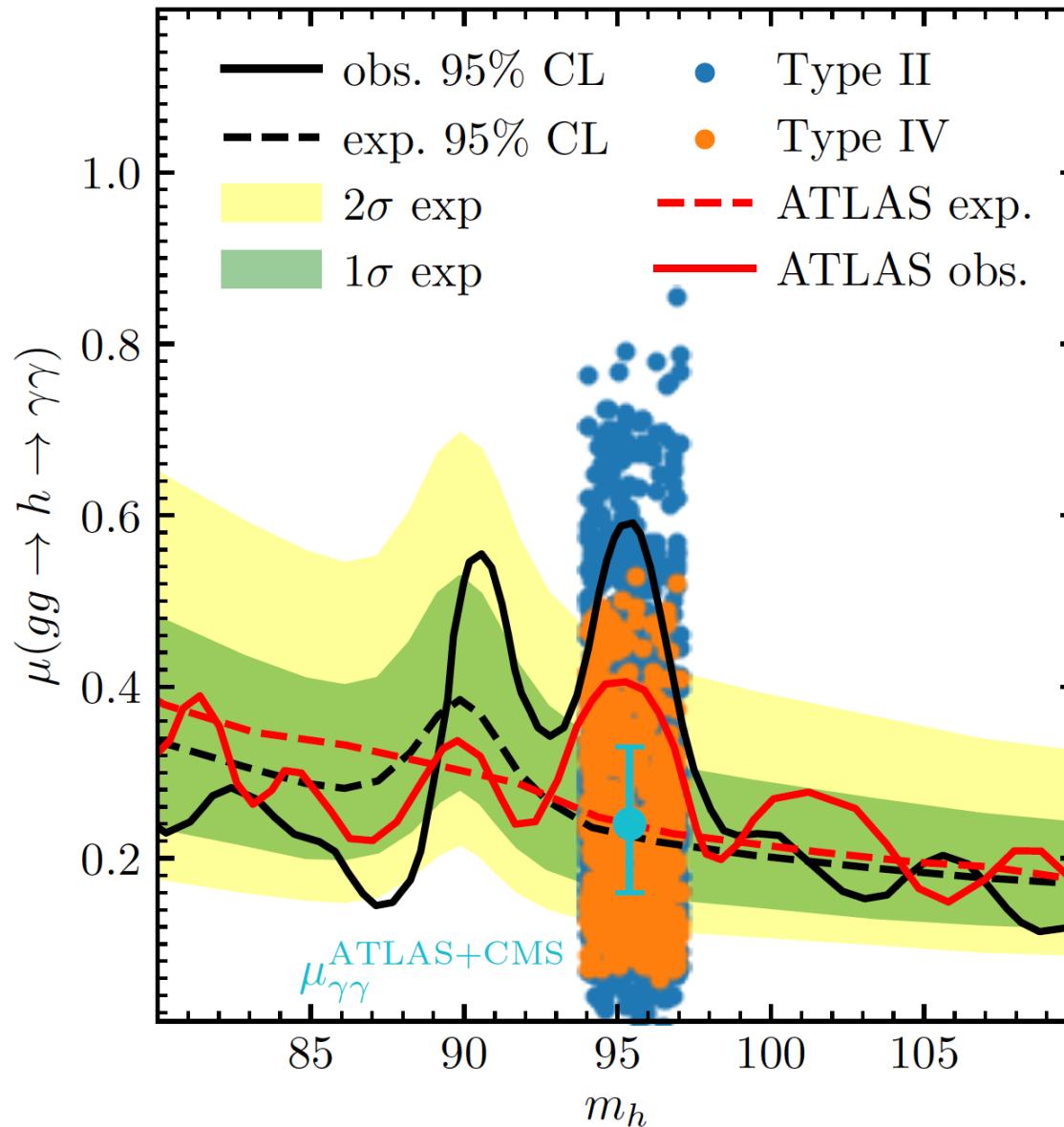
$$\mu_{\gamma\gamma}^{\text{CMS}} = [\sigma(gg \rightarrow h_{95}) \times \text{BR}(h_{95} \rightarrow \gamma\gamma)]_{\text{exp/SM}} = 0.33^{+0.19}_{-0.12} \text{ (2.9 } \sigma\text{)}$$

$$\mu_{\gamma\gamma}^{\text{ATLAS}} = 0.18 \pm 0.10 \text{ (1.7 } \sigma\text{)}$$





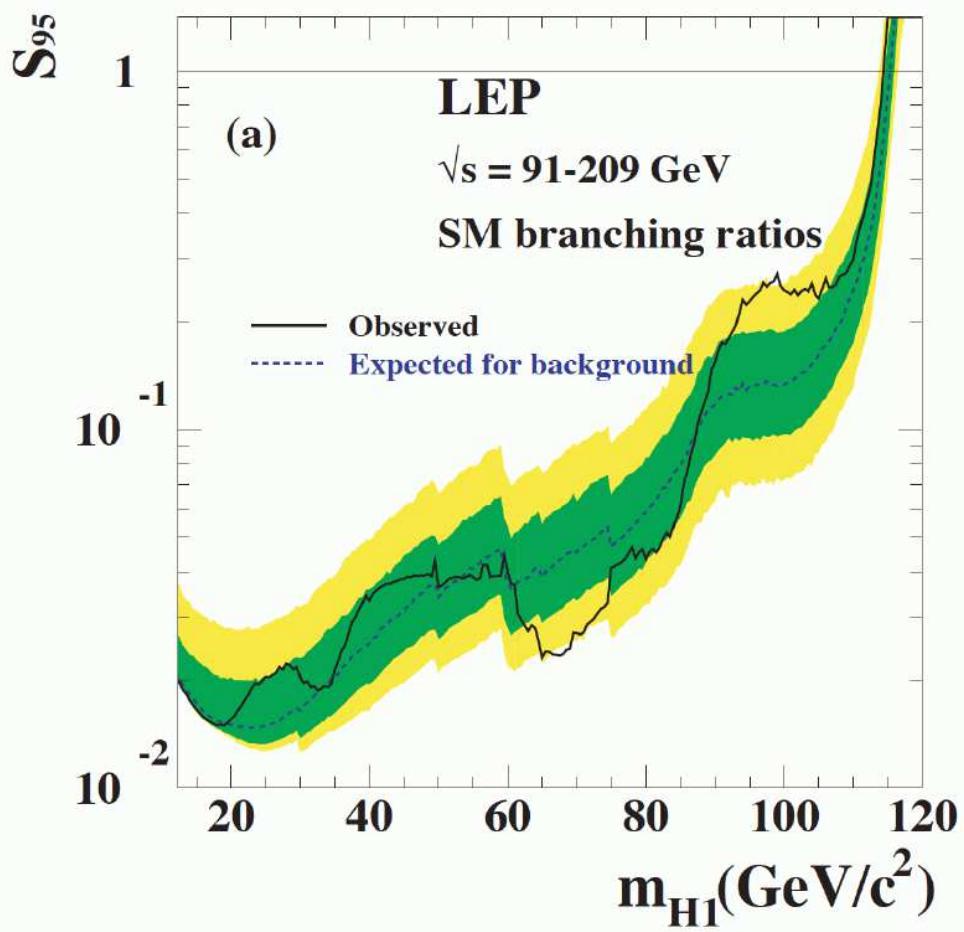
⇒ agreement between ATLAS and CMS!



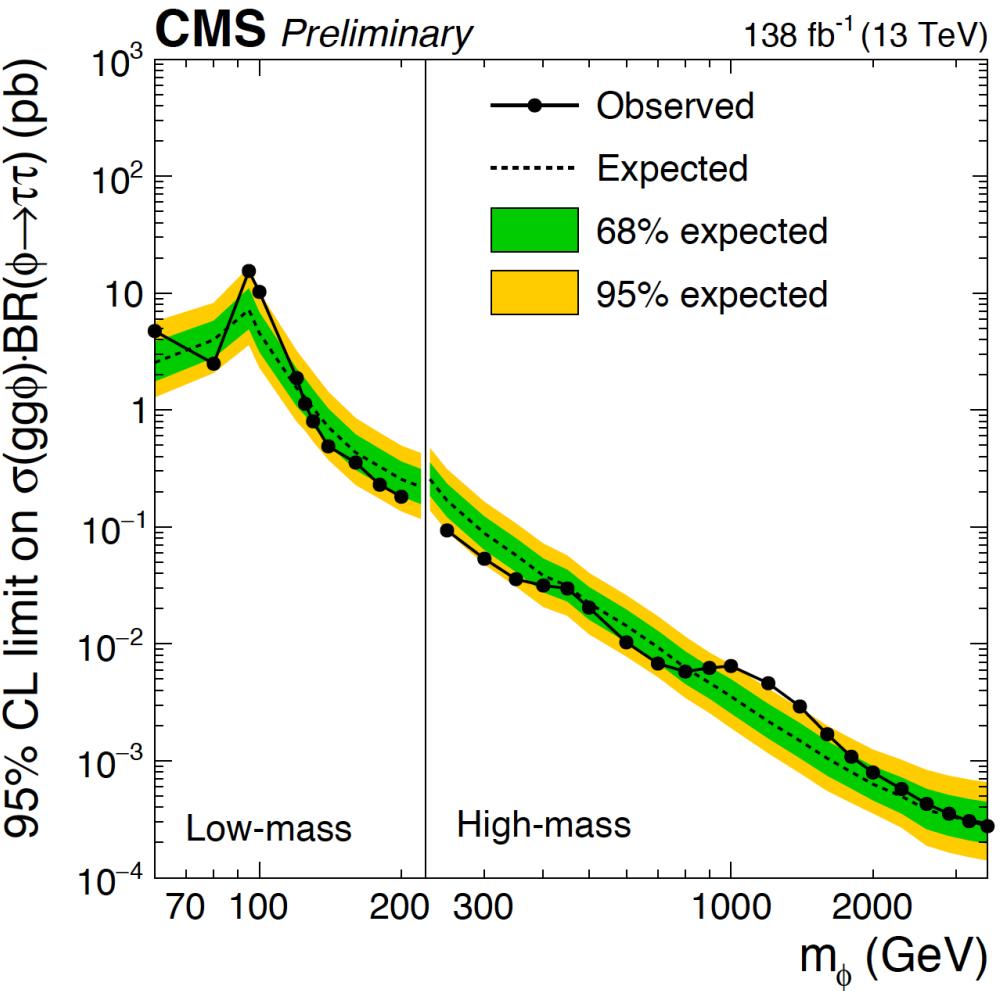
⇒ agreement between ATLAS and CMS!

$$\mu_{\gamma\gamma} = 0.24^{+0.09}_{-0.08} \text{ (3.1 }\sigma\text{)}$$

LEP: $e^+e^- \rightarrow Z\phi \rightarrow Zb\bar{b}$ (2σ)



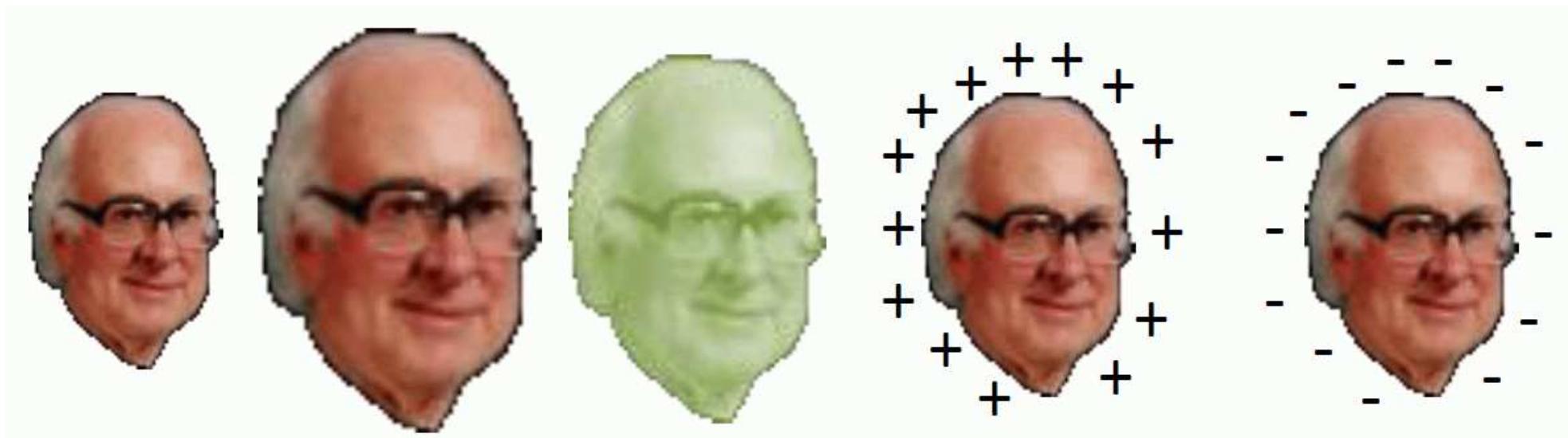
CMS: $pp \rightarrow \phi \rightarrow \tau^+\tau^-$ (2.4σ)



⇒ no LEE (as theorist I am allowed to add naively)

⇒ $\sim 4.6\sigma$

2. Possible model interpretation



Models: Summary

Authors	Model	arXiv	Excesses	Comments
Cao, Guo, He et al.	nNMSSM	1612.08522	$bb + \gamma\gamma$	
Fox, Weiner	2HDM + VL	1710.07649	$bb + \gamma\gamma$	
Haisch, Malinauskas	2HDM	1712.06599	$bb + (\gamma\gamma)$	
TB, Heinemeyer, Muñoz	$\mu\nu$ SSM	1712.07475	$bb + \gamma\gamma$	EW seesaw
Liu, Liu, Wagner, Wang	$U(1)_{L_\mu - L_\tau}$	1805.01476	$bb + \gamma\gamma$	B-anomalies
Domingo, Heinemeyer, Paßehr, Weiglein	NMSSM	1807.06322	$bb + \gamma\gamma$	
Hollik, Liebler, Moortgat-Pick et al.	μ NMSSM	1809.07371	$bb + \gamma\gamma$	Inflation
TB, Chakraborti, Heinemeyer	N2HDM	1903.11661	$bb + \gamma\gamma$	
Cline, Toma	pNG + squarks	1906.02175	$bb + \gamma\gamma$	DM
Choi, Hui Im, Sik Jeong et al.	gNMSSM	1906.03389	$bb + \gamma\gamma$	
Cao, Jia, Yue et al.	nNMSSM	1908.07206	$bb + \gamma\gamma$	Type-I seesaw
Aguilar-Saavedra, Joaquim	SM + $U(1)_Y'$	2002.07697	$bb + \gamma\gamma$	
TB, Olea-Romacho	S2HDM	2108.10864	$bb + \gamma\gamma$	DM, GC excess
TB, Grohsjean, Heinemeyer et al.	NMSSM	2109.01128	$\gamma\gamma$	400 GeV excess
Heinemeyer, Lika, Moortgat-Pick et al.	2HDM+s	2112.11958	$bb + \gamma\gamma$	
TB, Heinemeyer, Weiglein	N2HDM	2203.13180	$bb + (\tau\tau) + \gamma\gamma$	
TB, Heinemeyer, Weiglein	N2HDM	2204.05975	$bb + (\tau\tau) + \gamma\gamma$	CDF M_W
Benbrik, Boukidi, Moretti et al.	A2HDM-III	2204.07470	$bb + \gamma\gamma$	LFV

Green: 2HDM(+X), blue: Susy, red: Extra charged fields

Models: Summary

Authors	Model	arXiv	Excesses	Comments
TB, Heinemeyer, Weiglein Azevedo, TB, Ferreira	S2HDM C2HDM	2303.12018 2305.19716	$bb + (\tau\tau) + \gamma\gamma$ $bb + \tau\tau + \gamma\gamma$	DM
Bonilla, Carcamo, Kovalenko et al.	Left-Right model	2305.11967	$\gamma\gamma$	DM
TB, Heinemeyer, Weiglein	S2HDM	2306.03889	$bb + (\tau\tau) + \gamma\gamma$	ATLAS- $\gamma\gamma$
Escribano, Martín Lozano, Vicente	Scotogenic	2306.03735	$bb + \gamma\gamma$	DM, ν masses
Belyaev, Benbrik, Boukidi et al.	A2HDM	2306.09029	$bb + (\tau\tau) + \gamma\gamma$	
Ashanuman, Banik, Coloretti et al.	$Y = 0$ triplet	2306.15722	$\gamma\gamma$	CDF M_W
Aguilar-Saavedra, Camara, Joaquim et al.	UN2HDM	2307.03768	$(\tau\tau), \gamma\gamma$	

Green: 2HDM(+X), blue: Susy, red: Extra charged fields

Our example: N2HDM

[S2HDM similar with complex singlet]

Fields:

$$\Phi_1 = \begin{pmatrix} \phi_1^+ \\ \frac{1}{\sqrt{2}}(v_1 + \rho_1 + i\eta_1) \end{pmatrix}, \quad \Phi_2 = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}}(v_2 + \rho_2 + i\eta_2) \end{pmatrix}, \quad \Phi_S = v_S + \rho_S$$

Potential:

$$\begin{aligned} V = & m_{11}^2 |\Phi_1|^2 + m_{22}^2 |\Phi_2|^2 - m_{12}^2 (\Phi_1^\dagger \Phi_2 + h.c.) + \frac{\lambda_1}{2} (\Phi_1^\dagger \Phi_1)^2 + \frac{\lambda_2}{2} (\Phi_2^\dagger \Phi_2)^2 \\ & + \lambda_3 (\Phi_1^\dagger \Phi_1)(\Phi_2^\dagger \Phi_2) + \lambda_4 (\Phi_1^\dagger \Phi_2)(\Phi_2^\dagger \Phi_1) + \frac{\lambda_5}{2} [(\Phi_1^\dagger \Phi_2)^2 + h.c.] \\ & + \frac{1}{2} m_S^2 \Phi_S^2 + \frac{\lambda_6}{8} \Phi_S^4 + \frac{\lambda_7}{2} (\Phi_1^\dagger \Phi_1) \Phi_S^2 + \frac{\lambda_8}{2} (\Phi_2^\dagger \Phi_2) \Phi_S^2 \end{aligned}$$

Z_2 symmetry: $\Phi_1 \rightarrow \Phi_1$, $\Phi_2 \rightarrow -\Phi_2$, $\Phi_S \rightarrow \Phi_S$

Z'_2 symmetry: $\Phi_1 \rightarrow \Phi_1$, $\Phi_2 \rightarrow \Phi_2$, $\Phi_S \rightarrow -\Phi_S$ (broken by $v_S \Rightarrow$ no DM)

Physical states: h_1 , h_2 , h_3 (\mathcal{CP} -even), A (\mathcal{CP} -odd), H^\pm (charged)

Extension of the Z_2 symmetry to fermions determines four types:

	u -type	d -type	leptons
type I	Φ_2	Φ_2	Φ_2
type II	Φ_2	Φ_1	Φ_1
type III (lepton-specific)	Φ_2	Φ_2	Φ_1
type IV (flipped)	Φ_2	Φ_1	Φ_2

⇒ exactly as in 2HDM

Three neutral \mathcal{CP} -even Higgses:

$$\begin{pmatrix} h_1 \\ h_2 \\ h_3 \end{pmatrix} = R \begin{pmatrix} \rho_1 \\ \rho_2 \\ \rho_S \end{pmatrix}, \quad R = \begin{pmatrix} c_{\alpha_1}c_{\alpha_2} & s_{\alpha_1}c_{\alpha_2} & s_{\alpha_2} \\ -(c_{\alpha_1}s_{\alpha_2}s_{\alpha_3} + s_{\alpha_1}c_{\alpha_3}) & c_{\alpha_1}c_{\alpha_3} - s_{\alpha_1}s_{\alpha_2}s_{\alpha_3} & c_{\alpha_2}s_{\alpha_3} \\ -c_{\alpha_1}s_{\alpha_2}c_{\alpha_3} + s_{\alpha_1}s_{\alpha_3} & -(c_{\alpha_1}s_{\alpha_3} + s_{\alpha_1}s_{\alpha_2}c_{\alpha_3}) & c_{\alpha_2}c_{\alpha_3} \end{pmatrix}$$

Coupling to massive gauge bosons: (identical for all four types)

$$\begin{array}{c} \hline c_{h_i VV} = c_\beta R_{i1} + s_\beta R_{i2} \\ \hline h_1 & c_{\alpha_2} c_{\beta - \alpha_1} \\ h_2 & -c_{\beta - \alpha_1} s_{\alpha_2} s_{\alpha_3} + c_{\alpha_3} s_{\beta - \alpha_1} \\ h_3 & -c_{\alpha_3} c_{\beta - \alpha_1} s_{\alpha_2} - s_{\alpha_3} s_{\beta - \alpha_1} \\ \hline \end{array}$$

Coupling to fermions: (same pattern as in 2HDM)

	u -type ($c_{h_i tt}$)	d -type ($c_{h_i bb}$)	leptons ($c_{h_i \tau\tau}$)
type I	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i2}}{s_\beta}$
type II	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i1}}{c_\beta}$	$\frac{R_{i1}}{c_\beta}$
type III (lepton-specific)	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i1}}{c_\beta}$
type IV (flipped)	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i1}}{c_\beta}$	$\frac{R_{i2}}{s_\beta}$

“Physical” input parameters:

$$\alpha_{1,2,3}, \quad \tan \beta, \quad v, \quad v_S, \quad m_{h_{1,2,3}}, \quad m_A, \quad M_{H^\pm}, \quad m_{12}^2$$

Needed to fit the $\gamma\gamma$ and $b\bar{b}$ excesses: $m_{h_1} \sim 95$ GeV, $m_{h_2} \sim 125$ GeV

- $c_{h_1 VV}^2$ strongly reduced for $\mu_{b\bar{b}}$
- $c_{h_1 bb}$ reduced to enhance $\text{BR}(h_1 \rightarrow \gamma\gamma)$
- $c_{h_1 tt}$ not reduced for $\mu_{\gamma\gamma}$

	Decrease $c_{h_1 b\bar{b}}$	No decrease $c_{h_1 t\bar{t}}$	No enhancement $c_{h_1 \tau\bar{\tau}}$
type I	$(\frac{R_{12}}{s_\beta})$:-)	$(\frac{R_{12}}{s_\beta})$:-()	$(\frac{R_{12}}{s_\beta})$
type II	$(\frac{R_{11}}{c_\beta})$:-)	$(\frac{R_{12}}{s_\beta})$:-)	$(\frac{R_{11}}{c_\beta})$
type III	$(\frac{R_{12}}{s_\beta})$:-)	$(\frac{R_{12}}{s_\beta})$:-()	$(\frac{R_{11}}{c_\beta})$
type IV	$(\frac{R_{11}}{c_\beta})$:-)	$(\frac{R_{12}}{s_\beta})$:-)	$(\frac{R_{12}}{s_\beta})$

Type II and IV: $c_{h_1 bb}$ and $c_{h_1 tt}$ independent

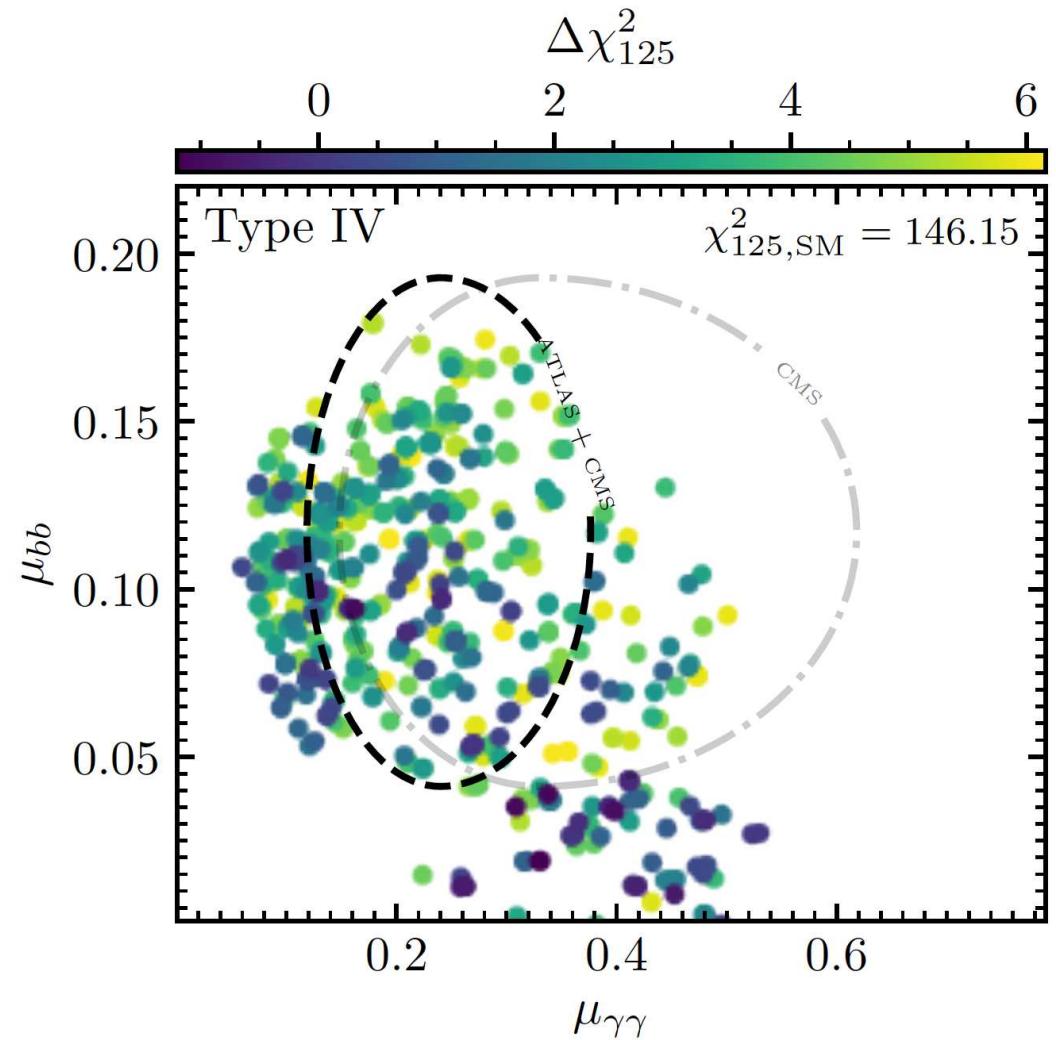
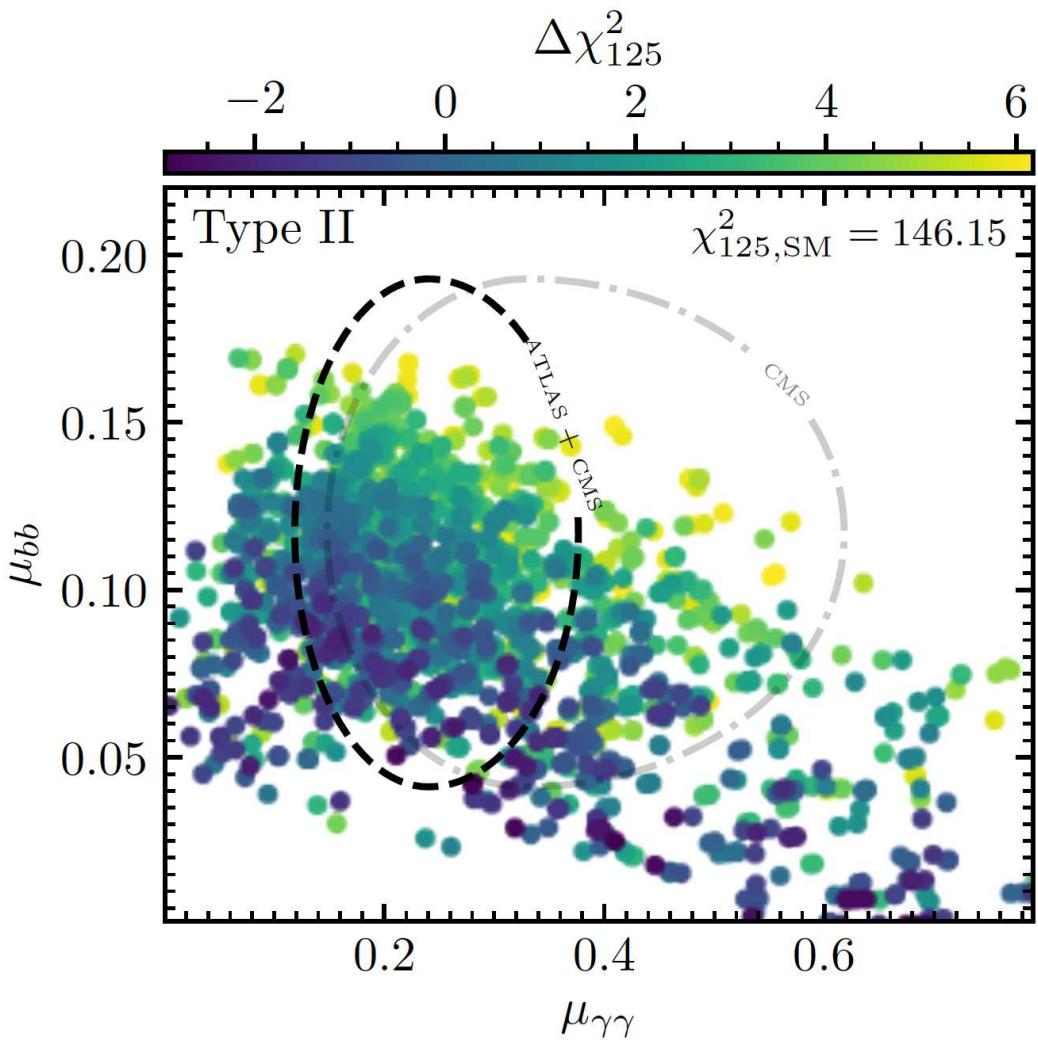
Type II vs. IV: $c_{h_1 \tau\bar{\tau}}$ can be suppressed or enhanced

⇒ only type II and IV can fit the $\gamma\gamma$ and $b\bar{b}$ excesses

⇒ $\tau\tau$ excess may decide between type II and IV

S2HDM type II vs. type IV

[T. Biekötter, S.H., G. Weiglein '23]

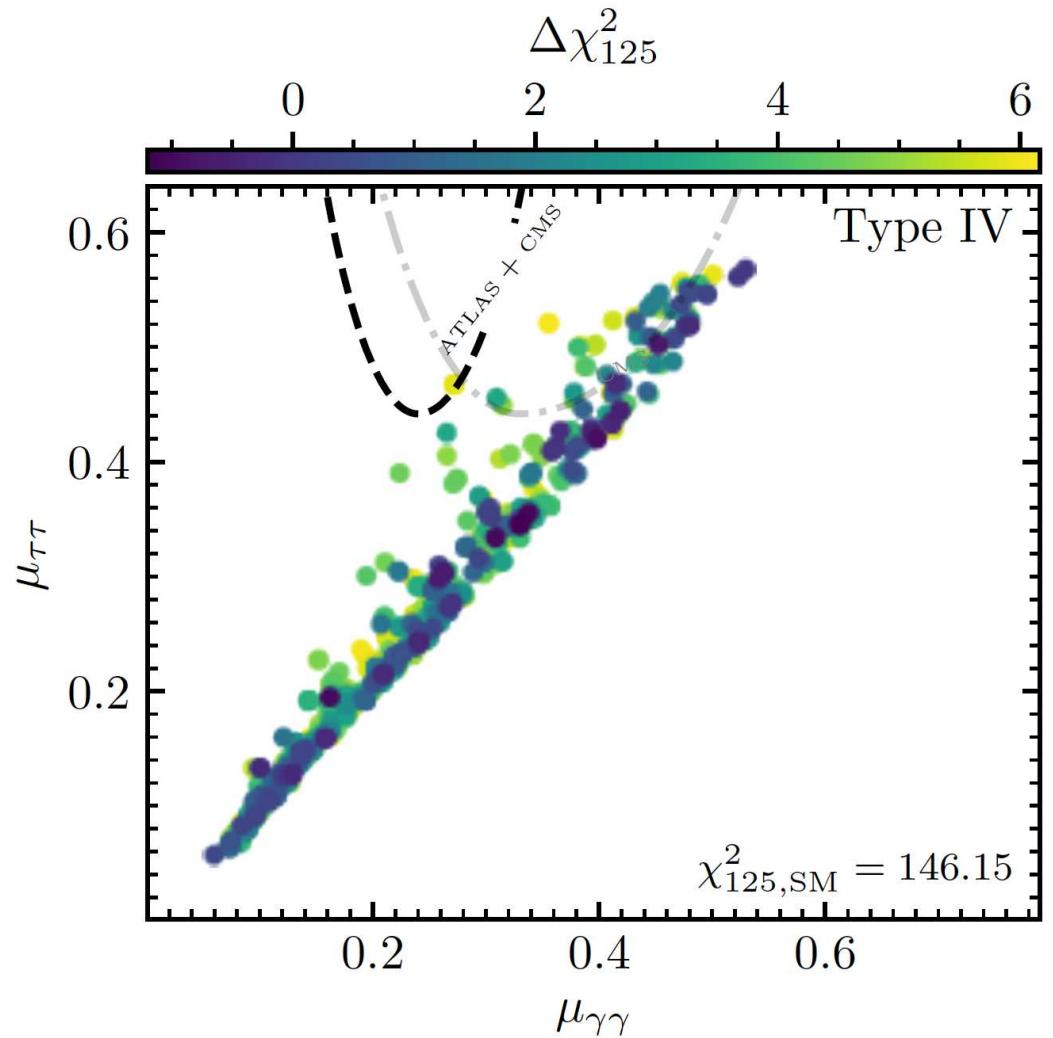
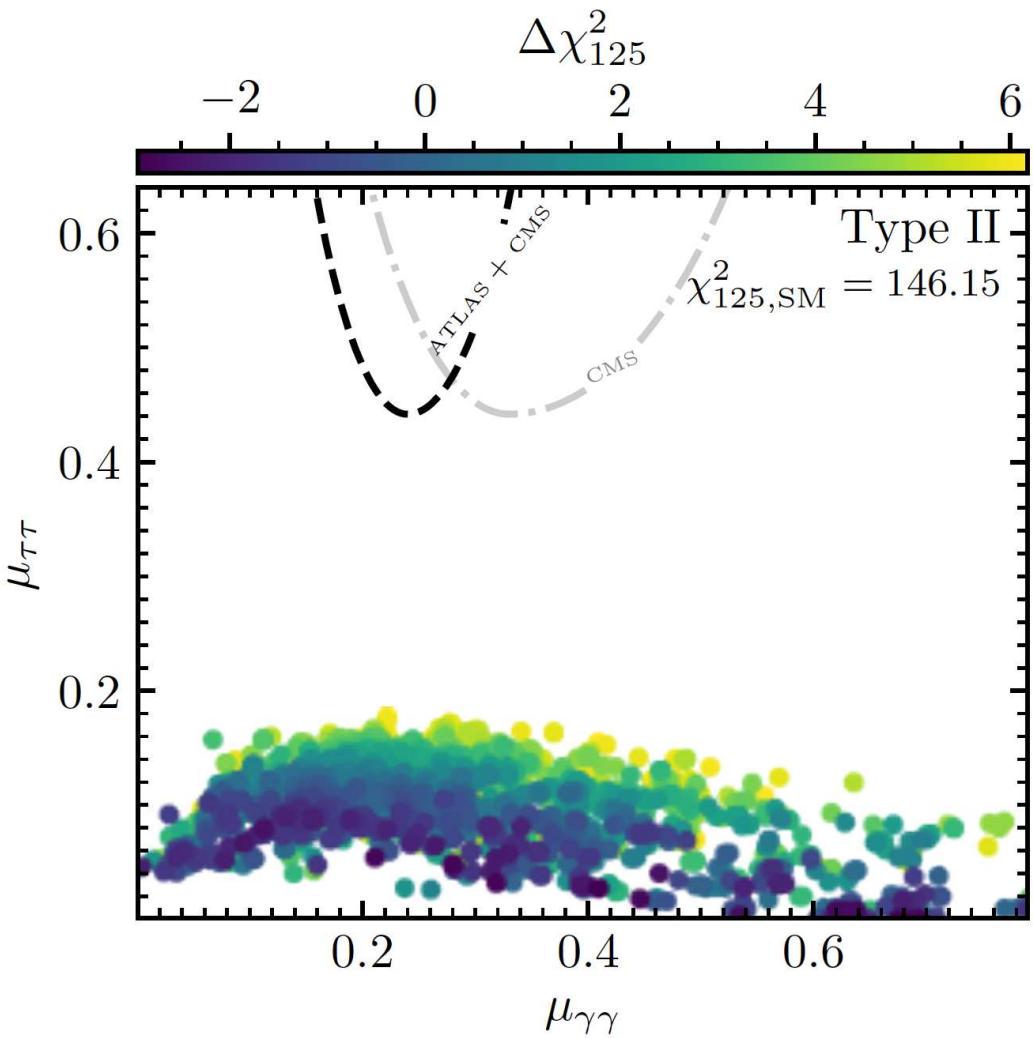


Color coding: χ^2_{125} from [HiggsSignals](#)

⇒ both type II and IV can fit the $\gamma\gamma$ and bb excesses

S2HDM type II vs. type IV

[T. Biekötter, S.H., G. Weiglein '23]



Color coding: χ^2_{125} from [HiggsSignals](#)

⇒ only type IV can fit marginally the $\gamma\gamma$ and $\tau\tau$ excesses

3. SUSY realizations

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Q: Can the models fit the excesses **despite** the additional SUSY constraints on the Higgs sector **???**

What about the NMSSM?

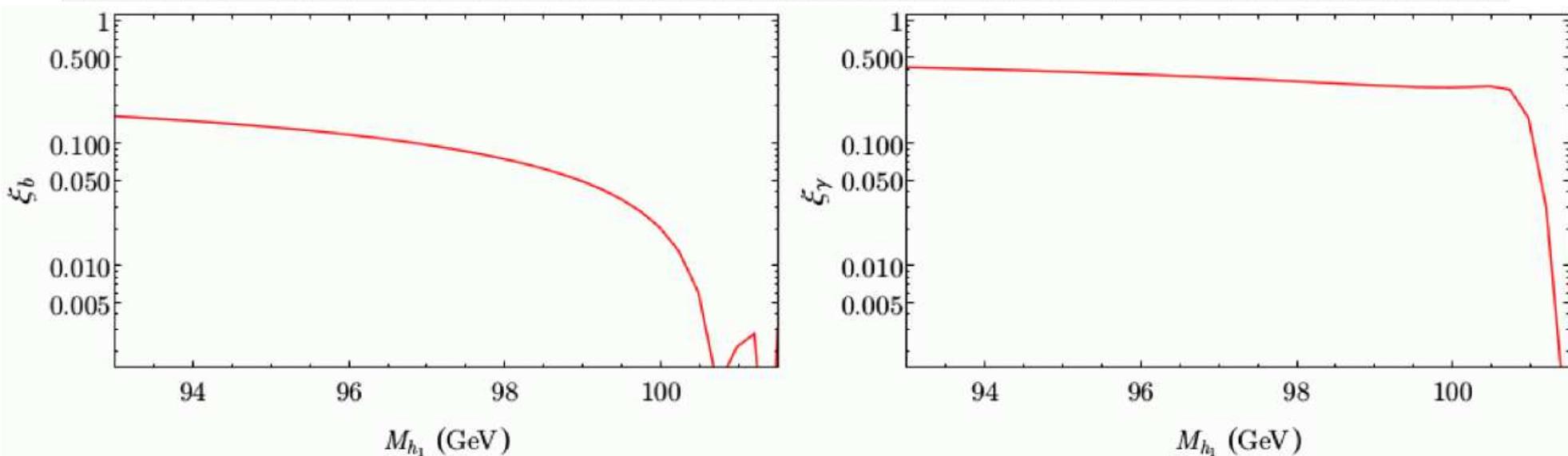
[F. Domingo, S.H., S. Passehr, G. Weiglein '18]

Parameters:

$\lambda = 0.6$, $\kappa = 0.035$, $\tan\beta = 2$, $\mu_{\text{eff}} = (397 + 15x)$ GeV, $M_{H^\pm} = 1$ TeV,
 $A_\kappa = -325$ GeV, $M_{\text{SUSY}} = 1$ TeV, $A_t = A_b = 0$

$$\xi_b \equiv \frac{\Gamma[h_1 \rightarrow ZZ] \cdot \text{BR}[h_1 \rightarrow b\bar{b}]}{\Gamma[H_{\text{SM}}(M_{h_1}) \rightarrow ZZ] \cdot \text{BR}[H_{\text{SM}}(M_{h_1}) \rightarrow b\bar{b}]} \sim \frac{\sigma[e^+e^- \rightarrow Z(h_1 \rightarrow b\bar{b})]}{\sigma[e^+e^- \rightarrow Z(H_{\text{SM}}(M_{h_1}) \rightarrow b\bar{b})]}$$

$$\xi_\gamma \equiv \frac{\Gamma[h_1 \rightarrow gg] \cdot \text{BR}[h_1 \rightarrow \gamma\gamma]}{\Gamma[H_{\text{SM}}(M_{h_1}) \rightarrow gg] \cdot \text{BR}[H_{\text{SM}}(M_{h_1}) \rightarrow \gamma\gamma]} \sim \frac{\sigma[gg \rightarrow h_1 \rightarrow \gamma\gamma]}{\sigma[gg \rightarrow H_{\text{SM}}(M_{h_1}) \rightarrow \gamma\gamma]}.$$



⇒ both excesses can be fitted simultaneously well with new $\mu_{\gamma\gamma}$!

What about the $\mu\nu$ SSM?

$\mu\nu$ SSM: [D. Lopez-Fogliani, C. Muñoz '06]

$\mu\nu$ SSM: NMSSM + well motivated RPV (in simple terms)
⇒ EW scale seesaw to reproduce the neutrino data

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⇒ EW scale seesaw to reproduce the neutrino data

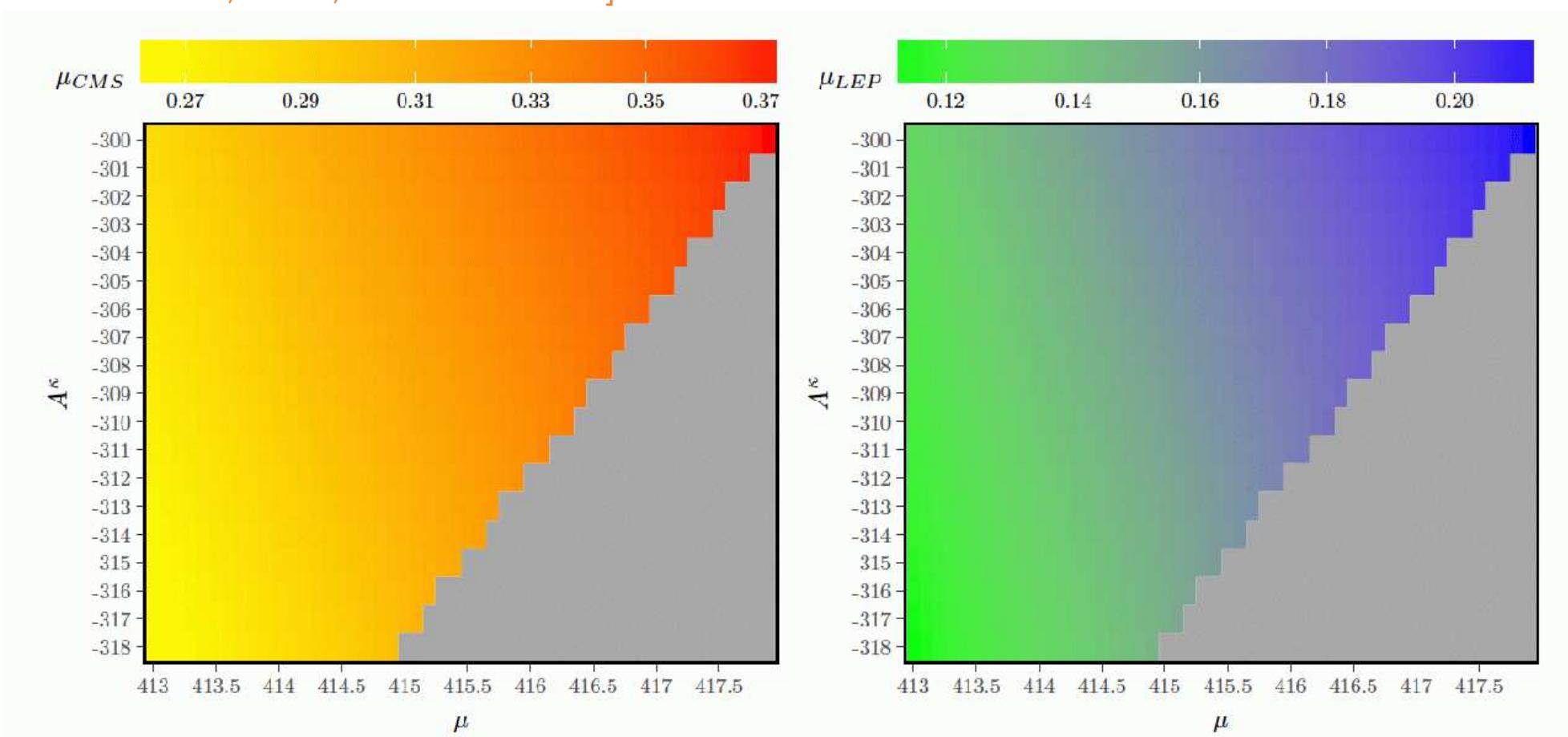
Can the $\mu\nu$ SSM explain the two excesses?

[T. Biekötter, S.H., C. Muñoz '17]

v_{iL}	Y_i^ν	A_i^ν	$\tan \beta$	μ	λ	A^λ	κ	A^κ	M_1
$\sqrt{2} \cdot 10^{-5}$	10^{-7}	-1000	2	[413; 418]	0.6	956.035	0.035	[-300; -318]	100
M_2	M_3	$m_{\tilde{Q}_{iL}}^2$	$m_{\tilde{u}_{iR}}^2$	$m_{\tilde{d}_{iR}}^2$	A_1^u	$A_{2,3}^{u,d}$	$(m_e^2)_{ii}$	A_{33}^e	$A_{11,22}^e$
200	1500	800^2	800^2	800^2	0	0	800^2	0	0

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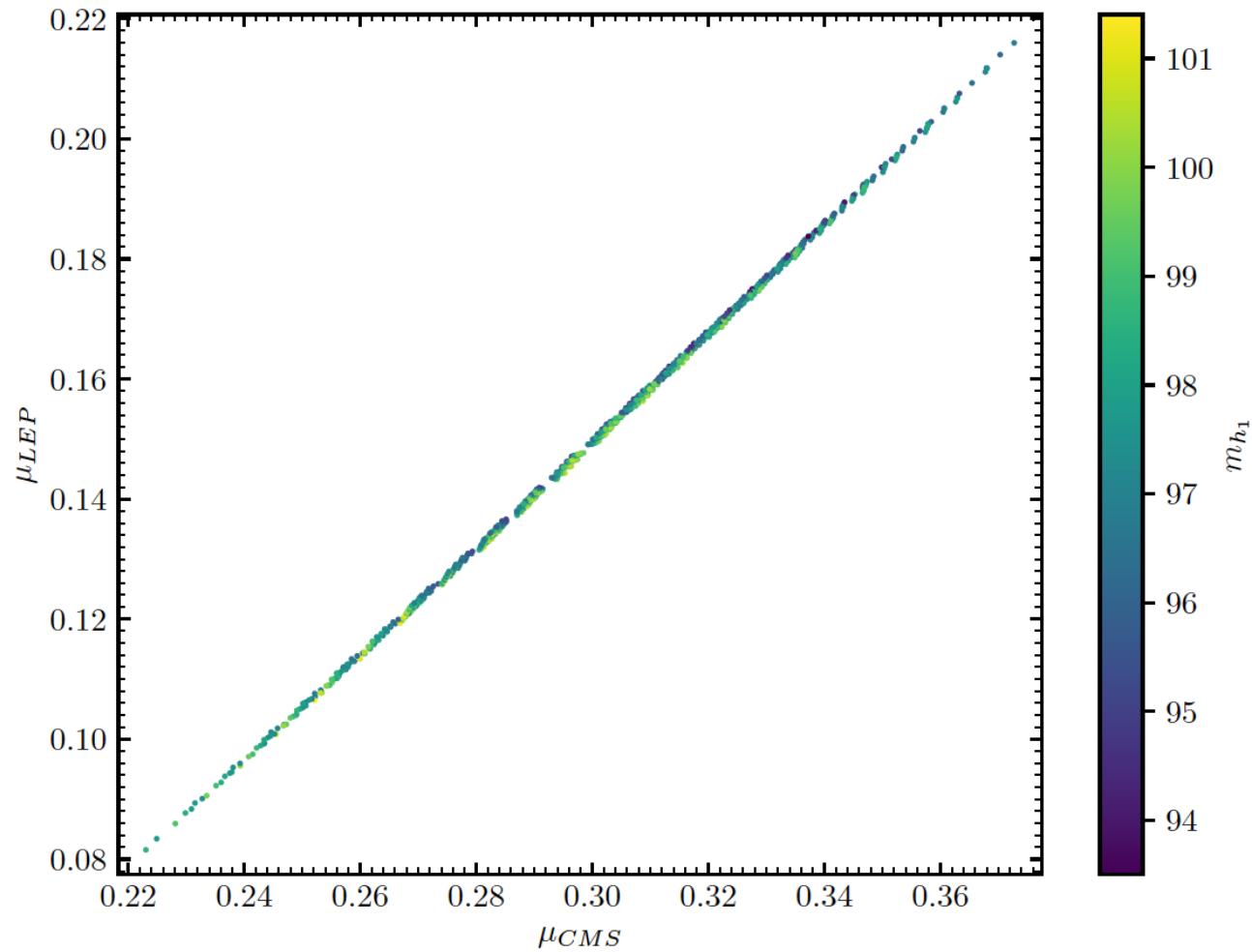


⇒ Yes! :-)

using the new $\mu_{\gamma\gamma}$!

Why does SUSY prefer the new $\mu_{\gamma\gamma}$?

[T. Biekötter, S.H., C. Muñoz '19]



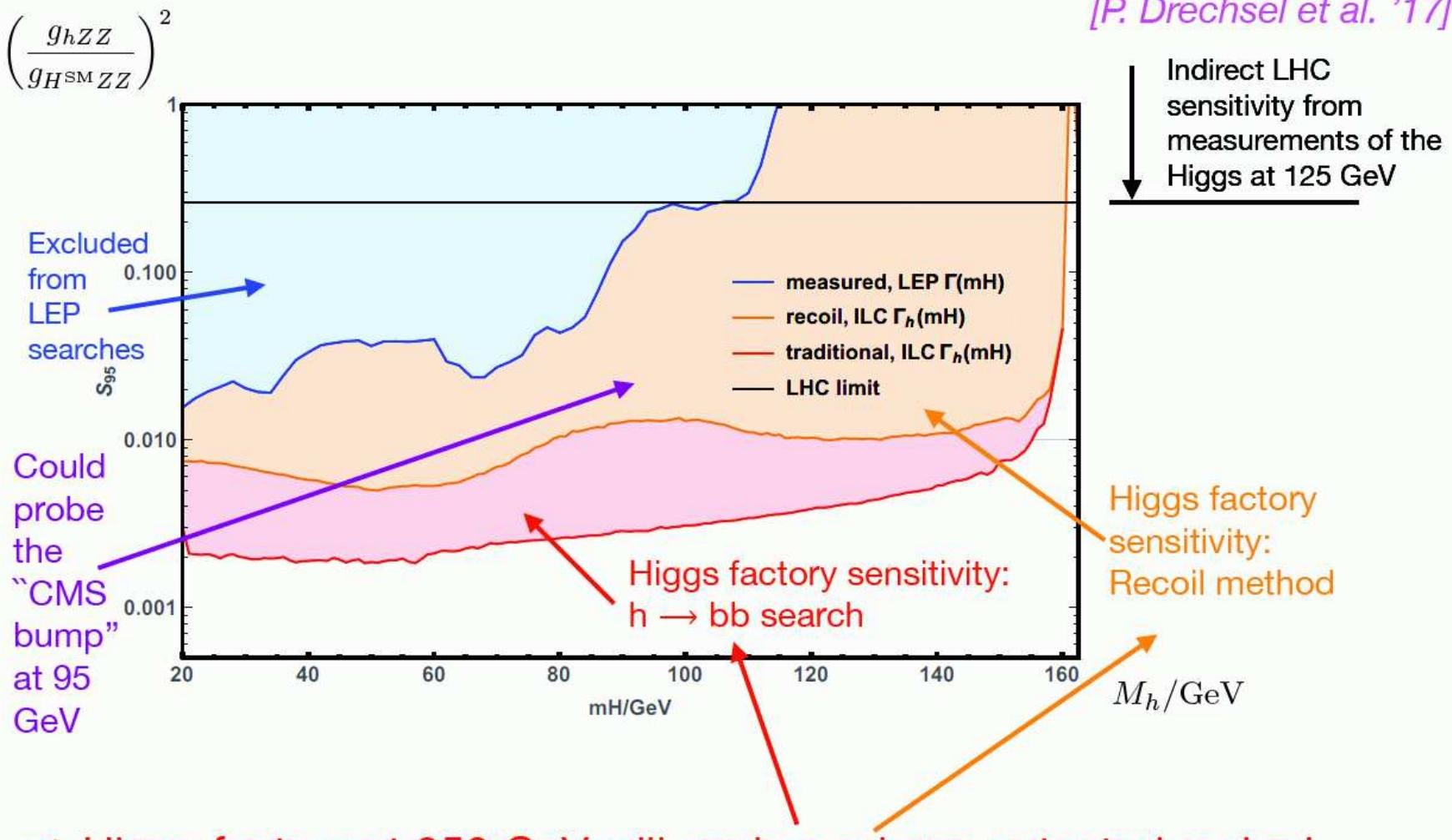
⇒ SUSY enforces strong correlation!

⇒ LEP excess enforces $\mu_{\gamma\gamma} \lesssim 0.35$

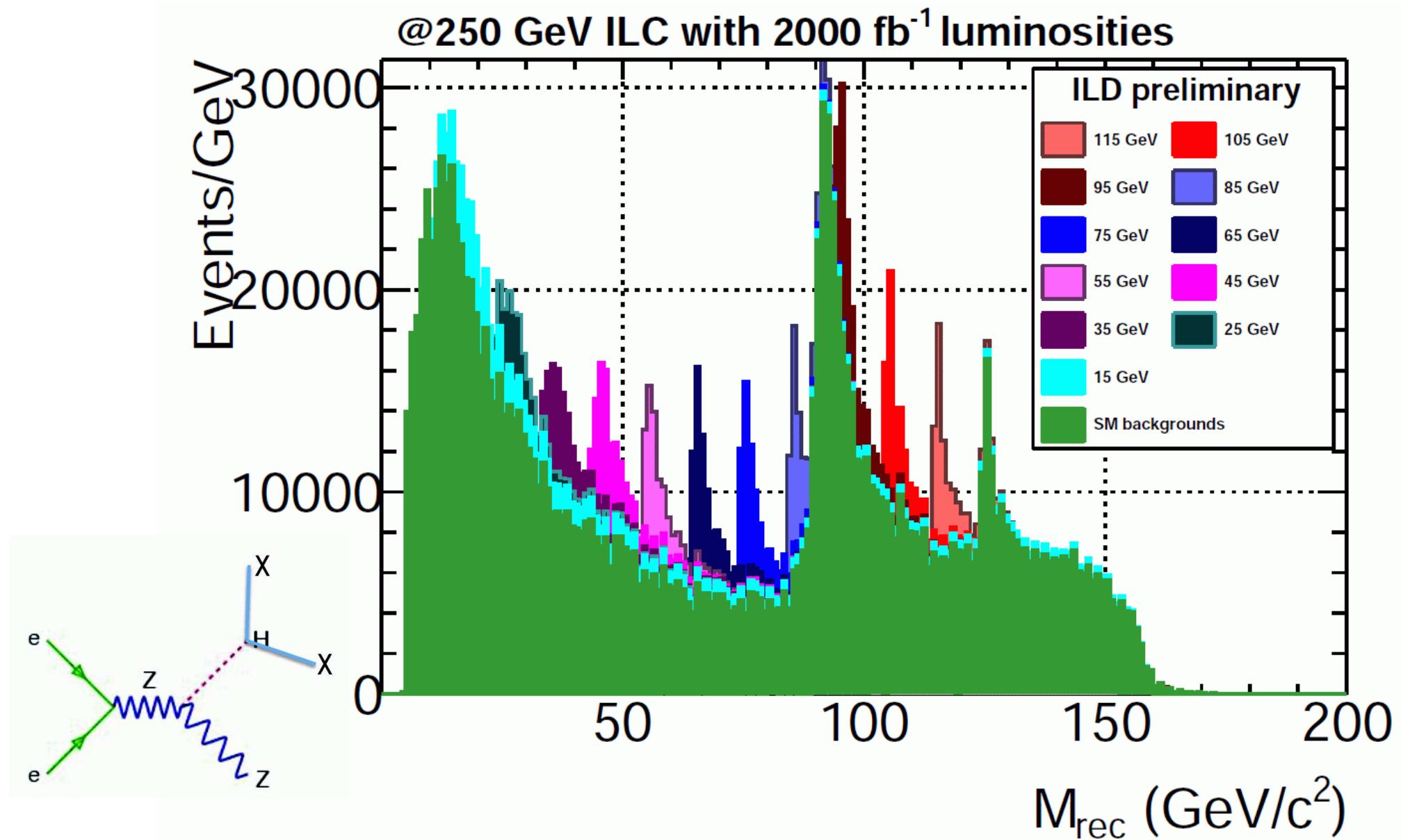
4. Physics opportunities at future e^+e^- colliders

Example for discovery potential for new light states:
Sensitivity at 250 GeV with 500 fb^{-1} to a new light Higgs

[P. Drechsel et al. '17]

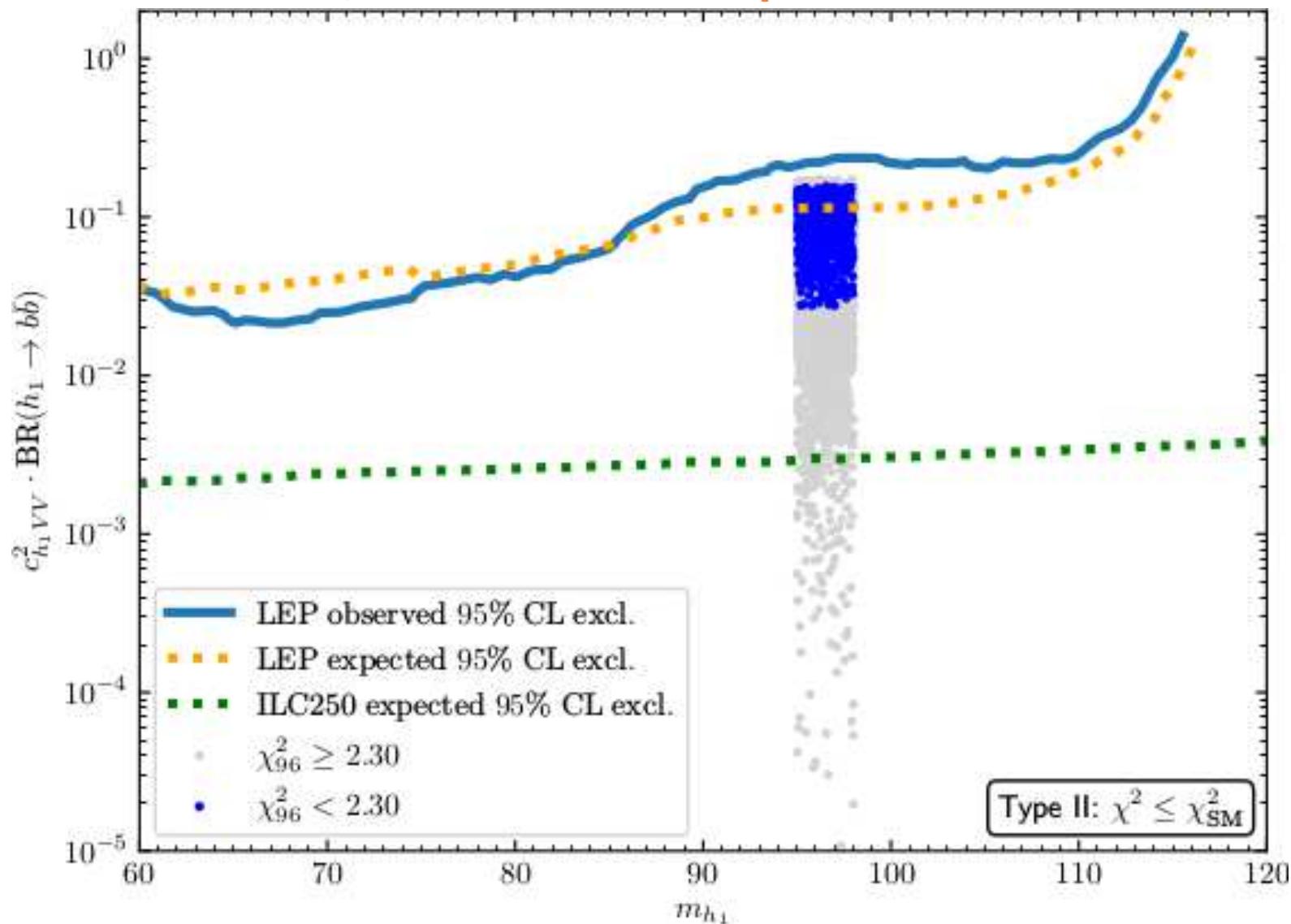


[Taken from G. Weiglein '18]



Production of the light Higgs at the ILC

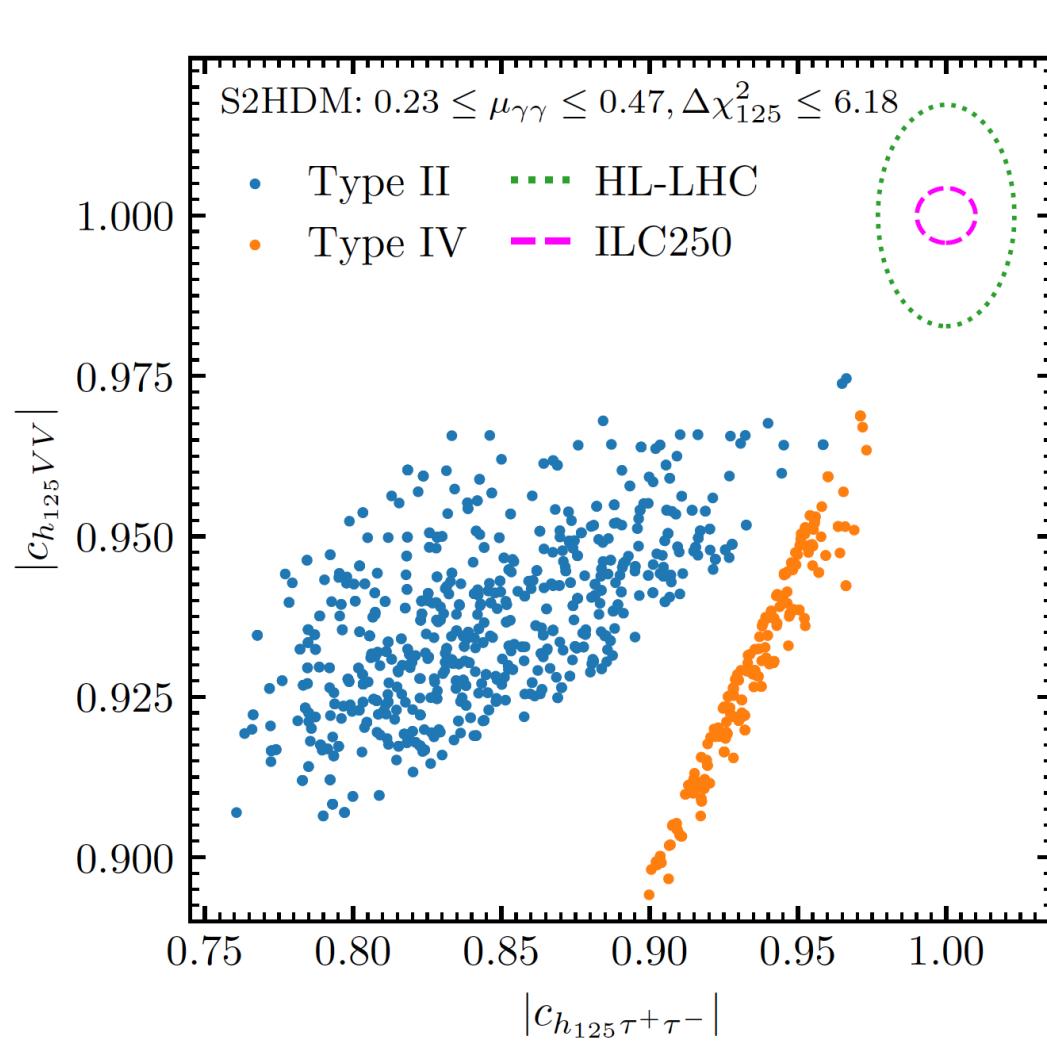
[T. Biekötter, S.H., G. Weiglein – PRELIMINARY]



⇒ new state easily in the reach of the ILC250 ⇒ coupling measurements

h_{125} coupling measurements at the HL-LHC/ILC

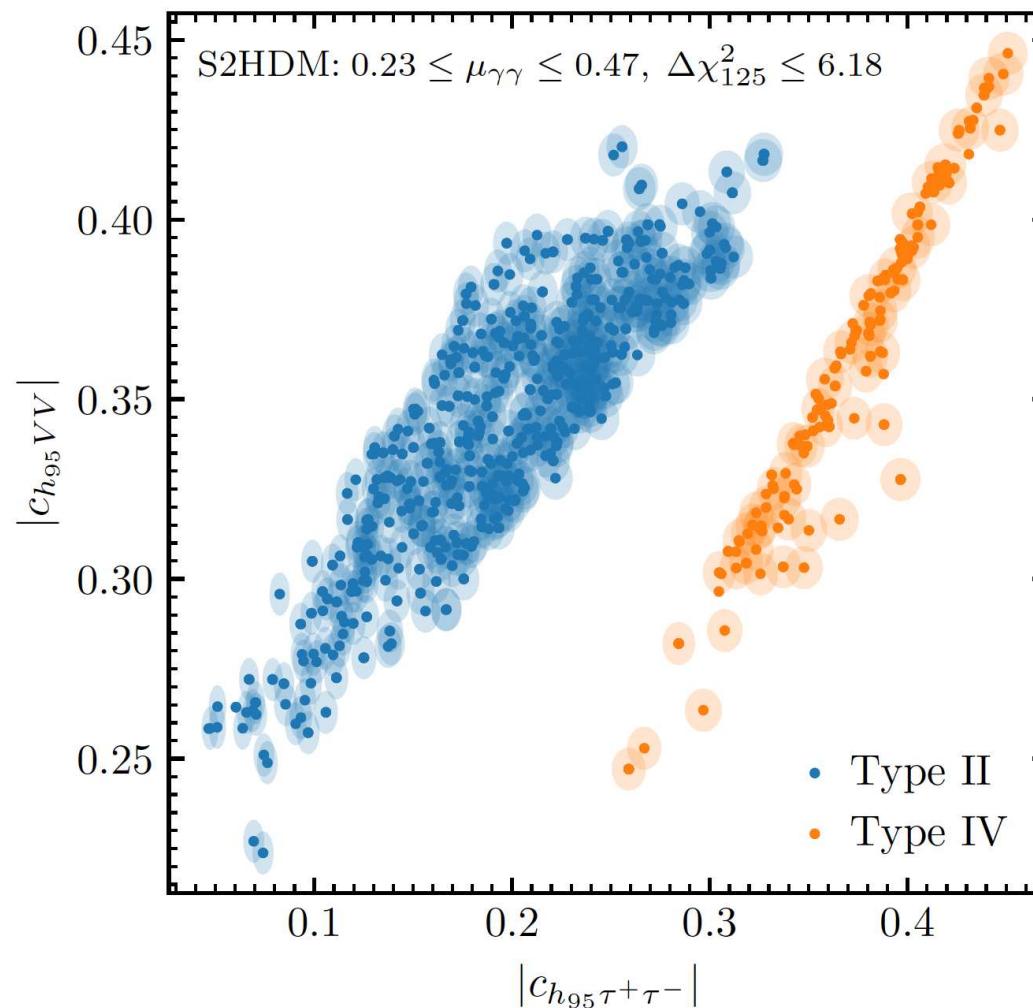
[T. Biekötter, S.H., G. Weiglein '23]



⇒ both types show some deviation from SM

h_{95} coupling measurements at the HL-LHC/ILC

[T. Biekötter, S.H., G. Weiglein '23]



⇒ models clearly distinguishable!

5. Conclusions

- Evidence for a Higgs boson at ~ 95.4 GeV
 - $pp \rightarrow h_{95} \rightarrow \gamma\gamma \Rightarrow \text{CMS: } 2.9\sigma, \text{ATLAS: } 1.7\sigma$
 - $e^+e^- \rightarrow Zh_{95} \rightarrow Zb\bar{b} \Rightarrow \text{LEP: } 2\sigma$
 - $pp \rightarrow h_{95} \rightarrow \tau\tau \Rightarrow \text{CMS: } 2.4\sigma$

\Rightarrow no LEE (as theorist I am allowed to add naively)

$\Rightarrow \sim 4.6\sigma$
- Possible model interpretation:
N2HDM or S2HDM: two Higgs doublets plus a real or complex singlet
Type II and IV: c_{h_1bb} and c_{h_1tt} independent
 \Rightarrow only type II and IV can fit the $\gamma\gamma$ and $b\bar{b}$ excesses
 $\Rightarrow \tau\tau$ excess may decide between type II and IV
- SUSY realizations:
NMSSM, $\mu\nu$ SSM, ... can describe $\gamma\gamma$ and $b\bar{b}$ using the new $\mu_{\gamma\gamma}$
- Physics opportunities at ILC250:
 - h_{125} couplings: both types clearly distinguishable from the SM
 - h_{95} can be produced in large numbers at CEPC
 - \Rightarrow high precision coupling measurements
 - \Rightarrow information about the underlying model



Further Questions?