

# First principles determination of bubble wall velocity and LTE approximation.

**Benoit Laurent**

Based on

J. Cline and BL (2204.13120)

W.-Y. Ai, BL and J. van de Vis (2303.10171)

SPCS 2023: Phase Transitions, Gravitational Waves, and Colliders

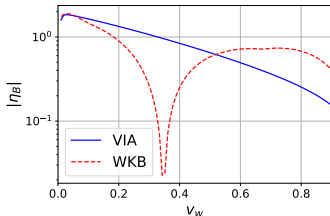
Tsung Dao Lee Institute, Shanghai

September 23, 2023

# Why is $\xi_w$ important?

## Baryogenesis:

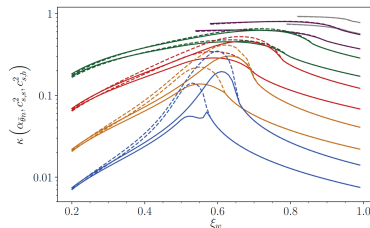
- FOPTs offer a mechanism to explain the observed baryon asymmetry of the Universe (e.g. electroweak baryogenesis),
- Baryogenesis requires a departure from equilibrium, which can be provided by the bubble wall.



Cline and Laurent (2108.04249)

## Gravitational waves:

- In the main mechanism, the energy in the wall is converted into sound waves which are then dissipated as GWs,
- These GWs could be detected by future space-based experiments such as LISA.



Giese, Konstantin, Schmitz and van de Vis  
(2010.09744)

- ▶ Wall velocity: General framework
- ▶ Total pressure and application to xSM
- ▶ Local thermal equilibrium approximation

# Wall velocity: General framework

A FOPT is triggered by the variation of a scalar field  $\phi$ 's VEV:

$$\phi(x^\mu) : \phi_1 \longrightarrow \phi_2.$$

This can be described by the **EOM**:

$$\partial^2 \phi + \frac{\partial V_0(\phi)}{\partial \phi} + \sum_i \frac{\partial m_i^2(\phi)}{\partial \phi} \int \frac{d^3 p}{(2\pi)^3 2E_i} f_i(p, x) = 0.$$

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The set of  $f_i$  is the distribution function of each species in the plasma. They follow the **Boltzmann equations**:

$$\left( p^\mu \partial_\mu + \frac{1}{2} \partial^\mu (m_i^2) \partial_{p_\mu} \right) f_i(p, x) = -\mathcal{C}_i[f].$$

To simplify the analysis,  $f_i$  is split into a “background” part (common to all species) and an “out-of-equilibrium” part:

$$f_i = f_{bg} + \delta f_i,$$

with

$$f_{bg} \equiv f_{eq} = \left[ \exp \left( \frac{1}{T(\mathbf{x})} p_\mu u_{pl}^\mu(\mathbf{x}) \right) \pm 1 \right]^{-1}.$$

$\delta f_i$  is assumed to be nonzero only for a few species that undergo a large mass variation (e.g. top quark).

The evolution of  $T(\mathbf{x})$  and  $v_{\text{pl}}(\mathbf{x})$  is derived from the conservation of the energy-momentum tensor

$$0 = \partial_\mu T^{\mu\nu} = \partial_\mu (T_\phi^{\mu\nu} + T_{\text{plasma}}^{\mu\nu}),$$

where

$$T_\phi^{\mu\nu} = \partial^\mu \phi \partial^\nu \phi - \eta^{\mu\nu} \left[ \frac{1}{2} \partial_\alpha \phi \partial^\alpha \phi - V_0(\phi) \right],$$
$$T_{\text{plasma}}^{\mu\nu} = \omega u_{\text{pl}}^\mu u_{\text{pl}}^\nu - p_T \eta^{\mu\nu} + \sum_i \int \frac{d^3 p}{(2\pi)^3 E_i} p^\mu p^\nu \delta f_i(\mathbf{x}, \mathbf{p}),$$



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In the wall frame, these equations can be integrated directly:

$$T^{30} = c_1(\xi_w) \quad \text{and} \quad T^{33} = c_2(\xi_w),$$

with  $c_{1,2}$  independent of  $z$ .

These equations can be solved for  $T$  and  $v_{\text{pl}}$ .

(More details in Cline and Laurent 2204.13120)

## 3 types of solutions

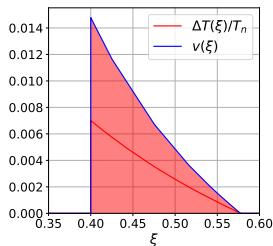
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## Deflagration:

- Subsonic:  $\xi_w < c_s$
- $v_- = \xi_w$  and  $T_+^{\text{SW}} = T_n$
- Shock wave



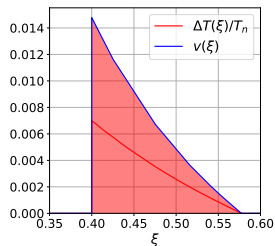
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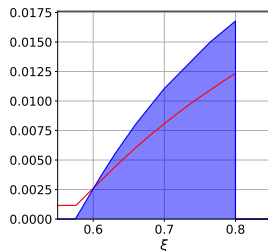
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## Detonation:

- Supersonic:  $\xi_w \geq \xi_J > c_s$
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- Rarefaction wave



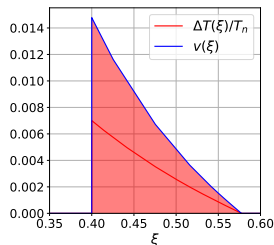
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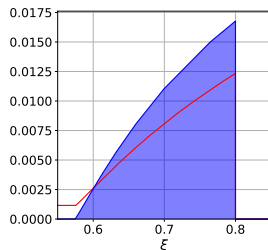
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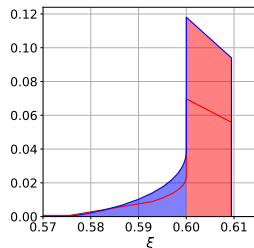
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## Hybrid:

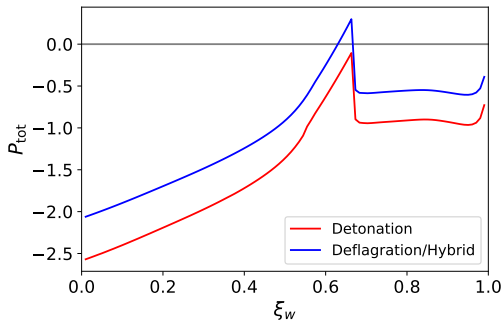
- $c_s < \xi_w < \xi_J$
- $v_- = c_s$  and  $T_+^{\text{sw}} = T_n$
- Rarefaction and shock waves



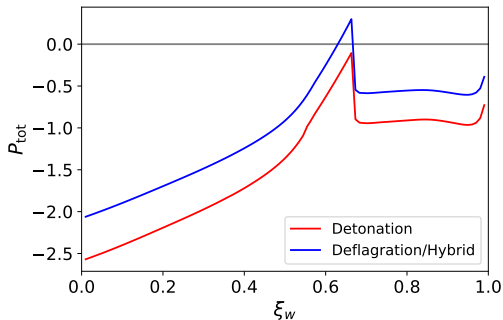
$$\xi_w = 0.6$$

# Total pressure and application to xSM

# Pressure on the wall

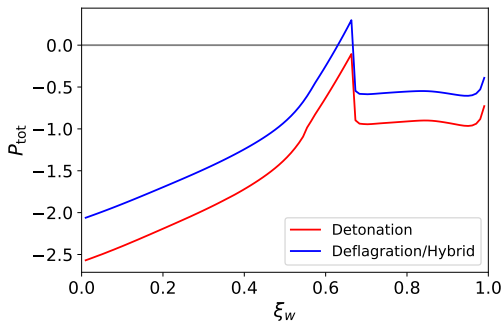


# Pressure on the wall

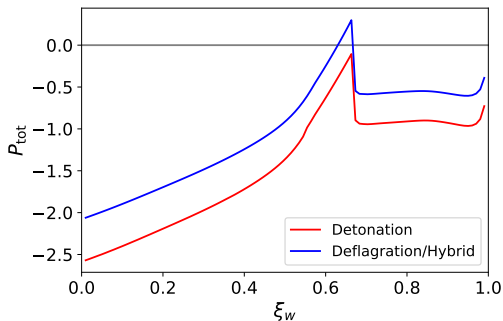


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- The pressure is maximized at  $\xi_J$ , so if there is no deflagration/hybrid solution, the wall runs away (or is stopped at  $\gamma_w \gg 1$ ),
- Solutions can be classified in **2 categories**: deflagrations/hybrids ( $\xi_w \sim c_s$ ) and ultrarelativistic detonations ( $\gamma_w \gg 1$ ).

## Example: Singlet scalar extension

We consider the singlet scalar extension with a  $Z_2$  symmetry:

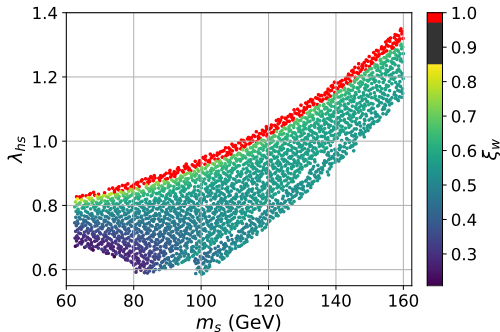
$$V_0(h, s) = \frac{\mu_h^2}{2} h^2 + \frac{\lambda_h}{4} h^4 + \frac{\mu_s^2}{2} s^2 + \frac{\lambda_s}{4} s^4 + \frac{\lambda_{hs}}{4} h^2 s^2.$$

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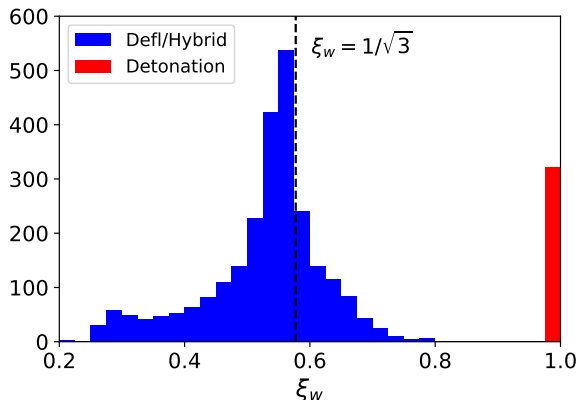
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Scan with  $\lambda_s = 1$ :



# Velocity distribution



- No models below 0.2,
- No models with  $\xi_J < \xi_w \ll 1$ ,
- $\xi_w > 0.5$  for 75% of the models.

# Local thermal equilibrium approximation

Far away from the wall, the fluid equations become:

$$\Delta (\omega \gamma^2 v) = \Delta (\omega \gamma^2 v^2 + p) = 0.$$

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Now, let's assume  $\delta f_i = 0$ :

We can contract the divergence of  $T^{\mu\nu}$  with  $u^\nu$  to get

$$0 = u_\nu \partial_\mu T^{\mu\nu} = u^\nu \partial_\nu \phi \underbrace{\left[ \partial^2 \phi + \frac{\partial V_T}{\partial \phi} \right]}_{\text{EOM}=0} + T \partial_\nu S^\nu,$$

with the entropy flux:

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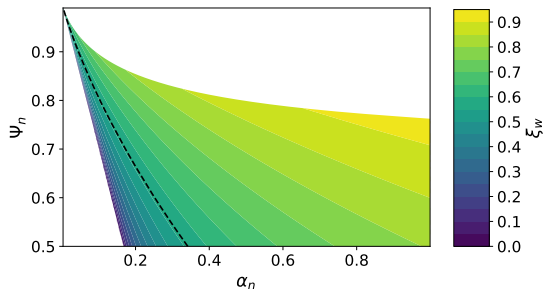
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**This new matching equation replaces the EOM!**

By fitting to the template model (constant sound speed),  $\xi_w$  only depends on 4 quantities:

$$\alpha_n \text{ (strength of PT)}, \quad \Psi_n = \frac{\omega_-}{\omega_+}, \quad \nu = 1 + 1/c_{s,-}^2, \quad \mu = 1 + 1/c_{s,+}^2.$$

See Ai, Laurent and van de Vis (2303.10171) for a Code snippet to compute  $\xi_w$ .



- The wall velocity can be calculated from a set of fluid equations (EOM, Boltzmann equation and conservation of EMT),
- These equations can be significantly simplified with the LTE approximation,
- Using the template model, the fluid equations can be written in a model-independent way which only depends on 4 free parameters,
- Solutions can be classified as deflagrations/hybrids ( $\xi_w \sim c_s$ ) or ultrarelativistic detonations ( $\gamma_w \gg 1$ ),
- For  $\Psi_n \lesssim 0.75$ , hybrid walls can be arbitrarily strong, so baryogenesis and GW production could be highly efficient.

Thank you!