

Could the neutrino emission of TXS 0506+056 come from the core of the active galactic nuclei?

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OUTLINES

01. Introduction

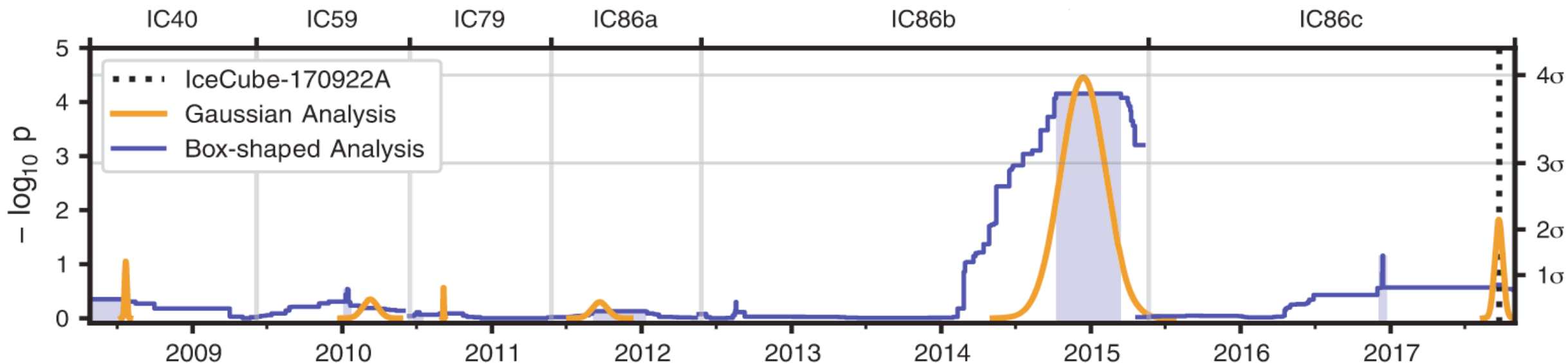
02. Steady-state neutrino emission

03. 2014-2015 flare neutrino emission

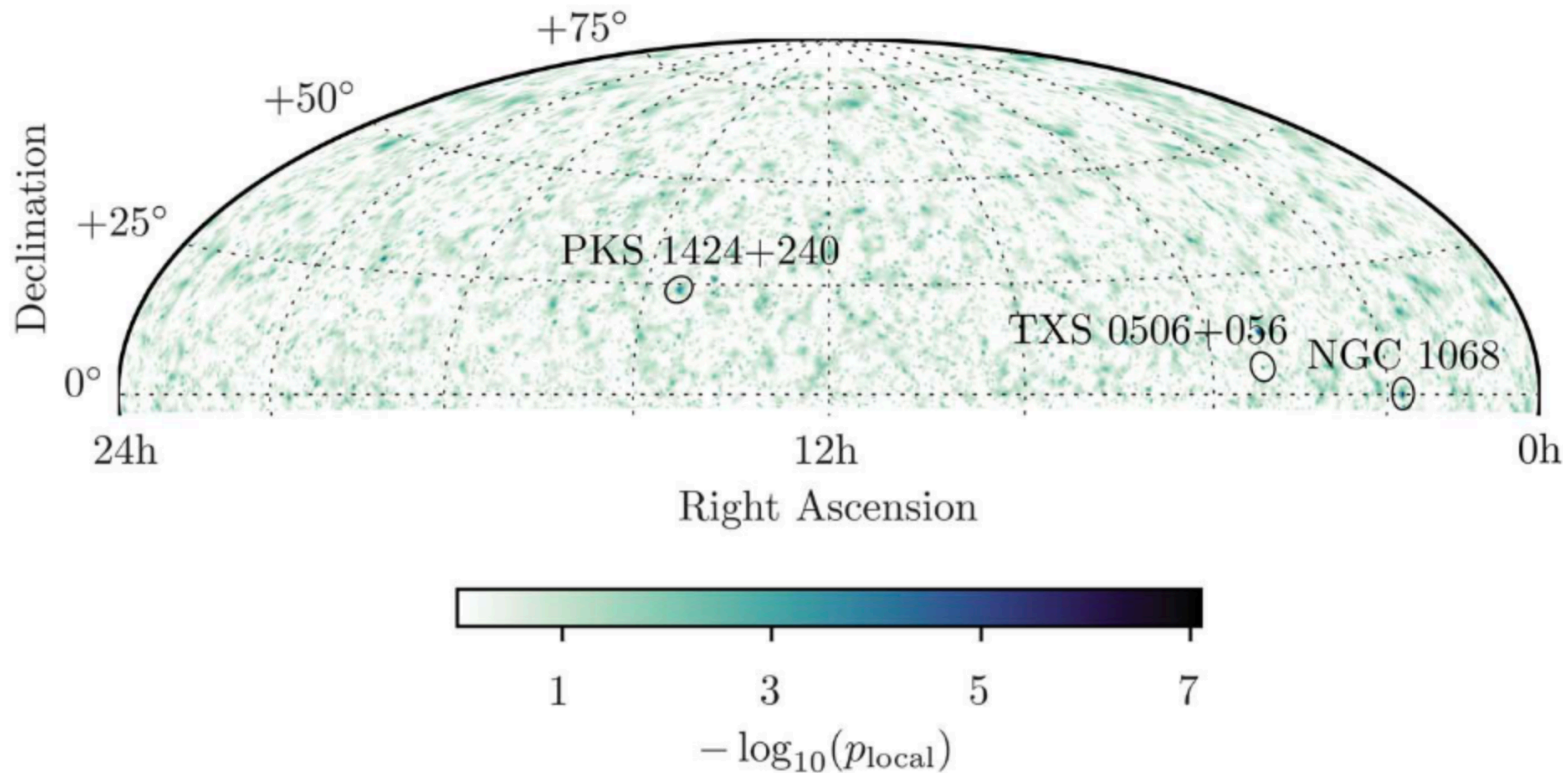
04. Summary



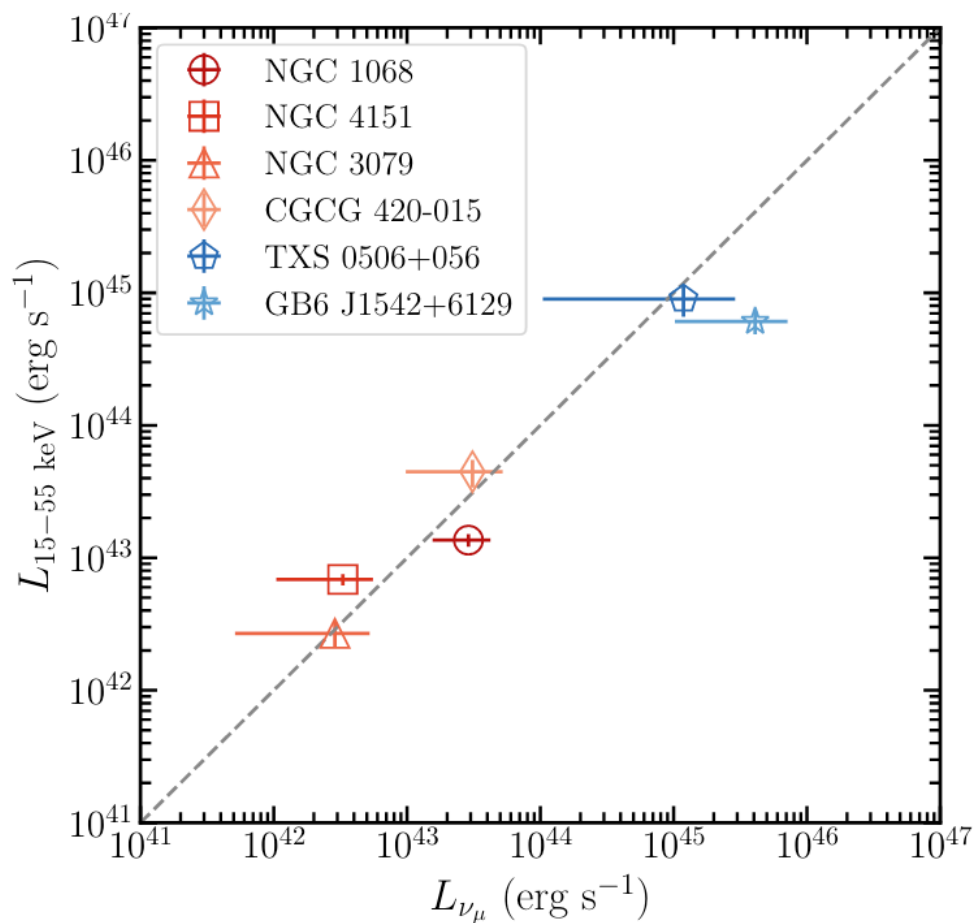
- Blazar, $z = 0.3365 \pm 0.0010$
- A high-energy (290 TeV) muon neutrino event, IceCube170922A, was detected in both spatial and temporal coincidence with the γ -ray flare of the known blazar TXS 0506+056.
- 3.5σ excess of 13 ± 5 high-energy neutrino events was discovered in the period between 2014 September and 2015 March, known as 2014-2015 neutrino flare.



IceCube Collaboration. 2022



- Sky map of the scan for point sources in the Northern Hemisphere. TXS 0506+056 relates to a time-integrated signal over the duration of 9.5yr, whereas previous analyses have found evidence for transient emission
- Local significances of 3.5σ ,

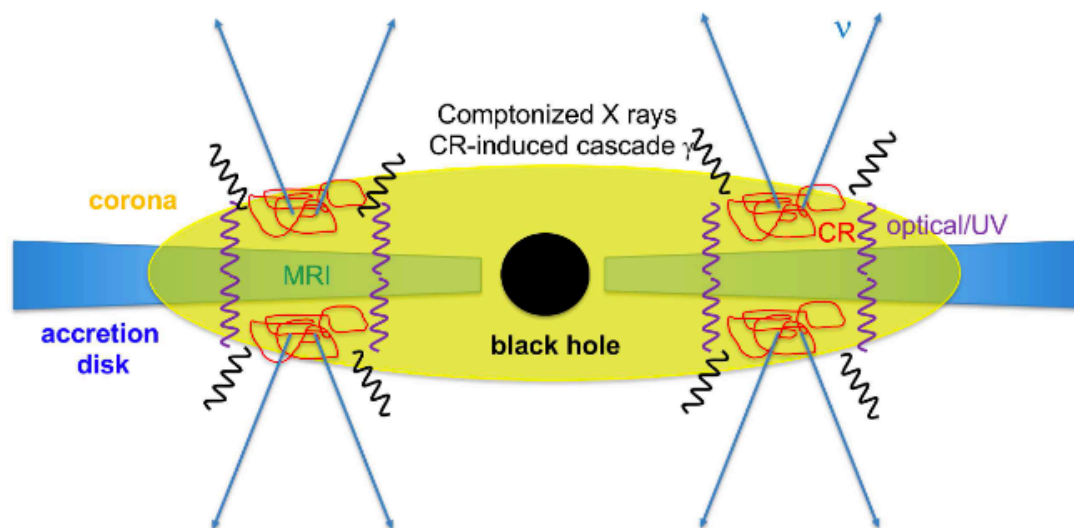


Kun et al. 2024

- Neutrino emission from radio-quiet Seyfert galaxies is unlikely to be related to their weak jets (NGC 1068)
- Neutrino flux is about 5 times higher than the average γ -ray flux for 14-15 flare
- A hint of gamma-absorption is also observed in the diffuse neutrino flux, pointing to the so-called "hidden" neutrino sources.
- **Core region can be a common origin, implying non-jet origin**

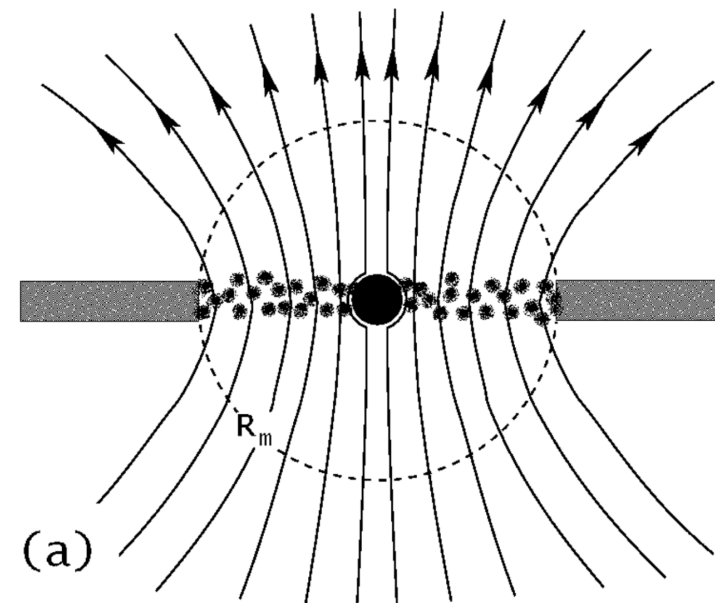
“Core Region” Disk+Corona

Murase et al. 2020

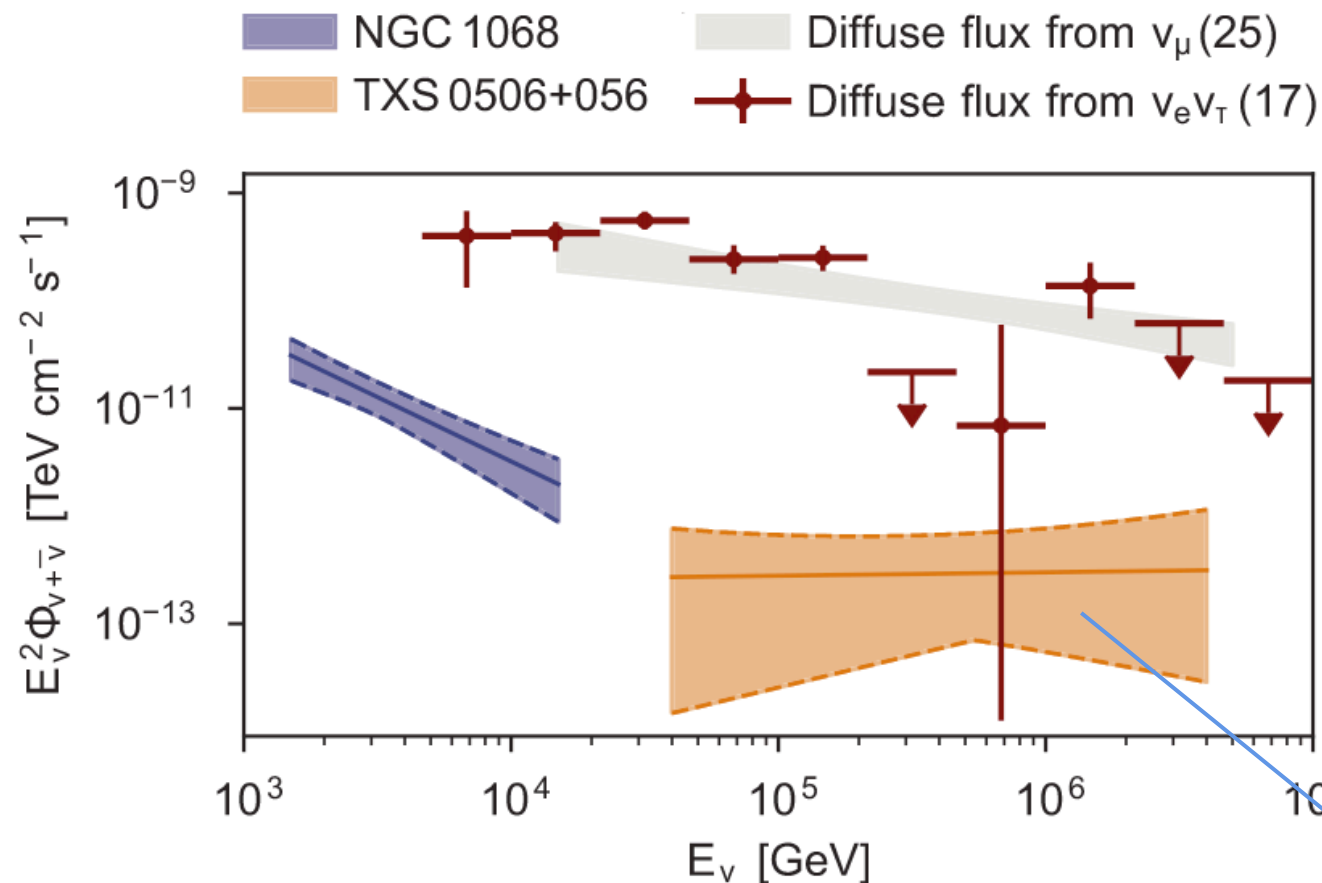


MAD (Magnetically arrested disk)

Narayan et al. 2003



- Photon Field: “Big Blue Bump”+ power-law X-ray component
- CR can be accelerated in the MAD (Magnetically Arrested Disk, $\beta \sim 1$)
- Emission region set within $10 R_g$



Energy Range: 40 TeV - 4 PeV

Luminosity upper limit : $1.3 \times 10^{45} \text{ erg s}^{-1}$

Luminosity lower limit: $7 \times 10^{43} \text{ erg s}^{-1}$

Central Value: $5.6 \times 10^{44} \text{ erg s}^{-1}$

Index: $\Gamma \approx 2.04$

Steady-state Emission

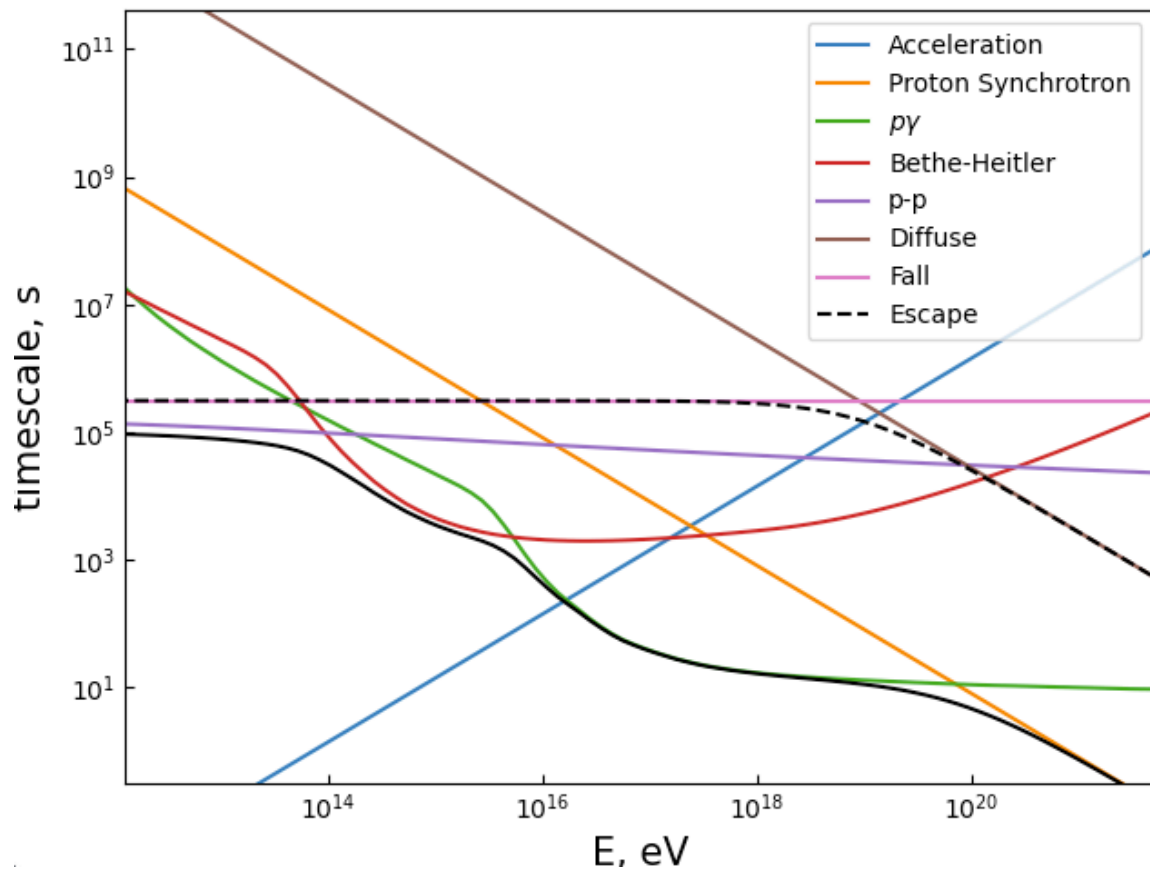
- For high-energy protons accelerated within MAD, we can estimate the CR luminosity from Eddington luminosity.

$$L_{CR} = \epsilon_{CR} \dot{m} L_{0506, \text{Edd}} = 6 \times 10^{45} \text{erg s}^{-1} \left(\frac{\epsilon_{CR}}{0.3} \right) \left(\frac{\dot{m}}{0.5} \right) \left(\frac{M_{BH}}{3 \times 10^8 M_{\odot}} \right)$$

$$L_{0506, \text{Edd}} \approx 4 \times 10^{46} \text{erg s}^{-1} \left(\frac{M_{BH}}{3 \times 10^8 M_{\odot}} \right)$$

- High-energy neutrinos are produced via pp and $p\gamma$ process. The estimated neutrino luminosity is illustrated as

$$L_{\nu} = \frac{1}{8} f_{pp, p\gamma} \epsilon \dot{m} L_{0506, \text{Edd}}$$



- We do not focus on specific acceleration mechanisms; instead, different acceleration channels are represented by the acceleration factor η

$$t_{acc} \approx \eta r_L / c \quad \eta \sim 300$$

- The maximum energy of accelerated protons is given by

$$E_{p,max} \approx 10 \text{ PeV} \left(\frac{\dot{m}}{0.5} \right)^{1/2} \left(\frac{R}{10R_g} \right)^{-5/4} \left(\frac{M_{BH}}{3 \times 10^8 M_\odot} \right)^{5/4} \left(\frac{\beta}{1} \right)^{-1/2} \left(\frac{\eta}{300} \right)^{-1} \left(\frac{n_\gamma}{4 \times 10^{15} \text{ cm}^{-3}} \right)^{-1}.$$

- For the neutrino energy range that we are interested (40 TeV - 4 PeV) , the interaction efficiencies are numerically obtained by $f_{p\gamma} \sim 0.2$ $f_{pp} \sim 0.07$
- Neutrino luminosity via *pγ*

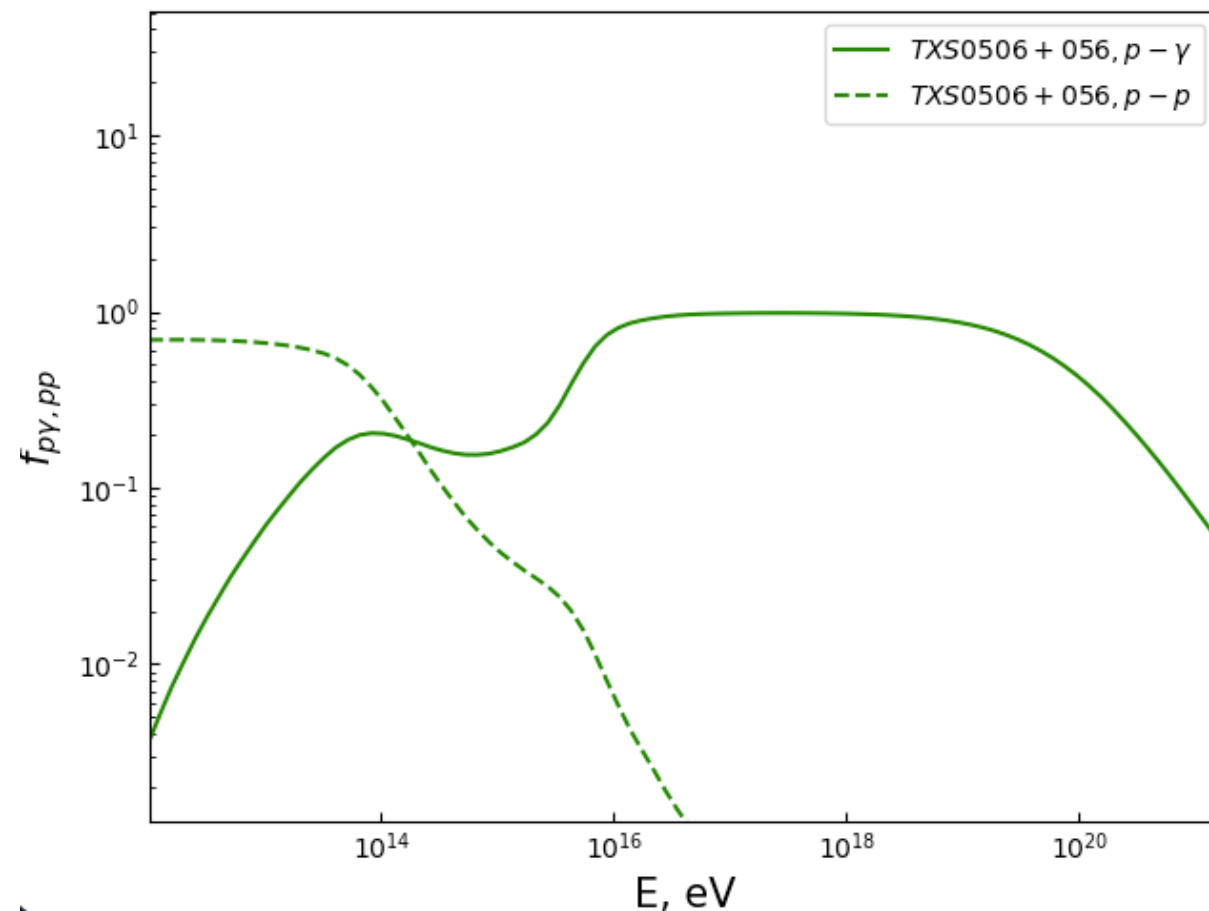
$$L_{\nu_{\mu}, p\gamma} = \frac{1}{8} \epsilon_{\text{CR}} f_{p\gamma} \dot{m} L_{\text{Edd}}$$

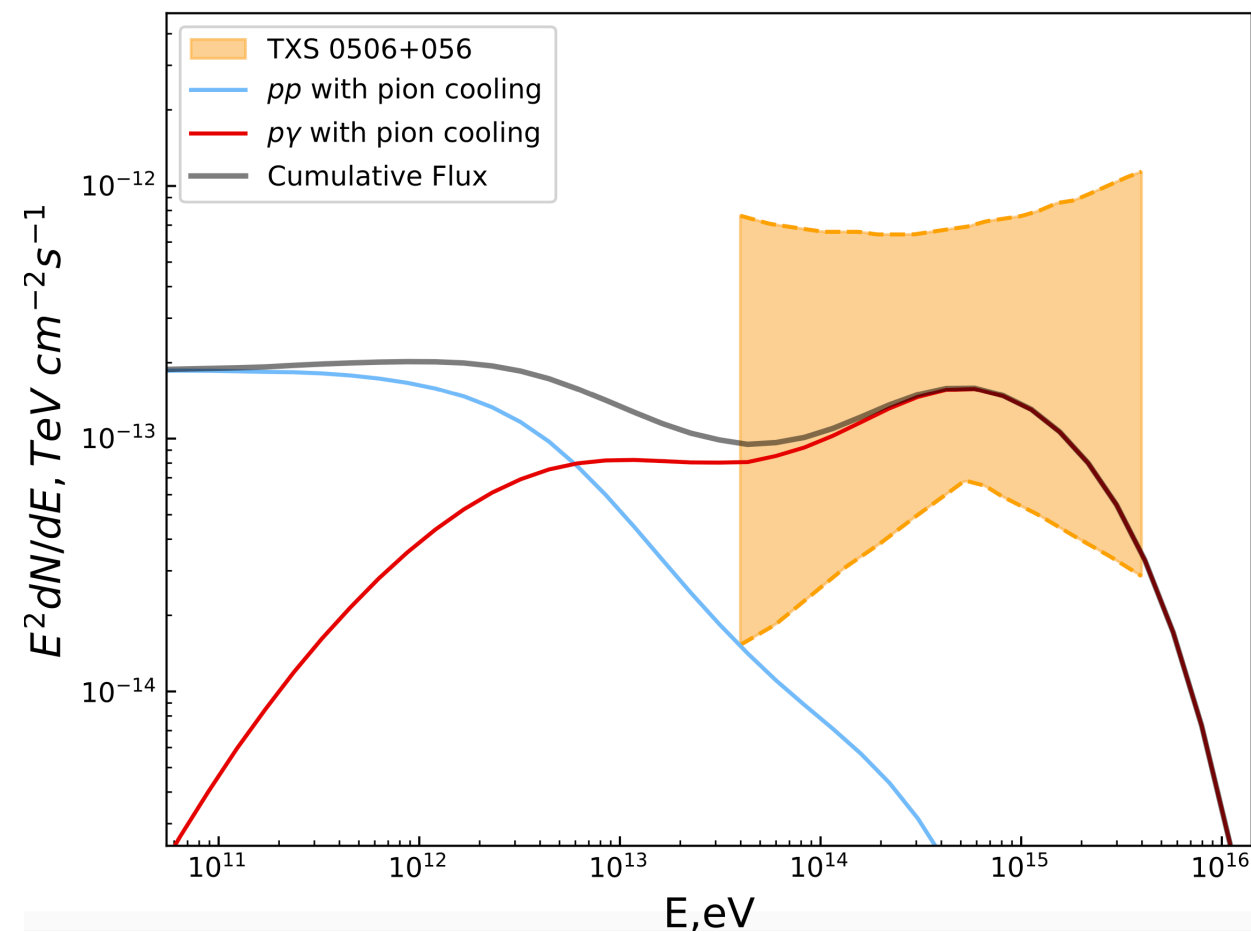
$$\sim 1.7 \times 10^{44} \text{ erg s}^{-1} \left(\frac{\epsilon_{\text{CR}}}{0.33} \right) \left(\frac{f_{p\gamma}}{0.2} \right) \left(\frac{\dot{m}}{0.5} \right)$$

- Neutrino luminosity via *pp*

$$L_{\nu_{\mu}, pp} = \frac{1}{8} \epsilon_{\text{CR}} f_{pp} \dot{m} L_{\text{Edd}}$$

$$\sim 5 \times 10^{43} \text{ erg s}^{-1} \left(\frac{\epsilon_{\text{CR}}}{0.33} \right) \left(\frac{f_{pp}}{0.07} \right) \left(\frac{\dot{m}}{0.5} \right)$$



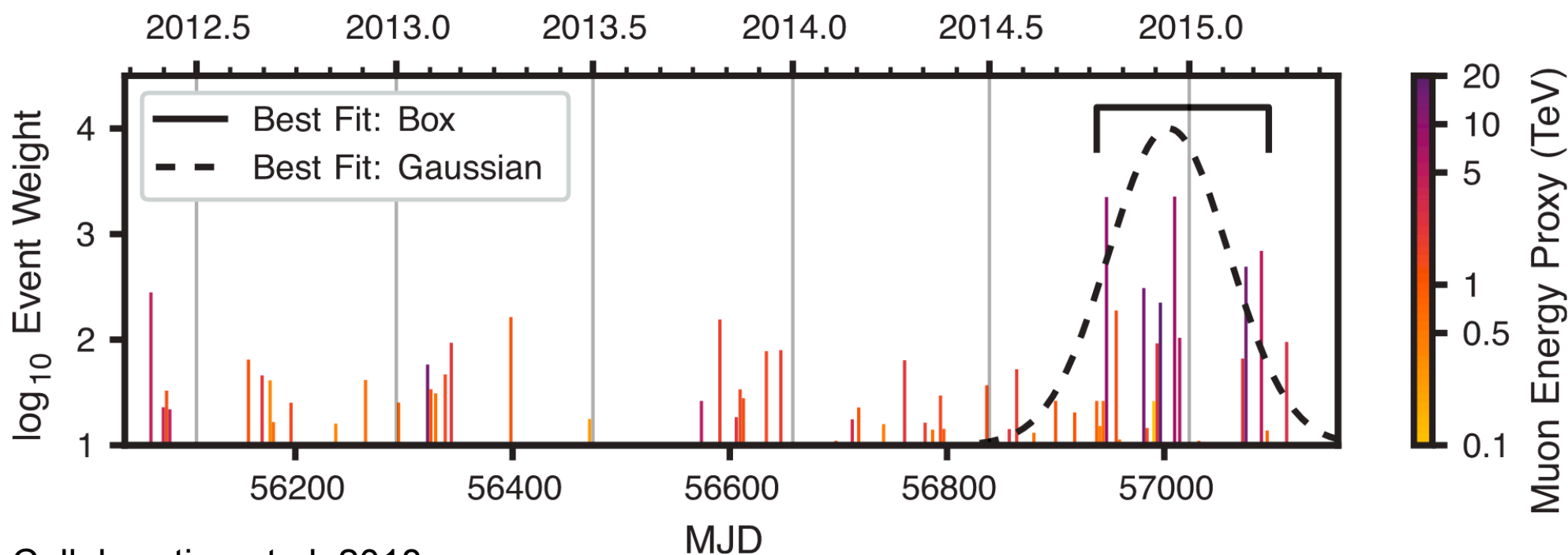


- With pion's cooling considered, the neutrinos' energy spectrum forms a natural cutoff
- $p\gamma$ dominates the steady-state neutrino emission.

$$\xi_{\pi} = \text{Min}\{t_{\pi,loss}/\tau_{\pi}, 1\}$$

$$\sim 0.6 \left(\frac{E_{\pi}}{10 \text{ PeV}} \right)^{-2/3} \left(\frac{\beta}{1} \right)^{1/3}$$

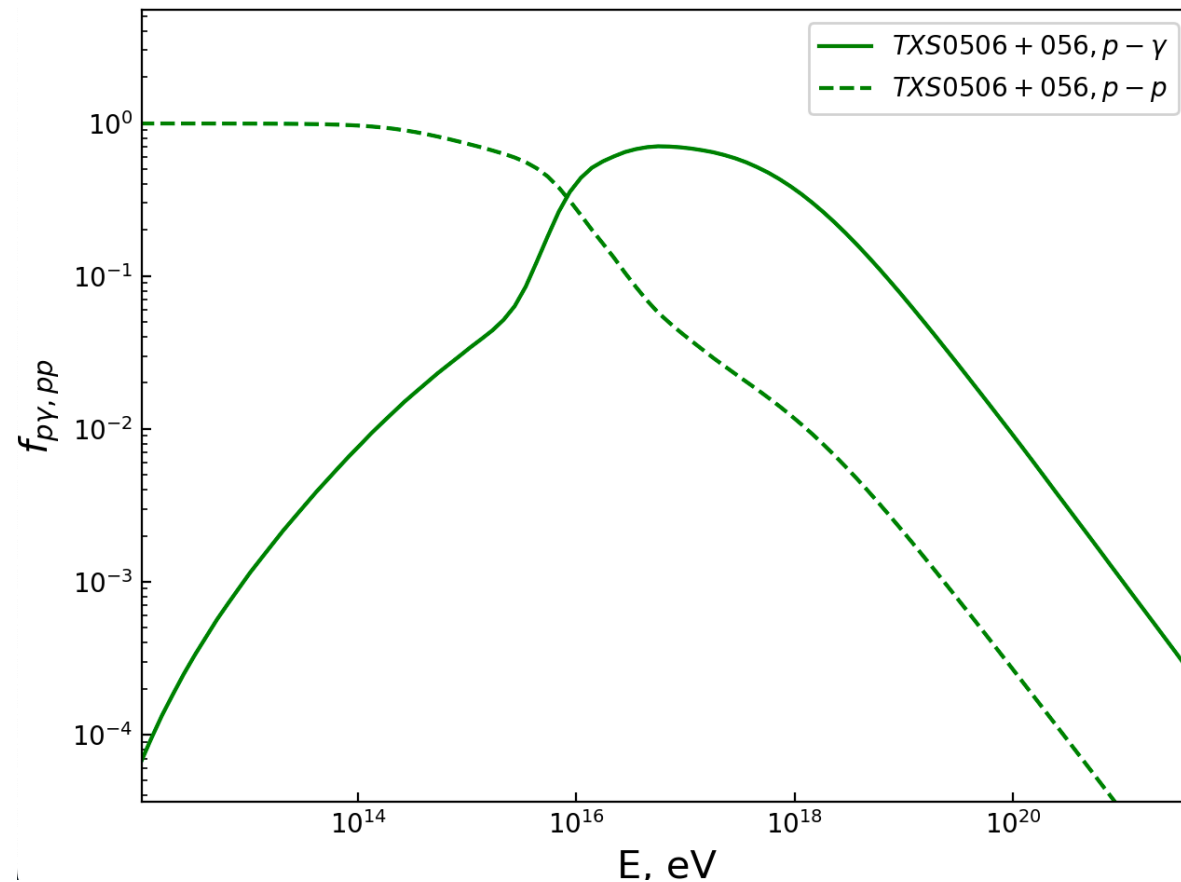
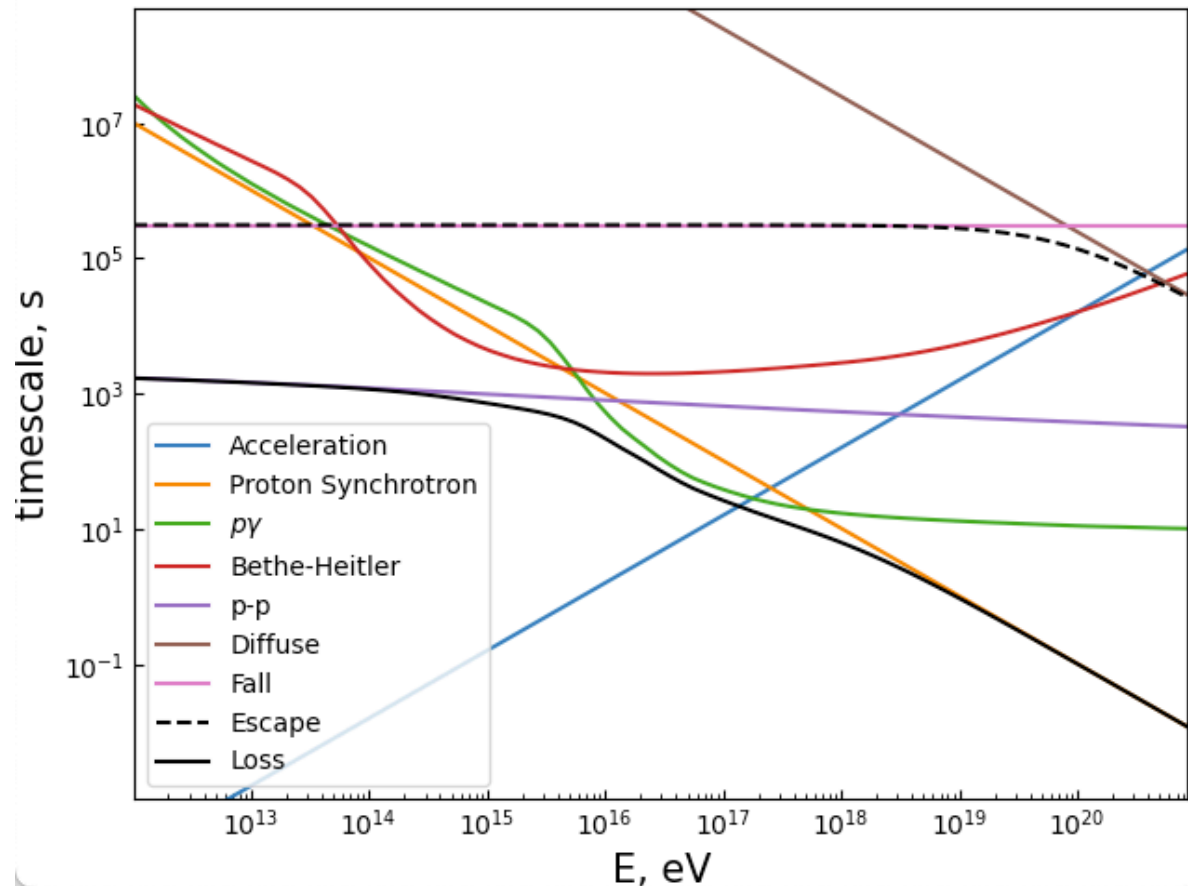
- The neutrino flare between 2014-2015 last 158 days
- Best-fitting spectral index $\Gamma = 2.1$
- Energy range: 32 TeV and 3.6 PeV.
- Isotropic neutrino luminosity $1.2^{+0.6}_{-0.4} \times 10^{47} \text{ erg s}^{-1}$
- 13 ± 5 muon-neutrino events that are present in addition to the expected background.

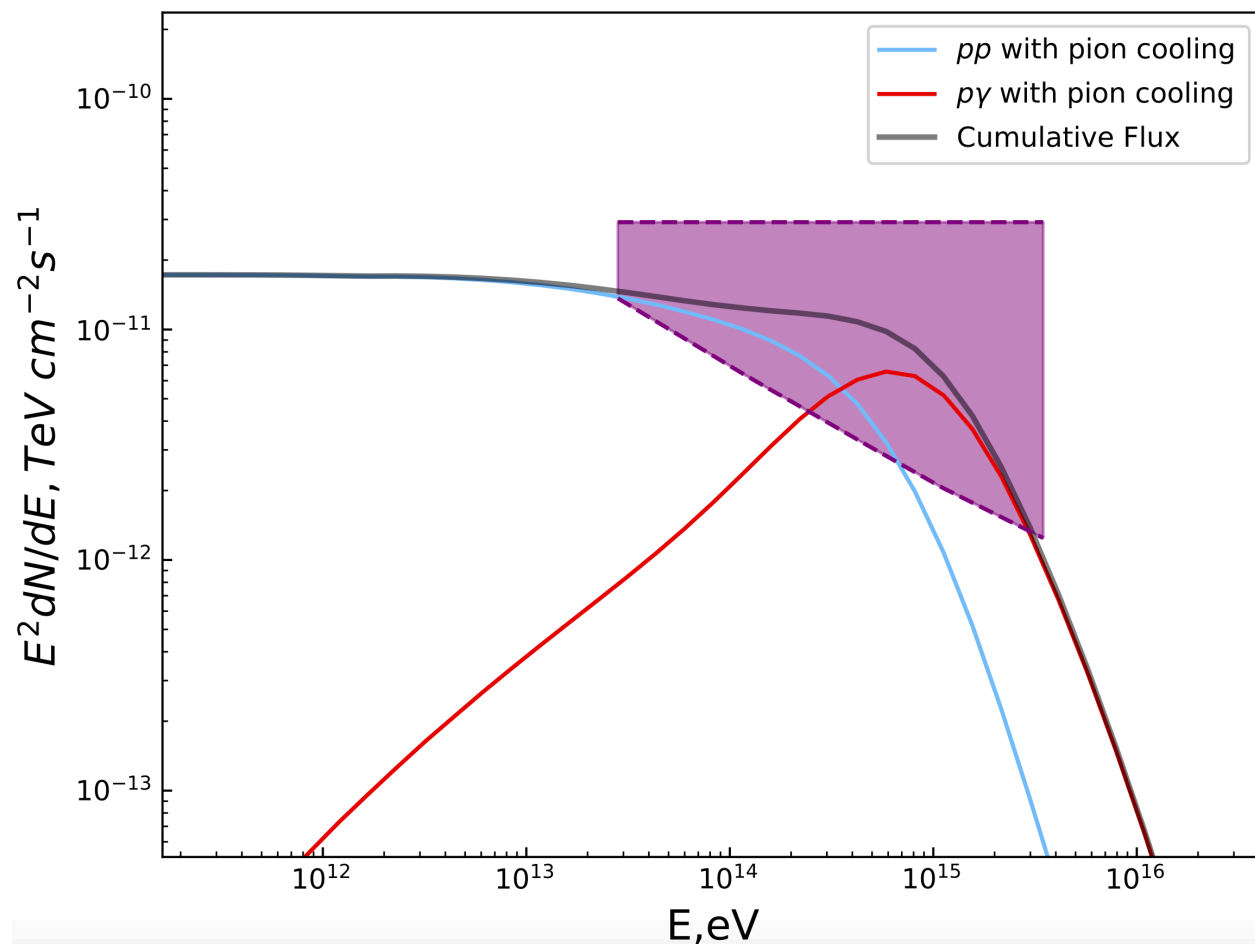


Timescale and efficiency in flare scenario

- We consider a Super-Eddington accretion with $\dot{m} \sim 40$ to explain the flare emission

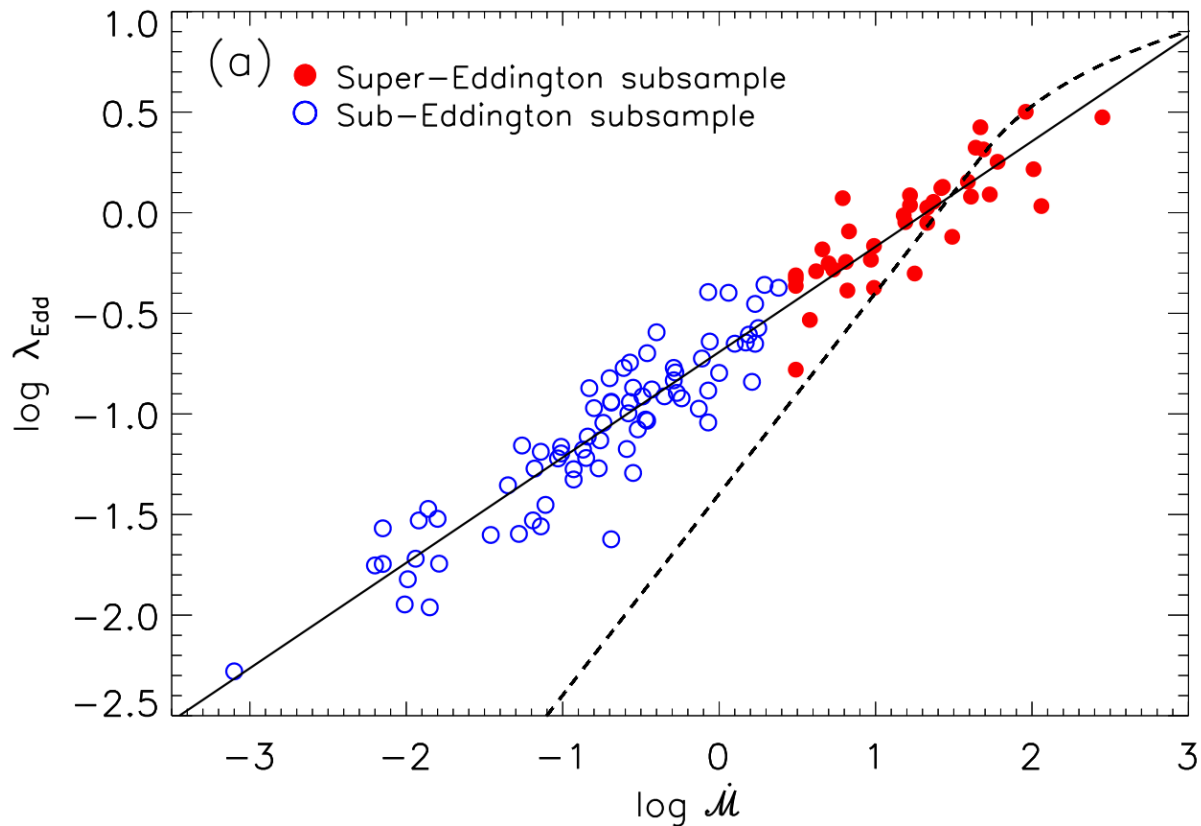
$$L_\nu = \frac{1}{8} f_{pp,p\gamma} \epsilon \dot{m} L_{0506,Edd} \sim 6 \times 10^{46} \text{ erg s}^{-1} \left(\frac{\epsilon_{CR}}{0.33} \right) \left(\frac{f}{1} \right) \left(\frac{\dot{m}}{40} \right)$$





- We consider a super-Eddington accretion $\dot{m} \sim 40$ to explain the 14-15 neutrino flare emission
- pp and $p\gamma$ both contribute the 2014-2015 flare neutrino emission.

- In the case of Sub-Eddington accretion ($\dot{m} \sim 0.5$) and MAD ($\beta \sim 1$), the steady-state neutrinos of TXS0506 may come from its core region.
- The flare neutrino emission in the period of 2014-2015 can explained by the core region emission with the Supper-Eddington accretion ($\dot{m} \sim 40$),



Huang et al. 2020

- 14-15 flare X-ray upper limit at 10 keV:
 $10^{-11} \text{ erg cm}^{-2} \text{ s}^{-1}$
- The Bolometric luminosity The evolution of X-ray and accretion rates is not linear with a power-law slope of 0.52
- The accretion rate increased by a factor of 80 from steady state to flare state; however, the X-ray luminosity likely rose by only a factor of 6-7, which would remain within the SWIFT upper limit of 2014-2015.