







Shaoming Hu

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Oct 2024 @ Shanghai

Transient Phenomena and Physical Processes Around Supermassive Black Holes



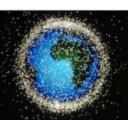
Outline

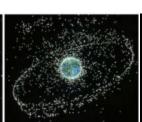




- Variability properties of blazars
- TDE in OJ 287
- Telescopes and Instruments of Shandong University





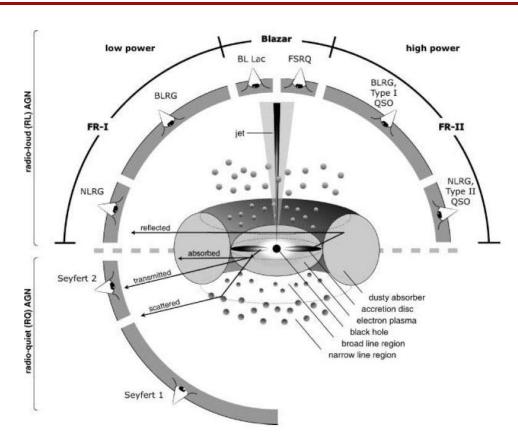




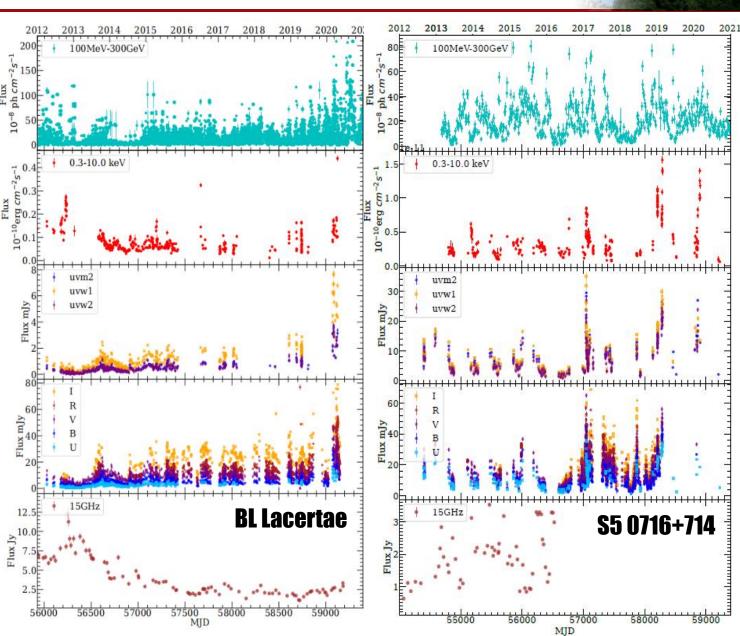


Multiband variability of Blazar





- What is physical source of the variability?
- ➤ Is there any TDE related variability?





Statistical analysis for blazar flares

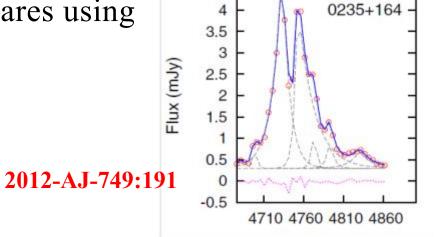
The light curves could be decomposed into individual flares using fitting methods.

$$\Delta S(t) = \begin{cases} \Delta S_{\text{max}} e^{(t - t_{\text{max}})/\tau}, & t < t_{\text{max}} \\ \Delta S_{\text{max}} e^{(t_{\text{max}} - t)/1.3\tau}, & t > t_{\text{max}} \end{cases}$$

 $= f_0 + f_{max} \times \exp[-(t - t_0)/T_d], t > t_0$

Abdo (2010)

$$F(t) = F_c + F_0 \frac{1}{e^{(\frac{t_0 - t}{T_r})} + e^{(\frac{t - t_0}{T_d})}}$$



Chatterjee (2012) $f(t) = f_0 + f_{max} \times \exp[(t - t_0)/T_r], t < t_0$

Flare duration time

$$T \simeq 2(T_r + T_d)$$

Skewness parameter

$$\xi = \frac{T_d - T_r}{T_d + T_r}$$

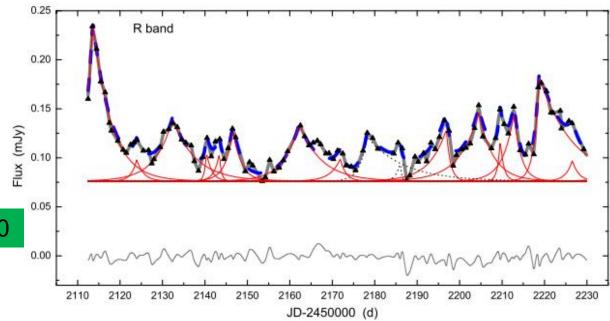
· Amplitude of flare

 f_{max}

Guo, Hu et al. 2016, MNRAS 460, 1790

Energy output

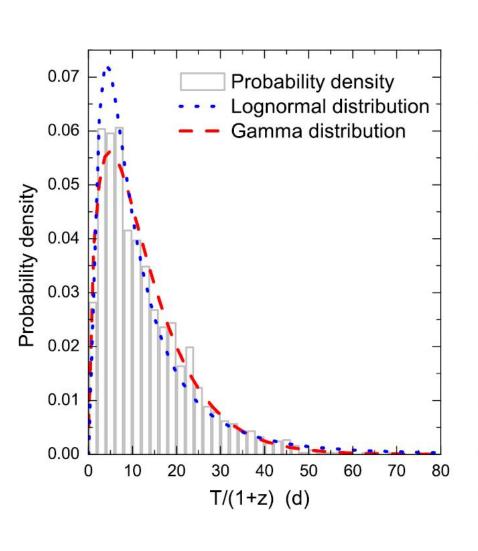
$$Area = \int_{-\infty}^{+\infty} f(t) - f_0 dt = f_{max}(t_r + t_d)$$

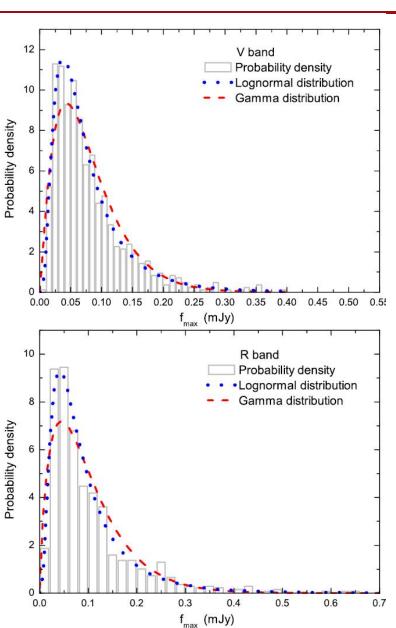


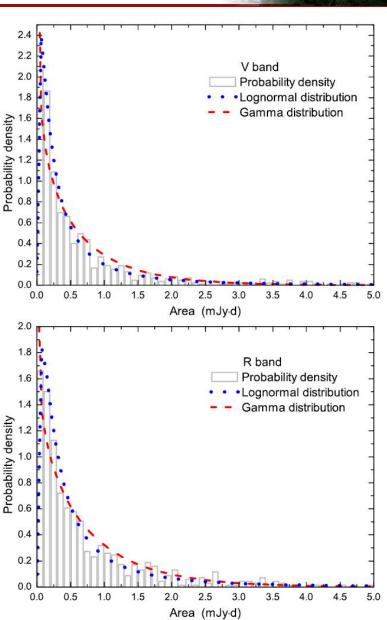


Optical Flares properties of BL Lac









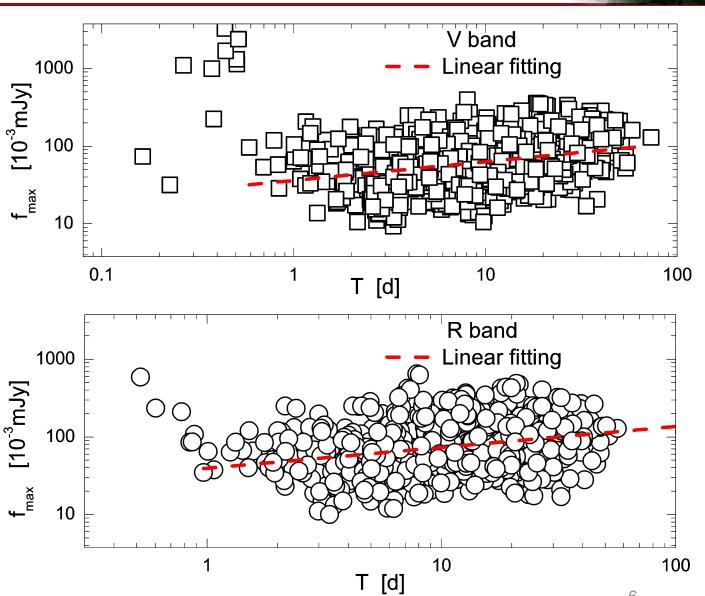
Guo, Hu et al. 2016, MNRAS 460, 1790



Statistical analysis for blazar flares



Seems there are at least two mechanisms of flares





JD-2454833 (Day)

Statistical analysis for blazar flares



2.5

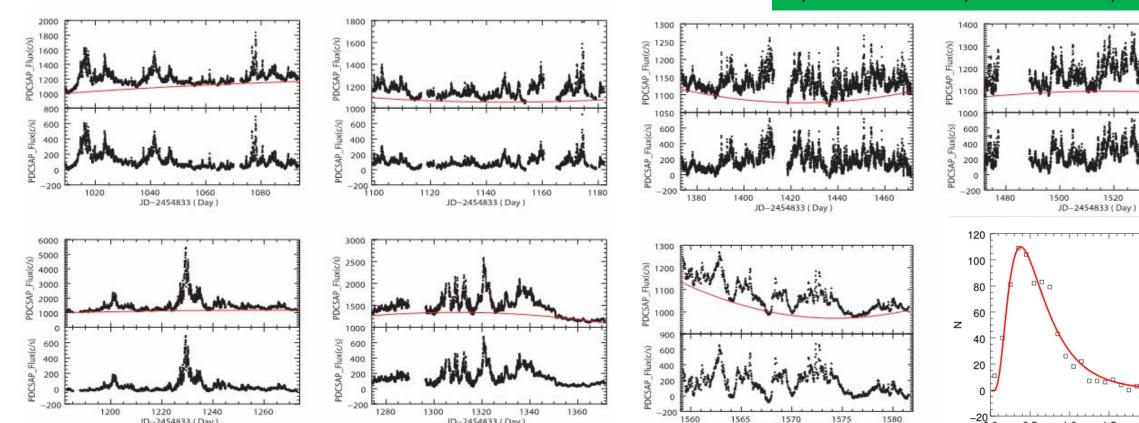
Duration(day)

Name	Q11	Q12	Q13	Q14	Q15	Q16	Q17
Beginning Date	29-Sep-2011	05-Jan-2012	29-Mar-2012	28-Jun-2012	05-Oct-2012	12-Jan-2013	09-Apr-2013
Ending Date	04-Jan-2012	28-Mar-2012	27-Jun-2012	03-Oct-2012	11-Jan-2013	08-Apr-2013	11-May-2013

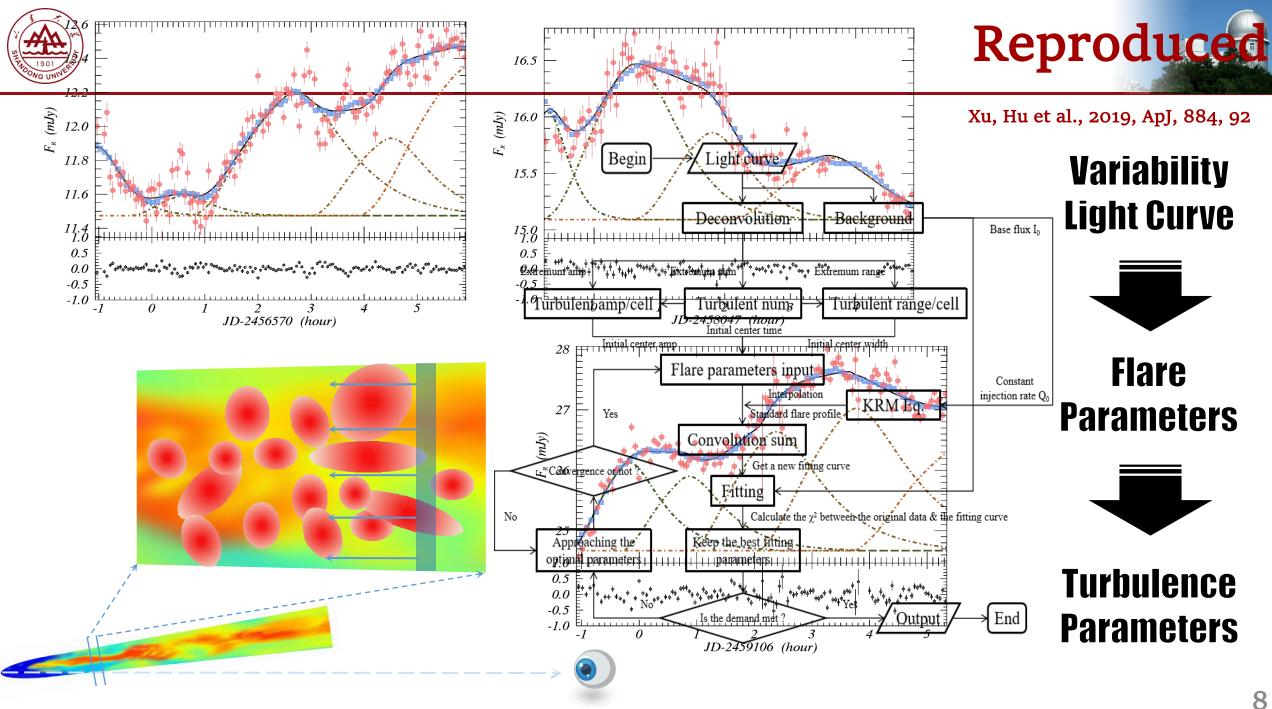
W2R 1926 + 42 **By Kepler**

Li, Hu et al. 2018, MNRAS 478, 172

JD-2454833 (Day)



JD-2454833 (Day)



upstream

Shock Acceleration





$$\frac{\partial N}{\partial t} + \frac{\partial}{\partial \gamma} \left[\left(\frac{\gamma}{t_{acc}} - \beta_s \gamma^2 \right) N \right] + \frac{N}{t_{esc}} = Q \delta(\gamma - \gamma_0)$$

$$\beta_s = \frac{4}{3} \frac{\sigma_r}{m_e c} \left(\frac{B^2}{2\mu_0} \right)$$

$$N(\gamma,t) = a \frac{1}{\gamma^2} \left(\frac{1}{\gamma} - \frac{1}{\gamma_{max}}\right)^{(t_{acc} - t_{esc})/t_{esc}} \times \theta(\gamma - \gamma_0) \theta(\gamma_1(t) - \gamma)$$

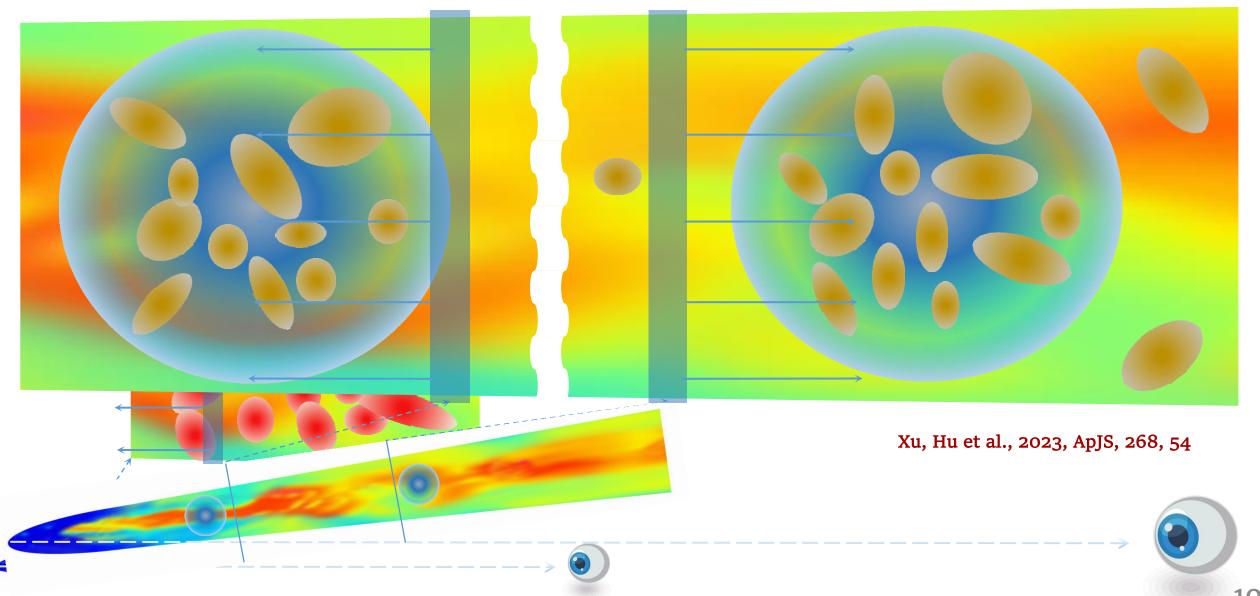
$$a = Q_0 t_{acc} \gamma_0^{t_{acc}/t_{esc}} (1 - \frac{\gamma_0}{\gamma_{max}})^{-t_{acc}/t_{esc}}$$

$$\gamma_1(t) = \left(\frac{1}{\gamma_{max}} + \left[\frac{1}{\gamma_0} - \frac{1}{\gamma_{max}}\right]e^{-t/t_{acc}}\right)^{-1}$$



Multi-zone Model

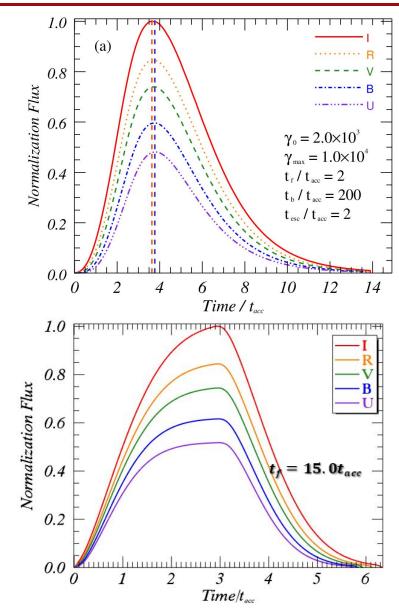


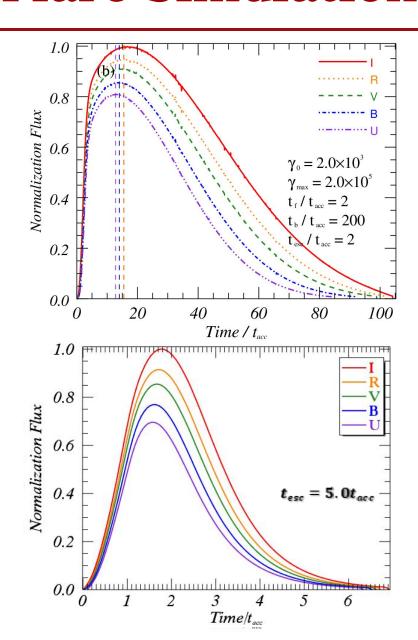


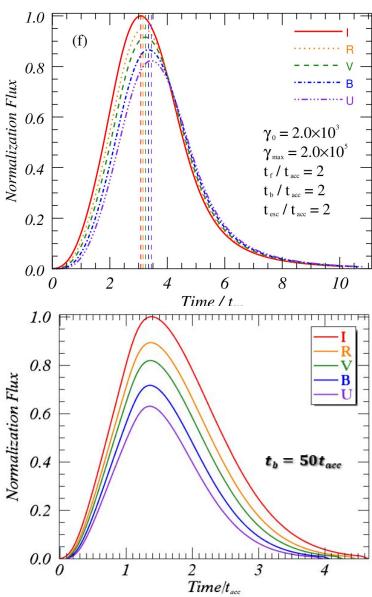


Flare Simulation





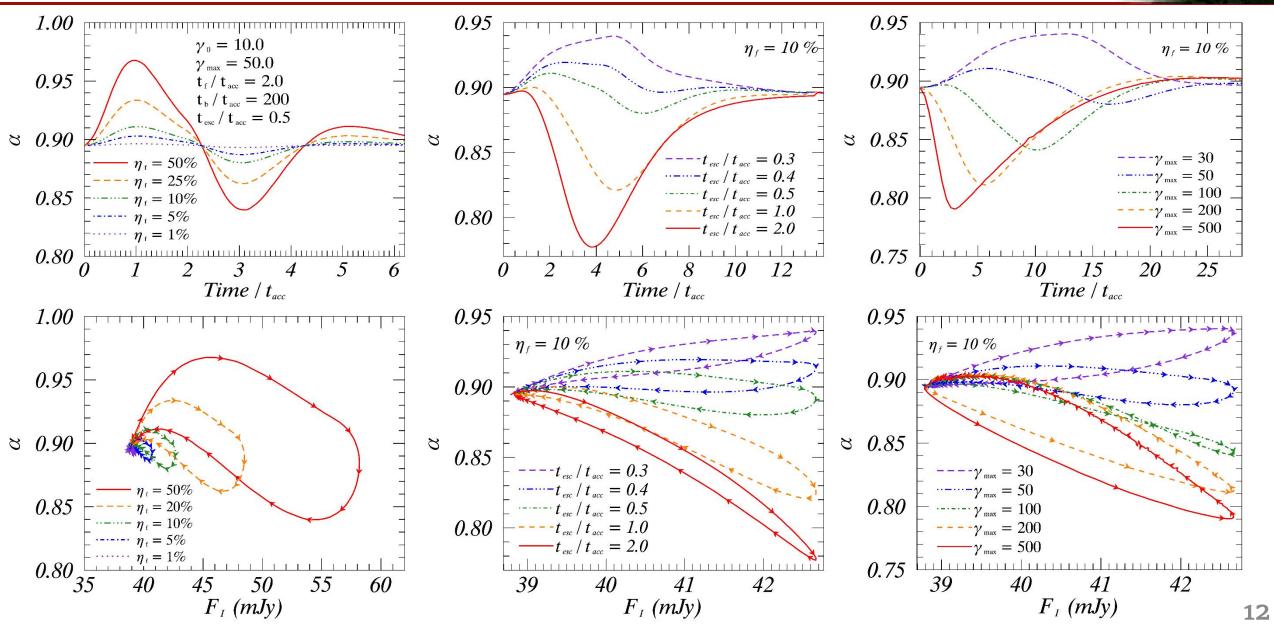






HTD/HID Evolution



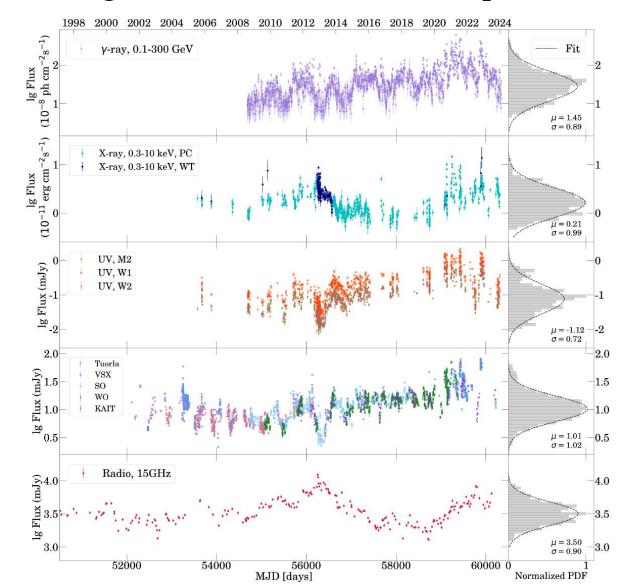




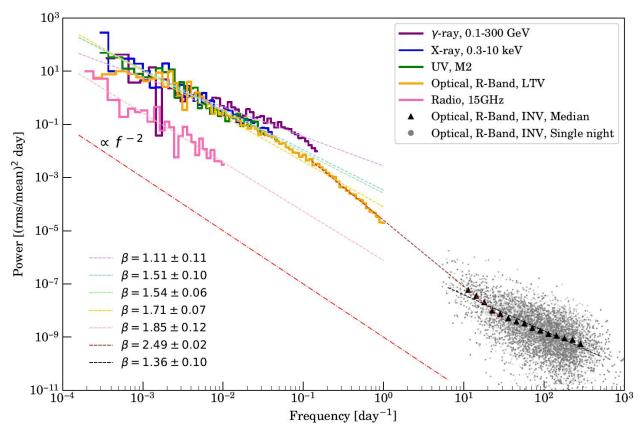
Variability → Power Spectrum



Taking BL Lacertae as an example



Classical Method — DFT



characteristic:

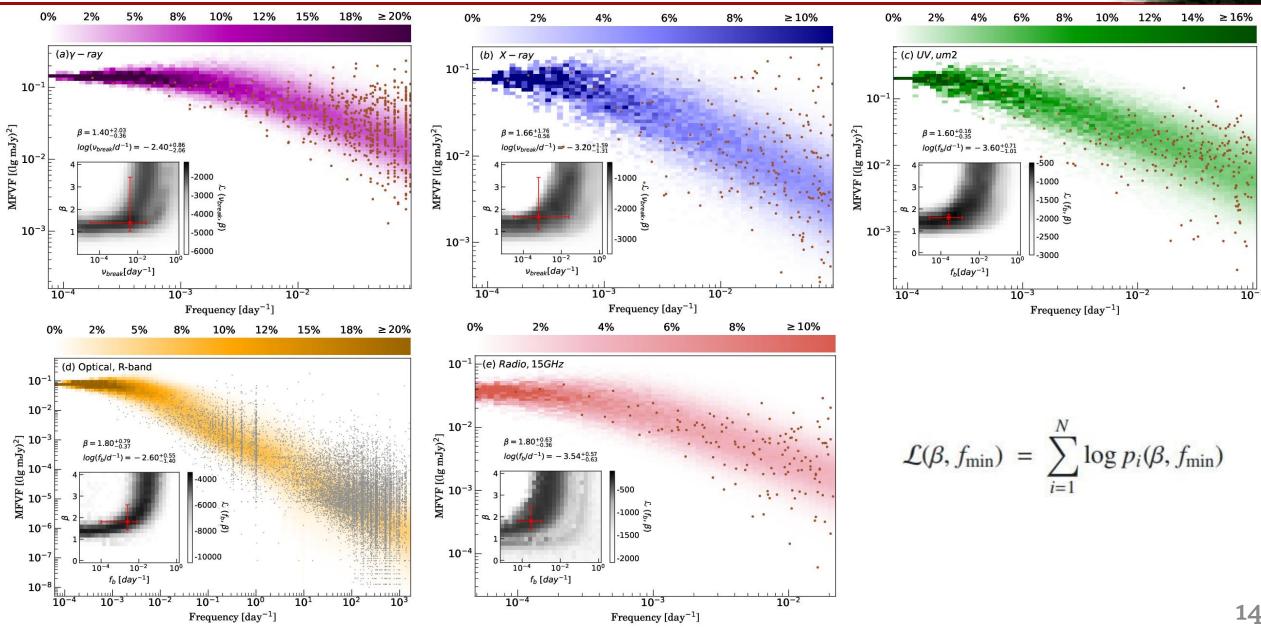
- Power-law index
- Broken Timescale
- Periodicity



MFVF

Broken Timescale Power-law index

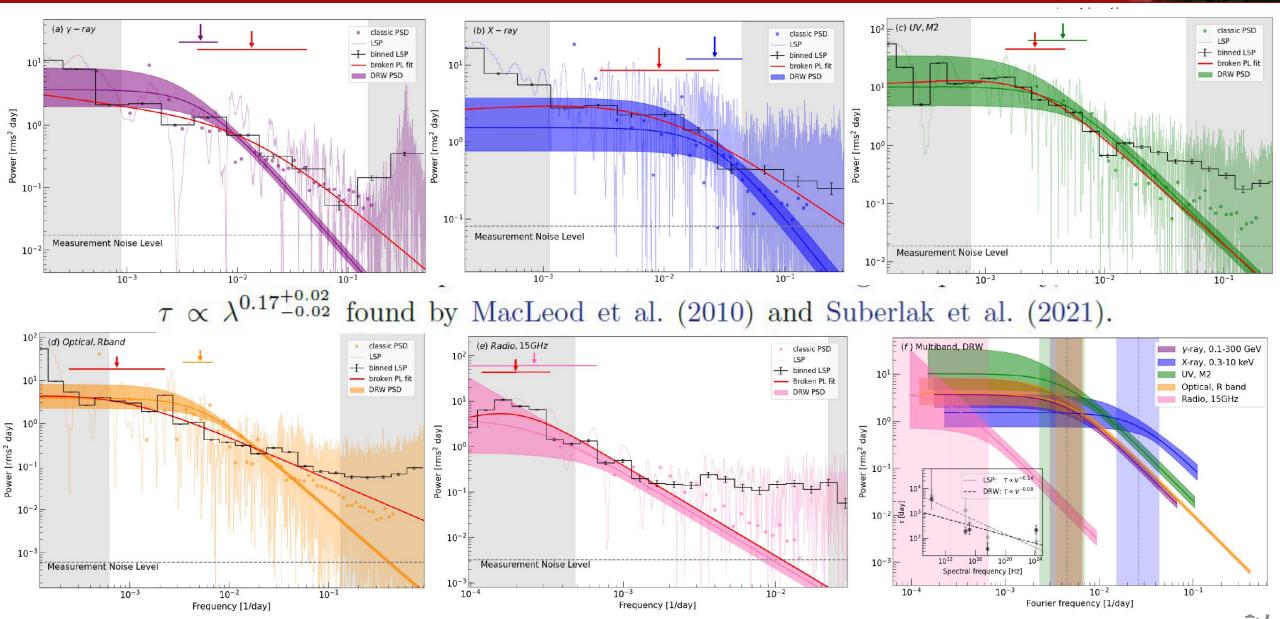






DRW / AR1 / OU

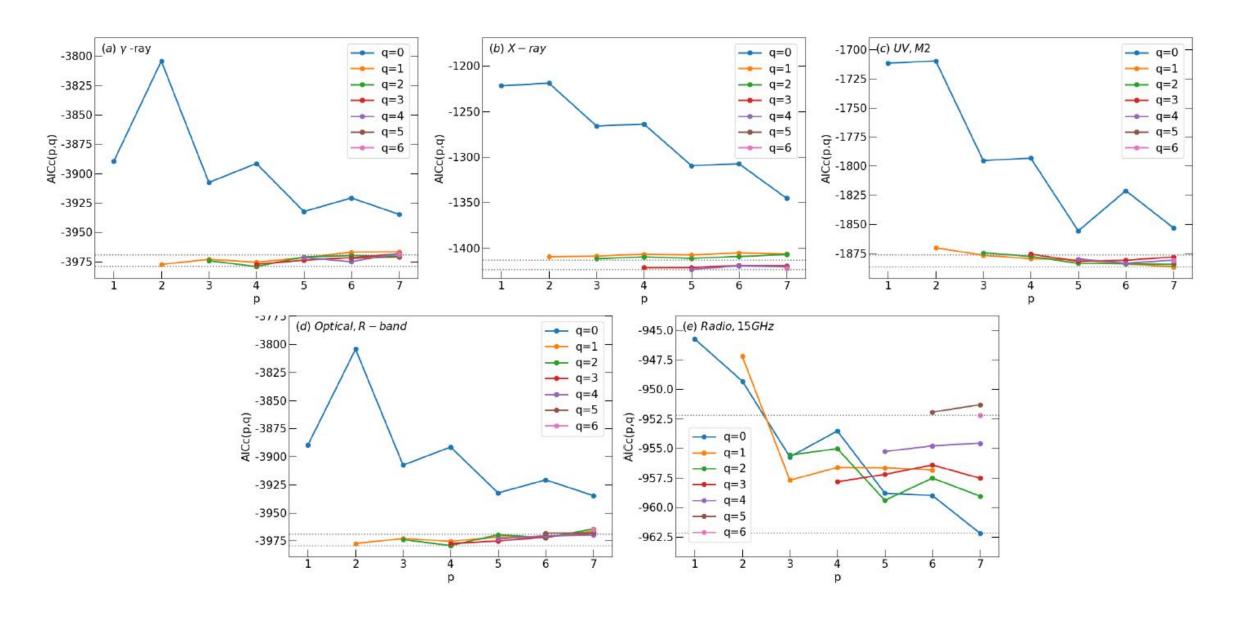






CARMA / Mixed OU



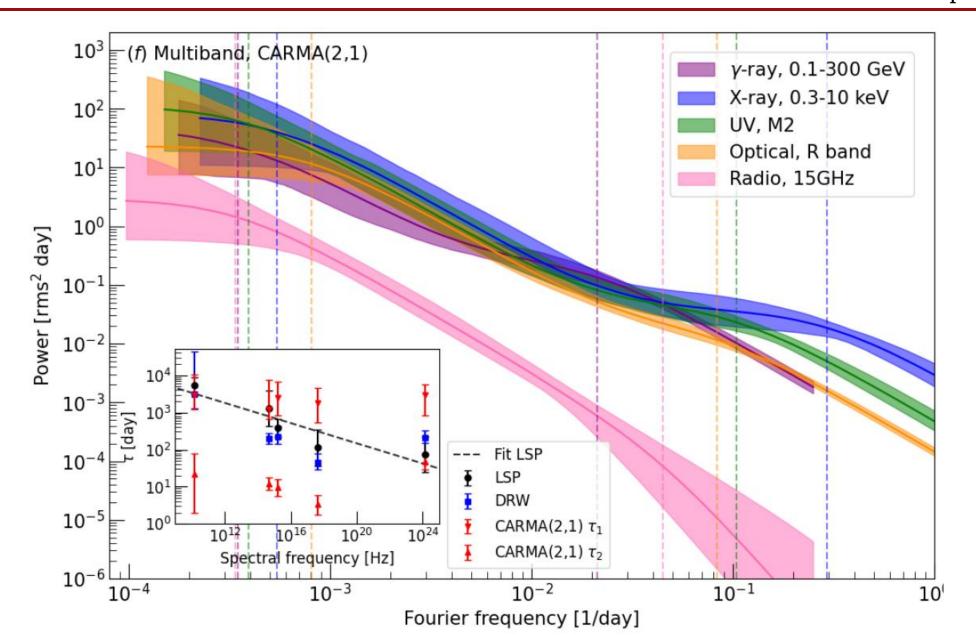




CARMA / Mixed OU

Taking **CARMA(2,1)** as an example







Summary



- The broken power-law model applied to the LSP reveals wavelength-dependent characteristic timescales, whereas the DRW model yields nearly identical timescales across optical, UV, and γ-ray bands. For higher-order CARMA models, the obtained timescales show even greater variability. Variability is better explained by the superposition of multiple stochastic processes;
- The DHO simulations reveal that the characteristic timescales and power spectra are similar in the low-frequency range, while notable differences emerge at high frequencies. We therefore infer that the variability in BL Lacertae is primarily governed by at least two stochastic processes: one that dominates at low frequencies and is weakly dependent on wavelength or even independent of it, and another that dominates at high frequencies and is wavelength-dependent.



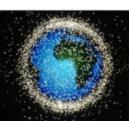
Outline

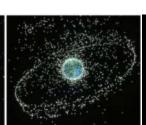




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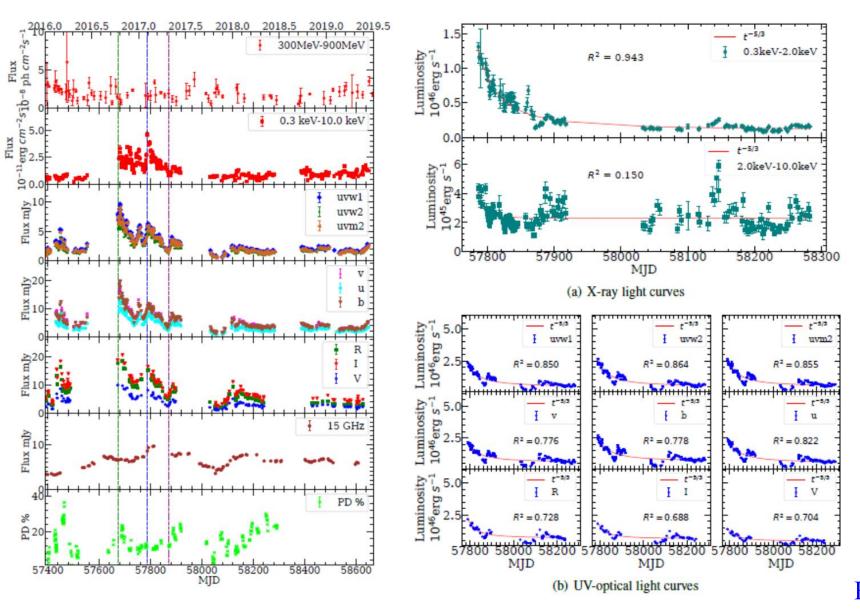


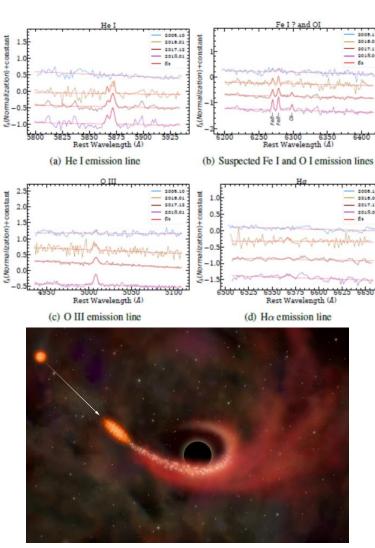




TDE --- the variability of OJ 287 in 2016-2017







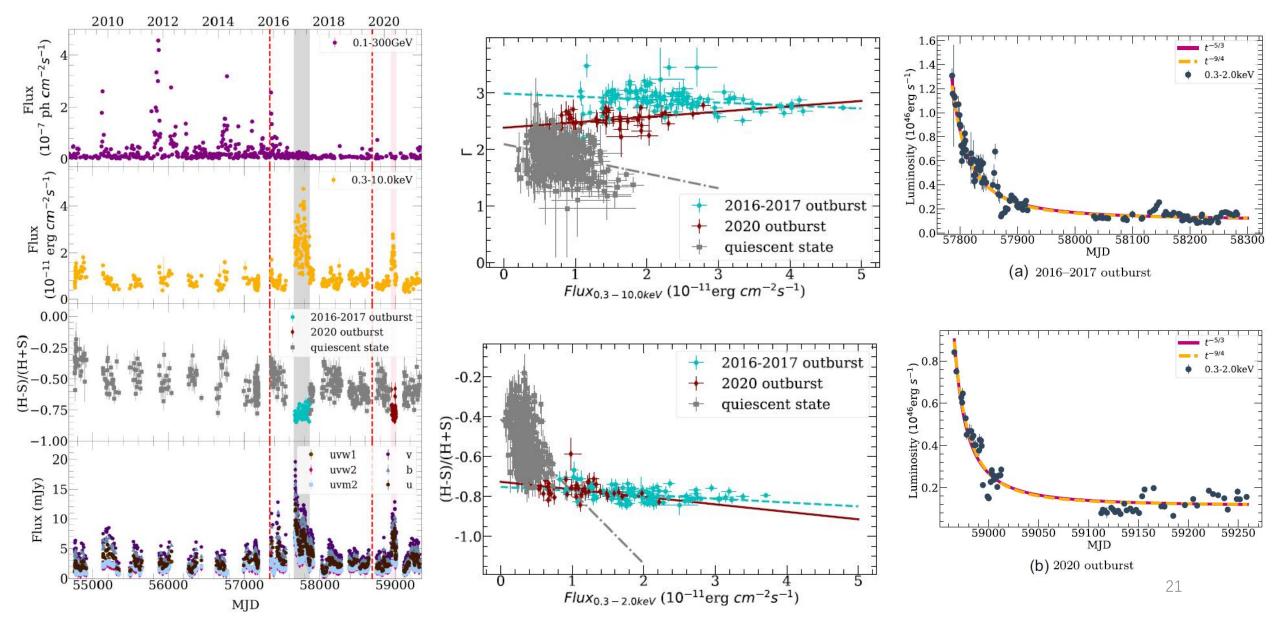
Huang, Hu et al., ApJ, 920, 12, 2021.



PTDE—Variability of OJ 287 in 2020

Huang, Hu et al., MNRAS, 515, 2778, 2022







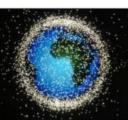
Outline

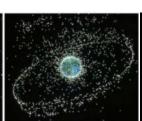




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1m Telescope



Optics: classic Cassegrain design, 1000mm optical working diameter, focal ratio: f/8

Mount: equatorial fork mount

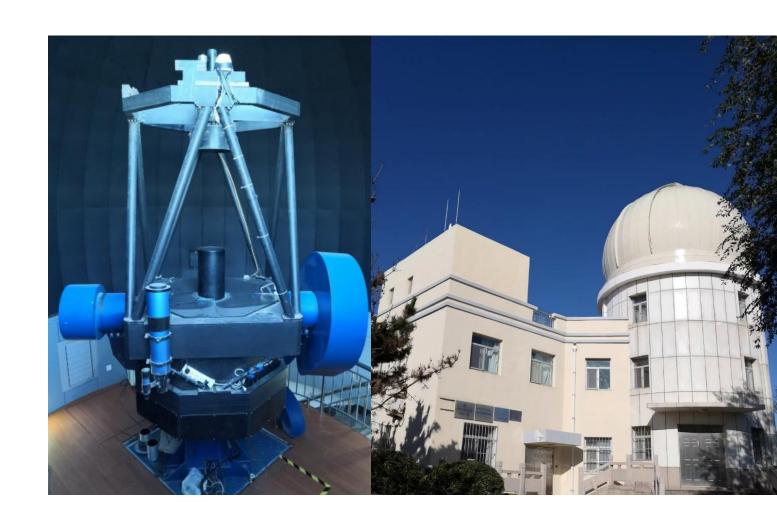
Image quality: 80% Energy in 0.65 arc sec within 15 arc minute FOV

Maximum slew speed is 4 degrees per second in both RA and DEC axis

Pointing accuracy: 5.4 arc sec RMS for 20-90 deg altitude

Tracking accuracy: 0.54 arc sec RMS in 10 minutes blind guiding

1m telescope and dome





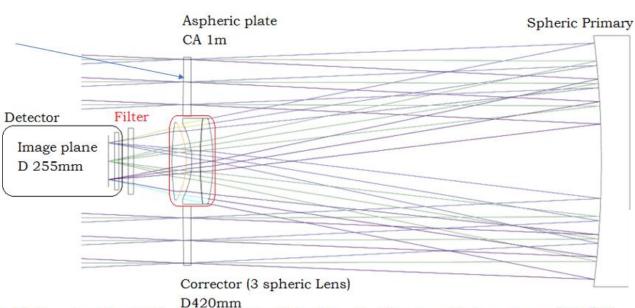
Wide field 1m telescope



□ Optical design for Sitian Project

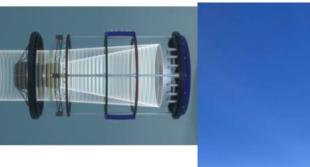
SITIAN- Prototype optics (Catadioptric)

Aperture 1m; Primary 1.3m; F# 2.06; FOV 5°×5°



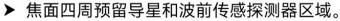
◆ Refer to the optical system of Antarctic Survey Telescopes (AST3)

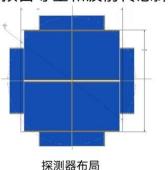
For Time domain Astronomy (Varaibility of AGNs, TDEs, Variables, SNs, Exoplanets, GRBs et al.), Installed in 2024

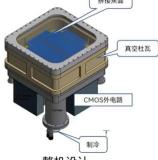




➤ 焦面尺寸为200mm×200mm,采用4片9K×9K、10μm背照式CMOS芯片拼接。



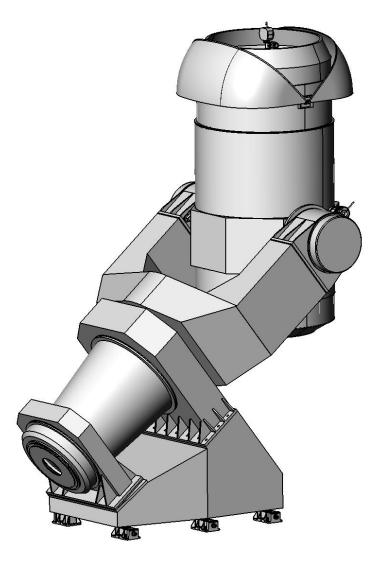






Progress





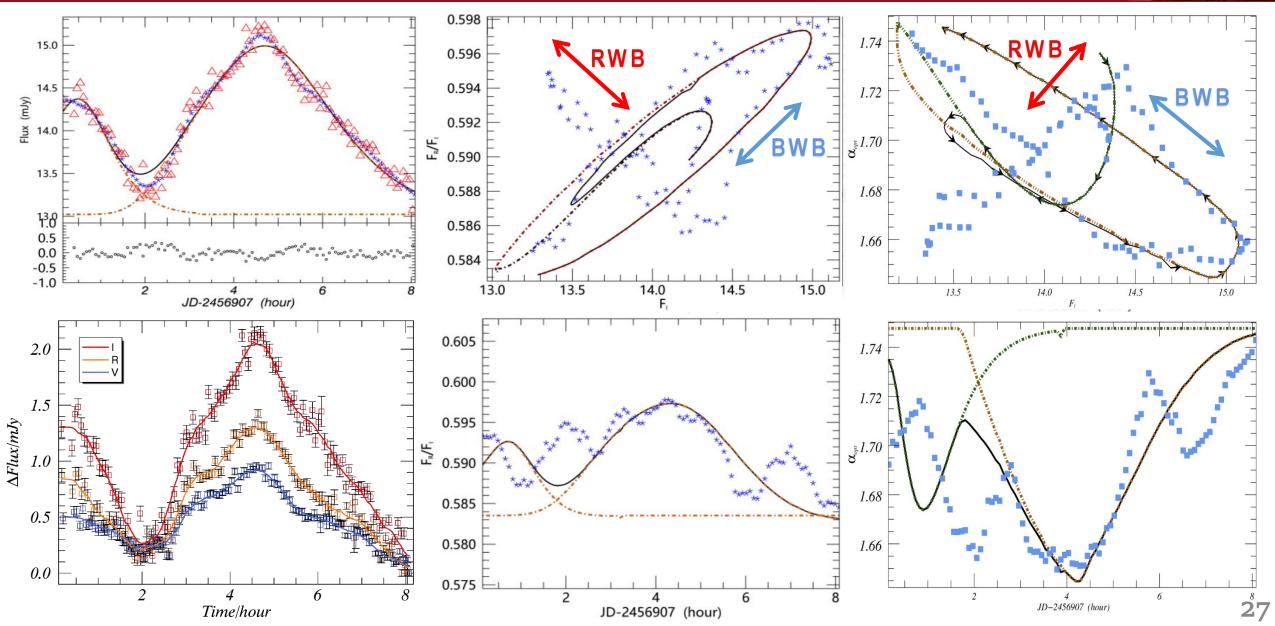






Fitting



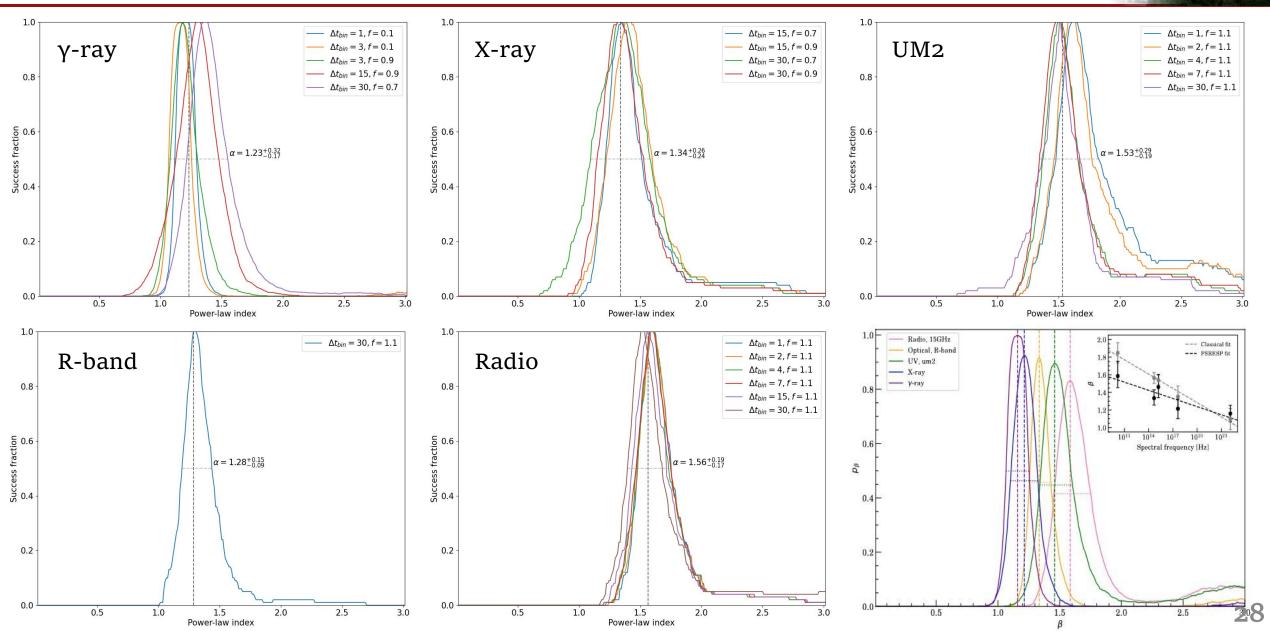




PSRESP



Power-law index

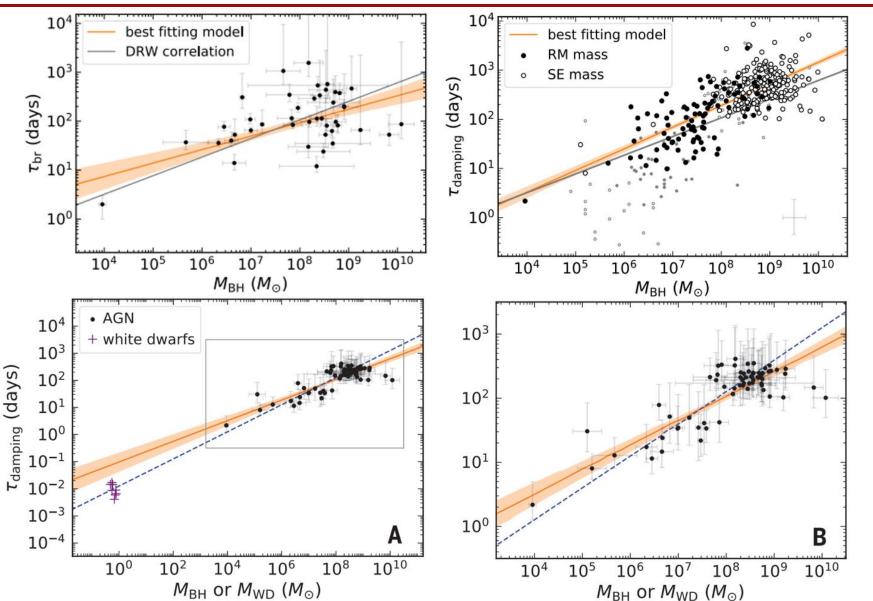




$au_{dampling}$ - $M_{accretor}$ Relation







$$\tau_{\text{damping}} = 199^{+11}_{-11} \text{ days } \left(\frac{M_{\text{BH}}}{10^8 M_{\odot}}\right)^{0.44^{+0.02}_{-0.02}}$$

$$\tau_{\rm damping} \propto \lambda^{0.17}$$

$$R_{\lambda} \propto M_{\rm BH}^{2/3} \longrightarrow t_{\rm orb}, t_{\rm th} \propto M_{BH}^{1/2}$$

$$\eta \propto L_{\rm bol}/\dot{M}_{\rm BH}$$

$$t_{th} \approx (2\pi\alpha)^{-1} t_{orb} = \alpha^{-1} t_{dyn} \approx (H/R)^2 t_{vis}$$

$$au_{
m damping} = 107^{+11}_{-12} {
m days} \left(rac{M_{
m BH}}{10^8\,M_{\odot}}
ight)^{0.38^{+0.05}_{-0.04}} \ ag{M_{
m BH}} = 10^{7.97^{+0.14}_{-0.14}} \, M_{\odot} \, \left(au_{
m damping}/100 \, {
m days}
ight)^{2.54^{+0.34}_{-0.35}} \ ag{100}$$

$$t_{
m orb} = 100 igg(rac{M_{
m BH}}{10^8\,M_{\odot}} igg) igg(rac{R}{100\,R_S} igg)^{3/2} {
m days}$$

$$t_{
m th} = 1680 \Big(rac{lpha}{0.01}\Big)^{-1} imes \Big(rac{M_{
m BH}}{10^8\,M_\odot}\Big) \Big(rac{R}{100\,R_S}\Big)^{3/2} {
m days}$$

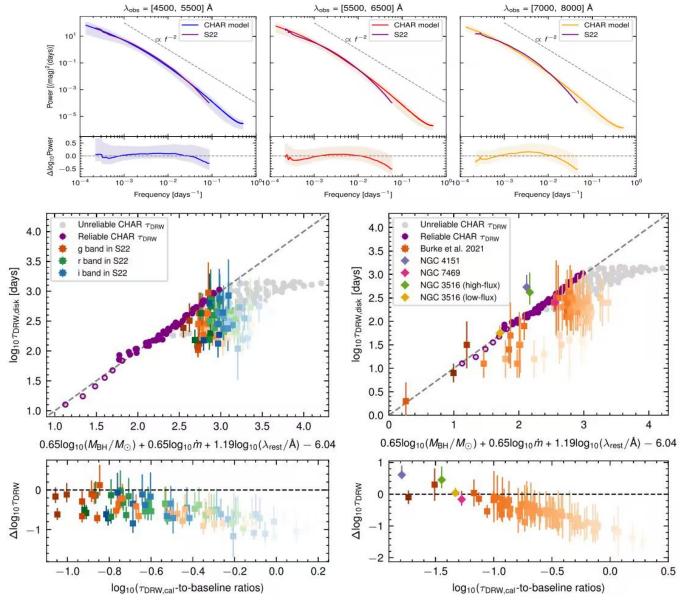


Accretion Disk Model

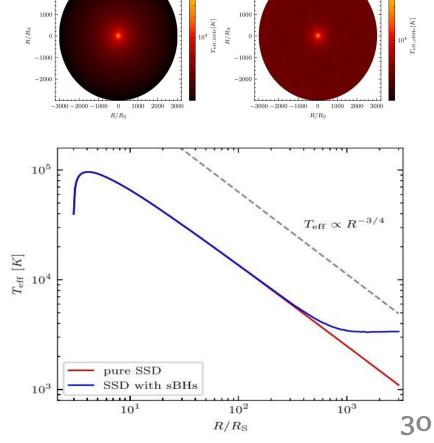
CHAR model



Corona Heated Accretion disk Reprocessing Model



Stellar Black Holes Can "Stretch" Supermassive Black Hole Accretion Disks



Cite from: Zhou Shuying et al., 2024, ApJL, 966, L9

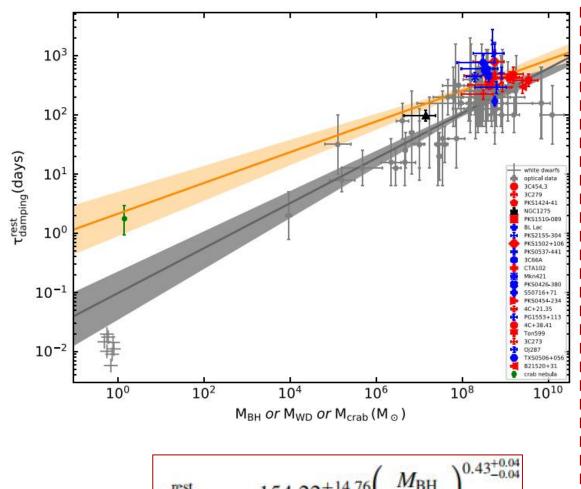
Cite from: Zhou Shuying et al., 2024, ApJ, 966, 8



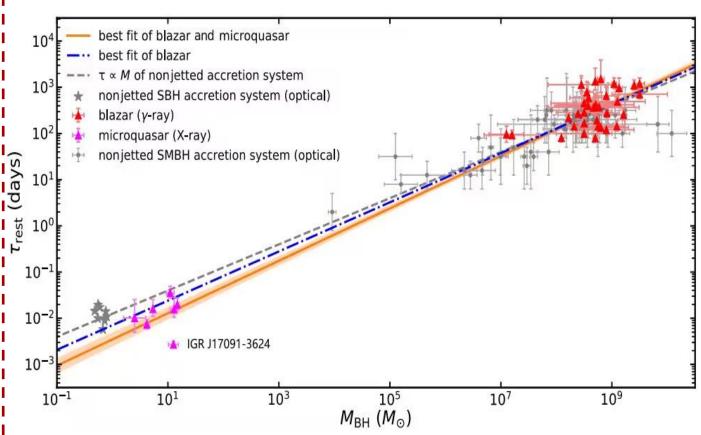
Disk-Jet Coupling Model







$$\tau_{\text{damping}}^{\text{rest}} = 154.22_{-15.75}^{+14.76} \left(\frac{M_{\text{BH}}}{10^8 M_{\odot}} \right)^{0.43_{-0.04}^{+0.04}}$$



$$\tau_{rest} = 120.47^{+14.84}_{-17.62} \text{ days} \left(\frac{M_{BH}}{10^8 M_{\odot}}\right)^{0.57^{+0.02}_{-0.02}}$$



Turbulence Jet Model DISA Diffusive shock acceleration



