

Cosmic-Ray boosted DM and its effect on the DM-electron direct detection

Ding Ran

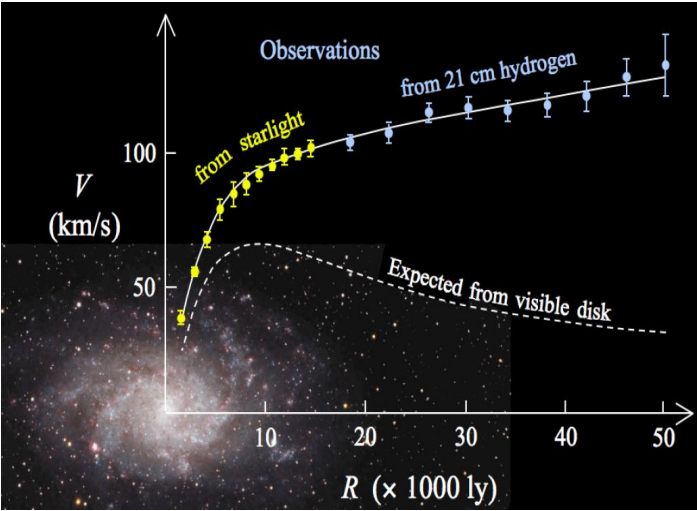
CHEP, PKU

Collaborated with Xiang Qianfei & Cao Qinghong

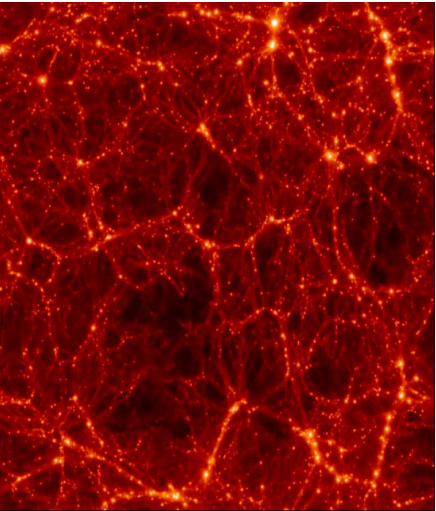
2nd TDLI Mini-Workshop on "New Physics at the Tera Scale"
2019.08.04

Evidences for DM

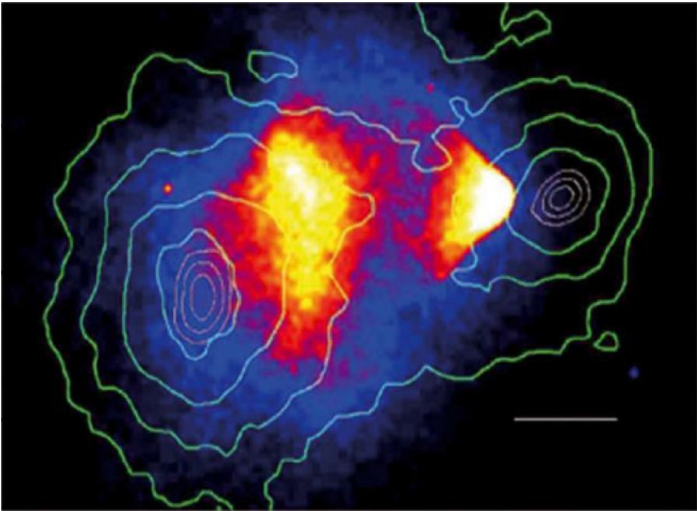
- Rotation curves



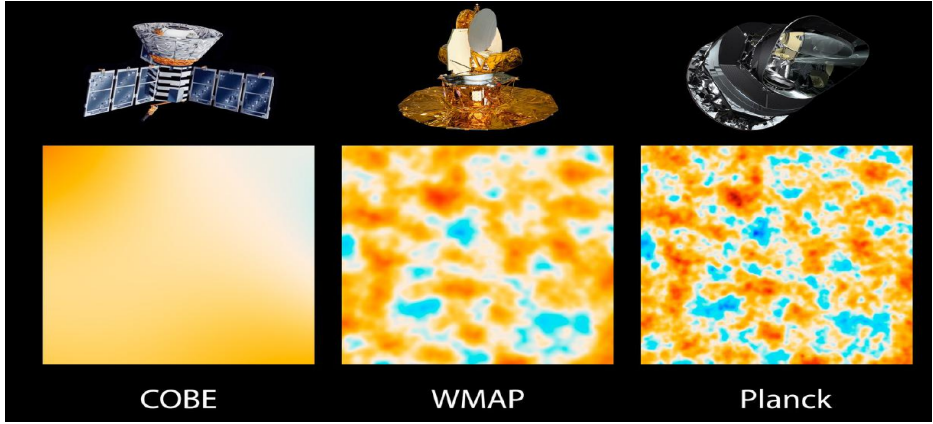
- LSS



- Bullet cluster



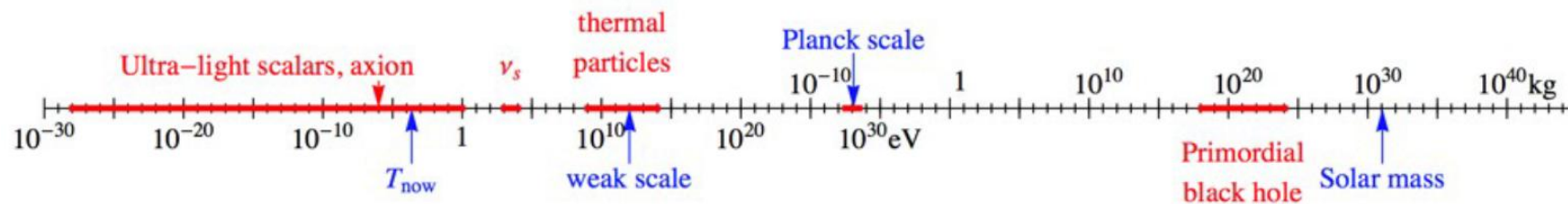
- CMB

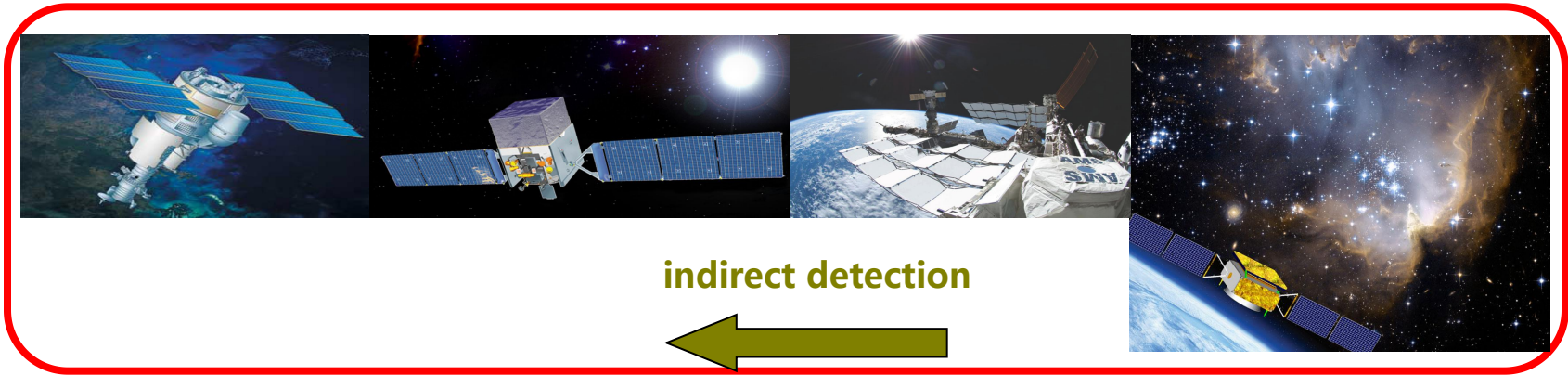


Dark sector candidates



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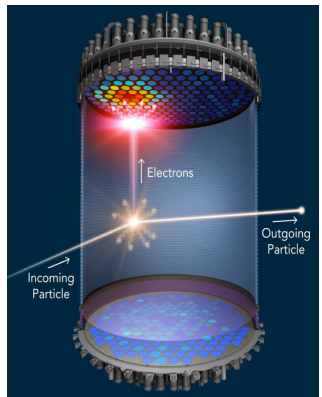


indirect detection

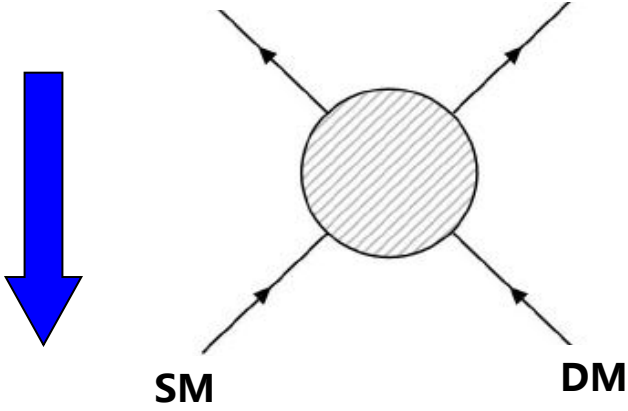


Cosmic-Ray (CR) physics

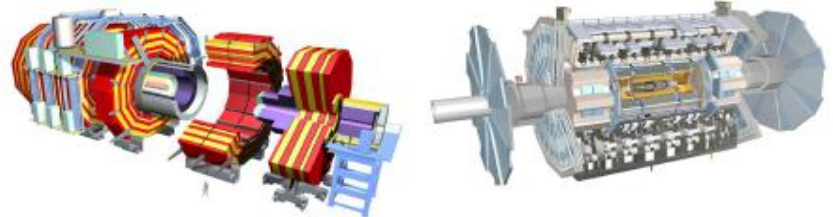
direct detection

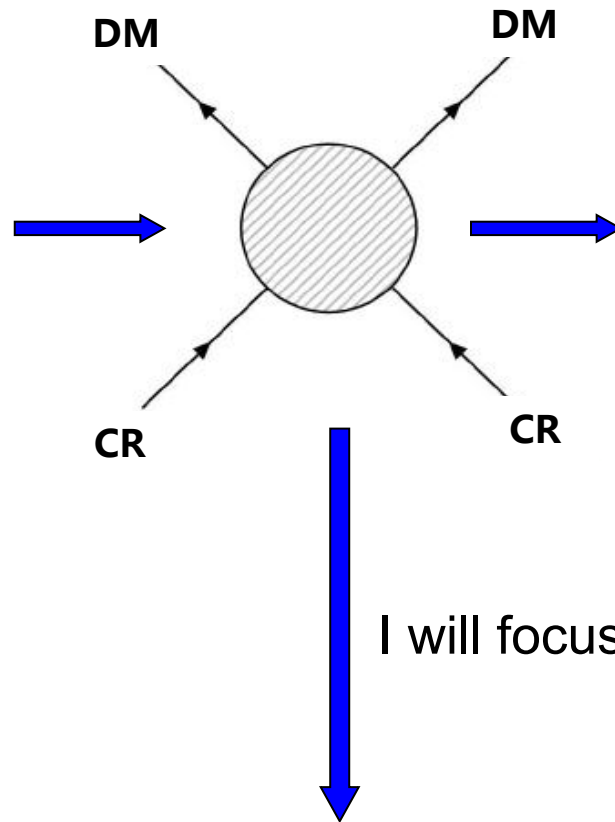
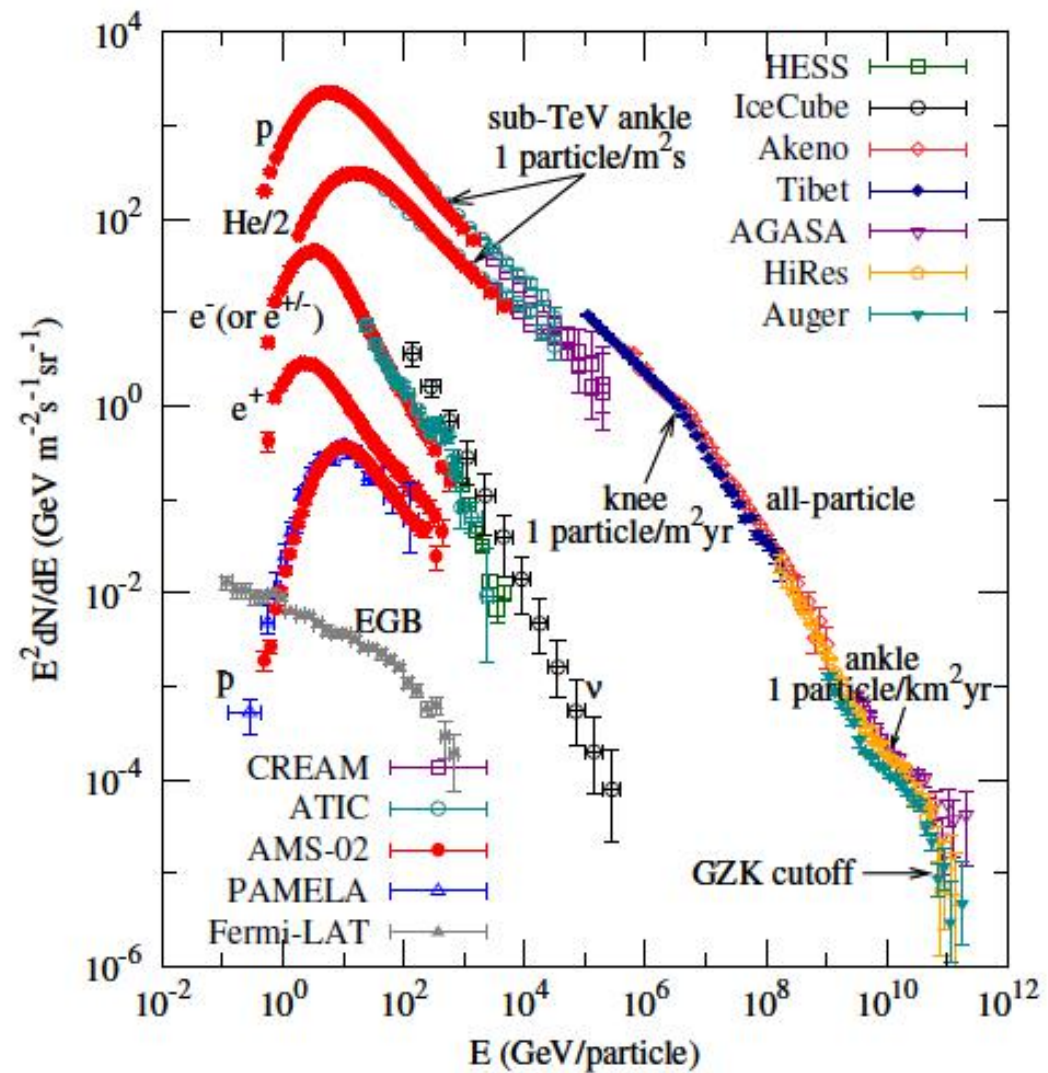


SM DM

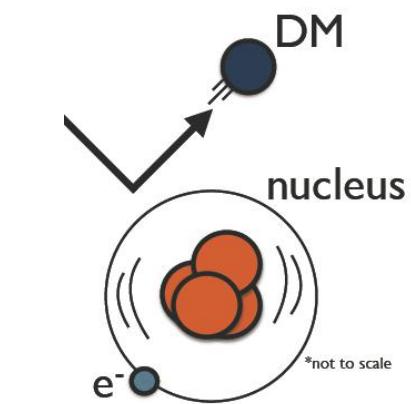


collider detection

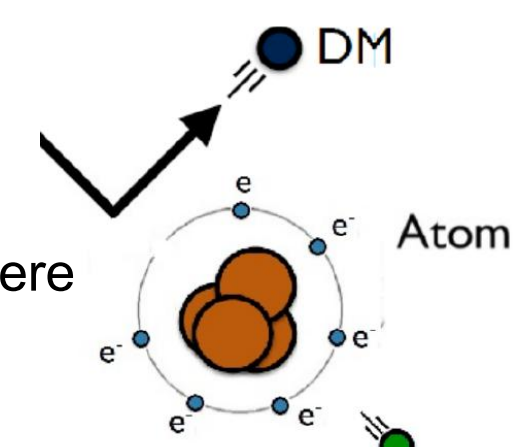




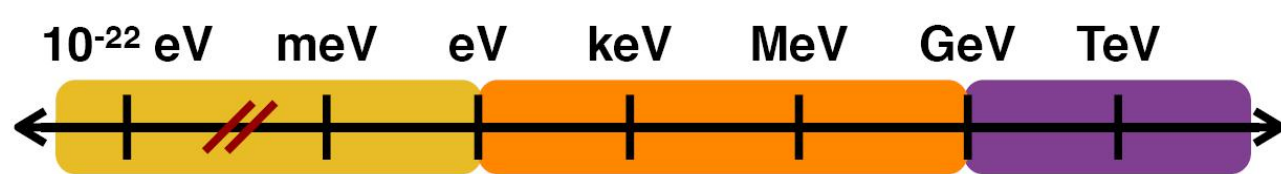
DM-nucleon scattering



DM-electron scattering

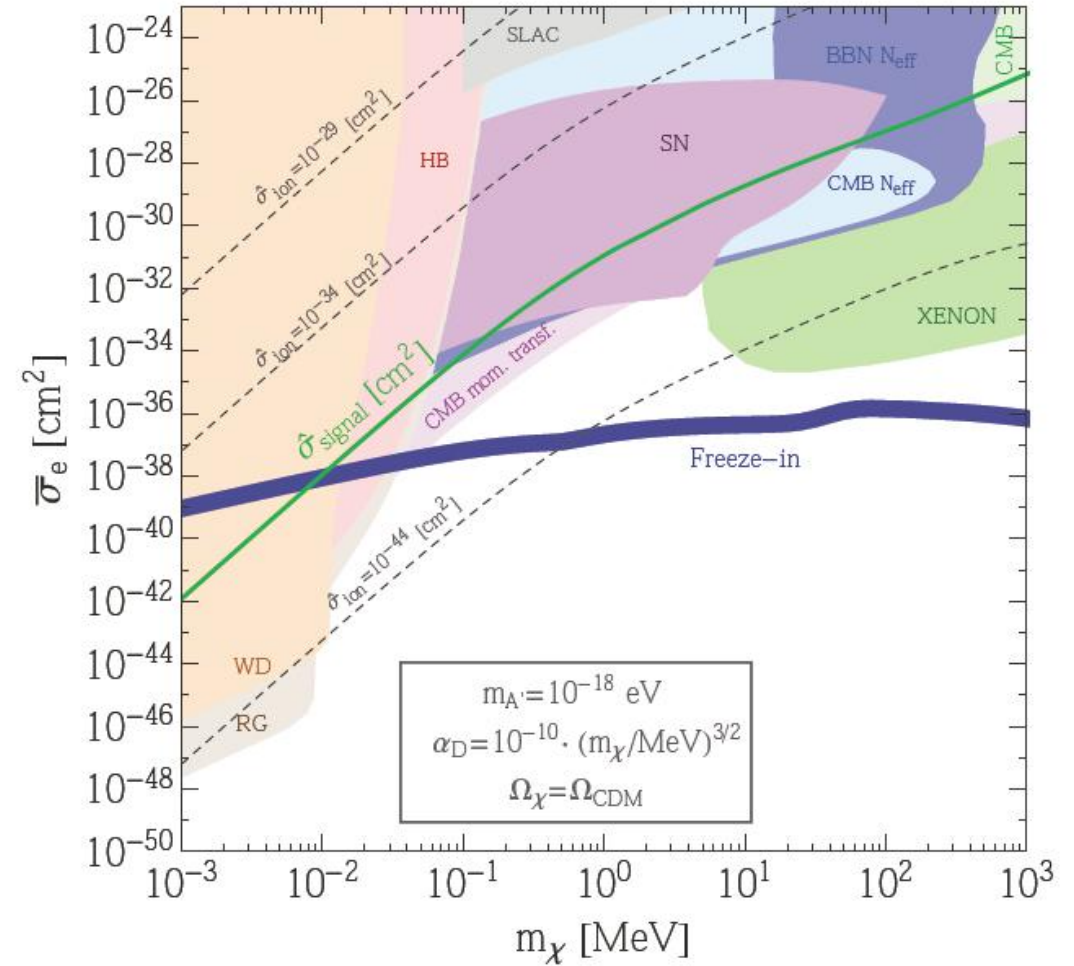
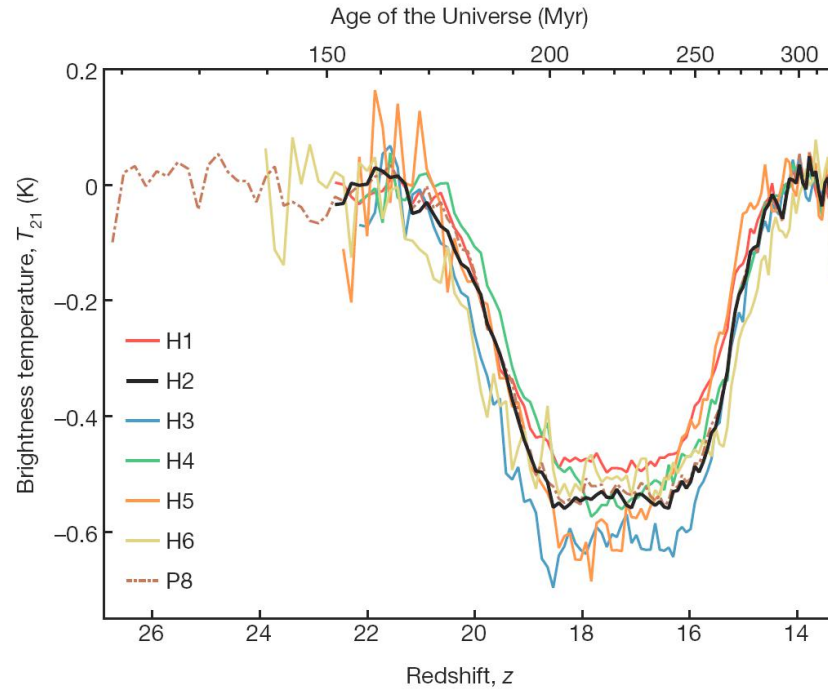
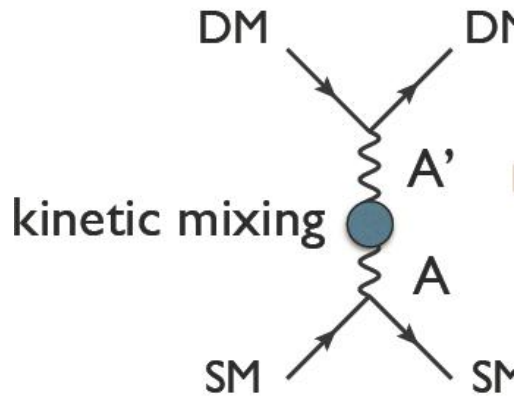


I will focus here



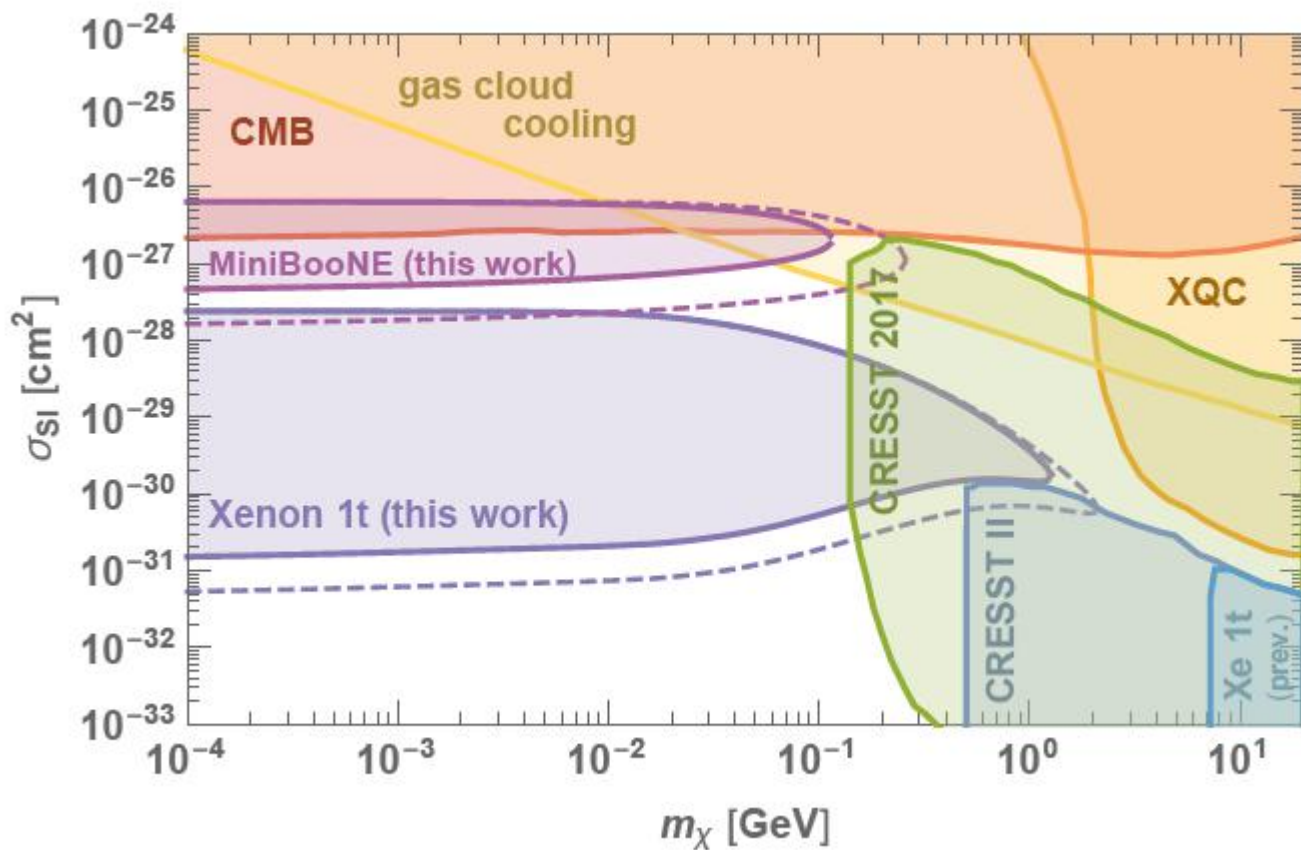
Theoretical motivation of light DM

- Hidden Photon Mediator



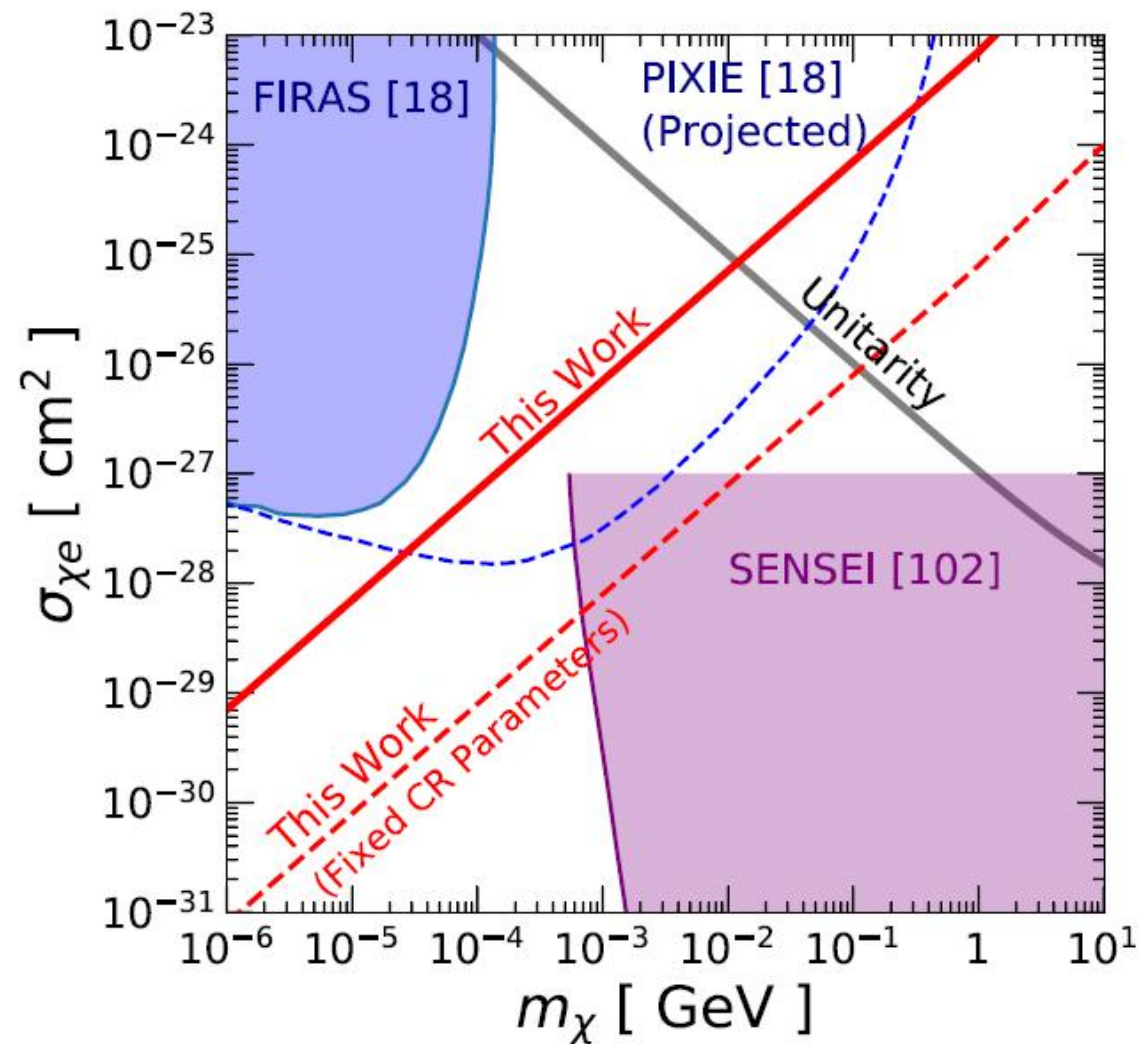
Limits from DM-nucleon direct detection

T. Bringmann and M. Pospelov, arXiv:1810.10543



Reverse Direct Detection

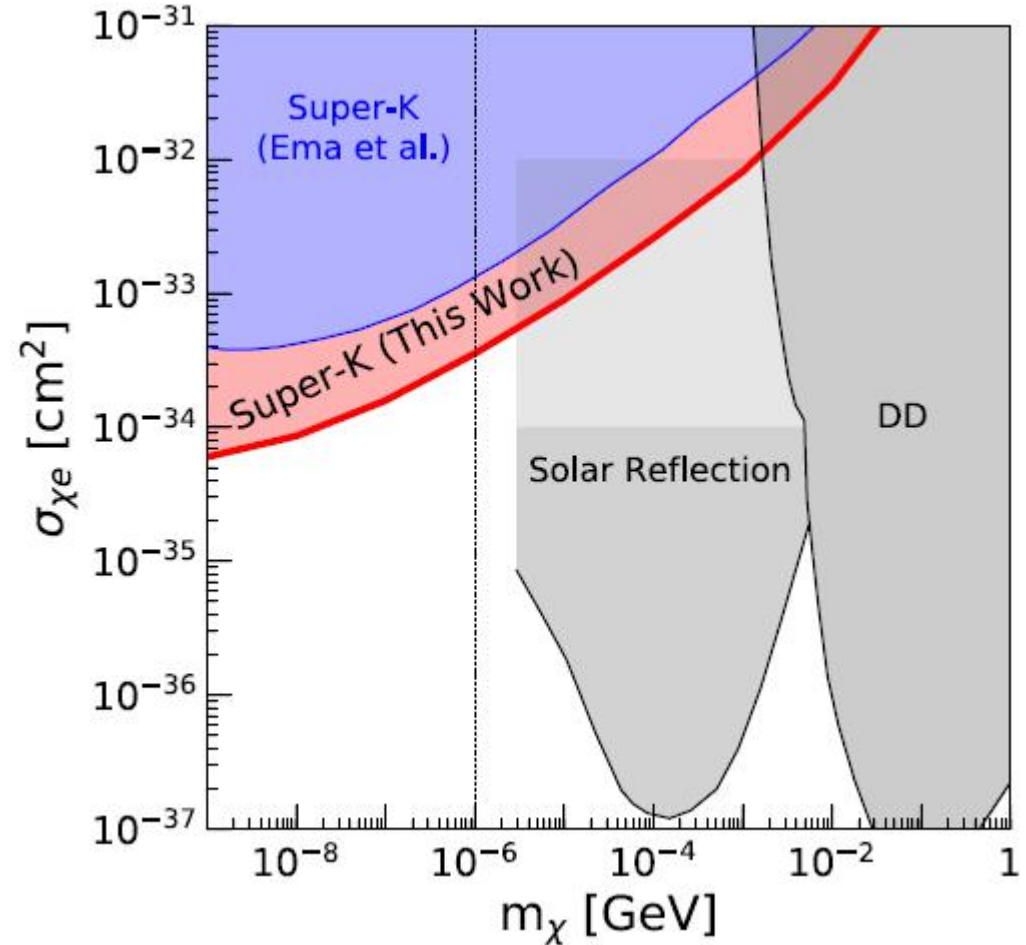
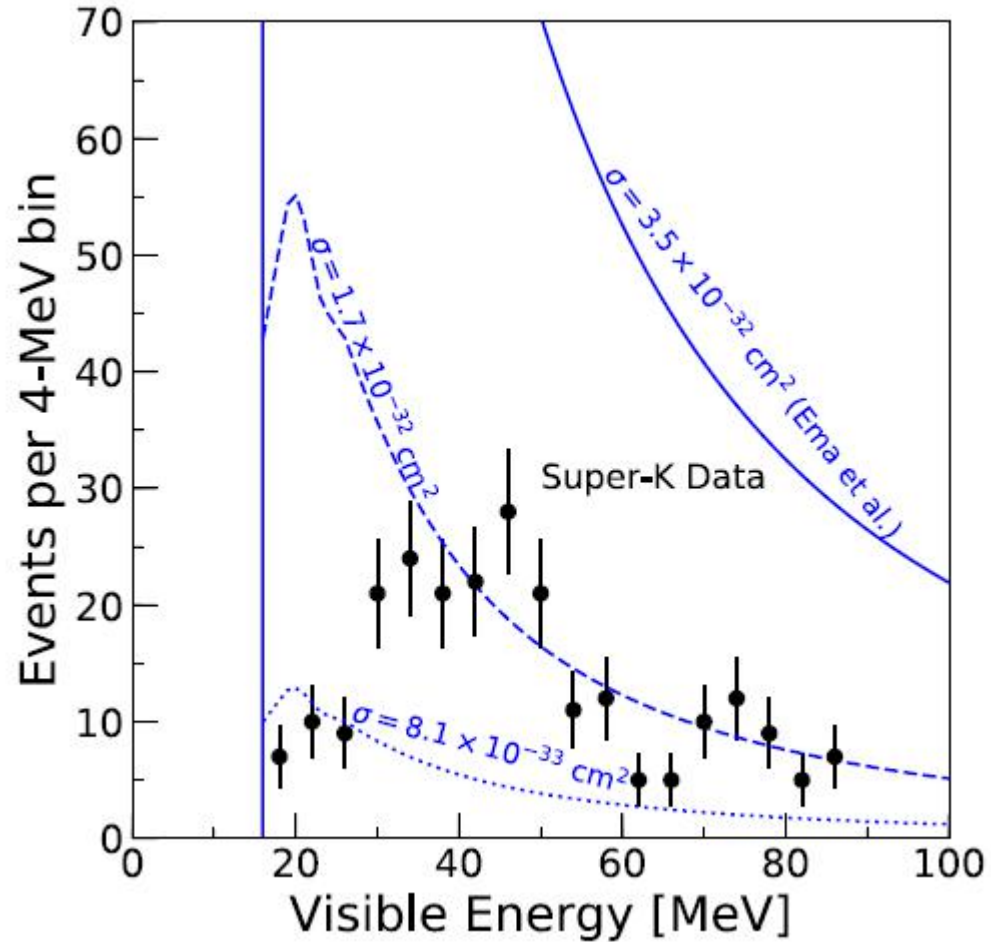
C. V. Cappiello, K. C. Y. Ng and J. F. Beacom, arXiv:1810.07705



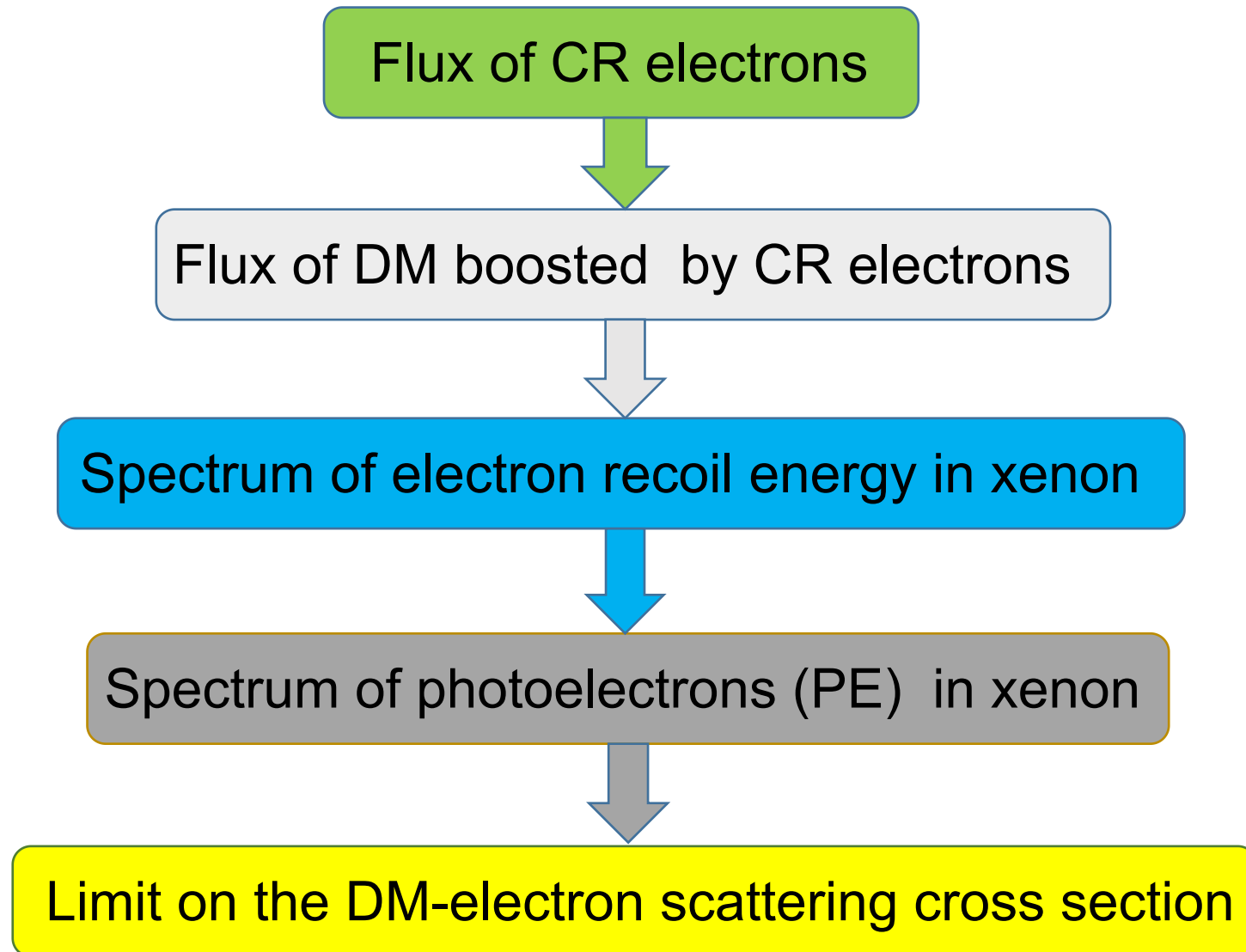
Limits from Neutrino Experiments

Y. Ema, F. Sala and R. Sato, arXiv:1811.00520

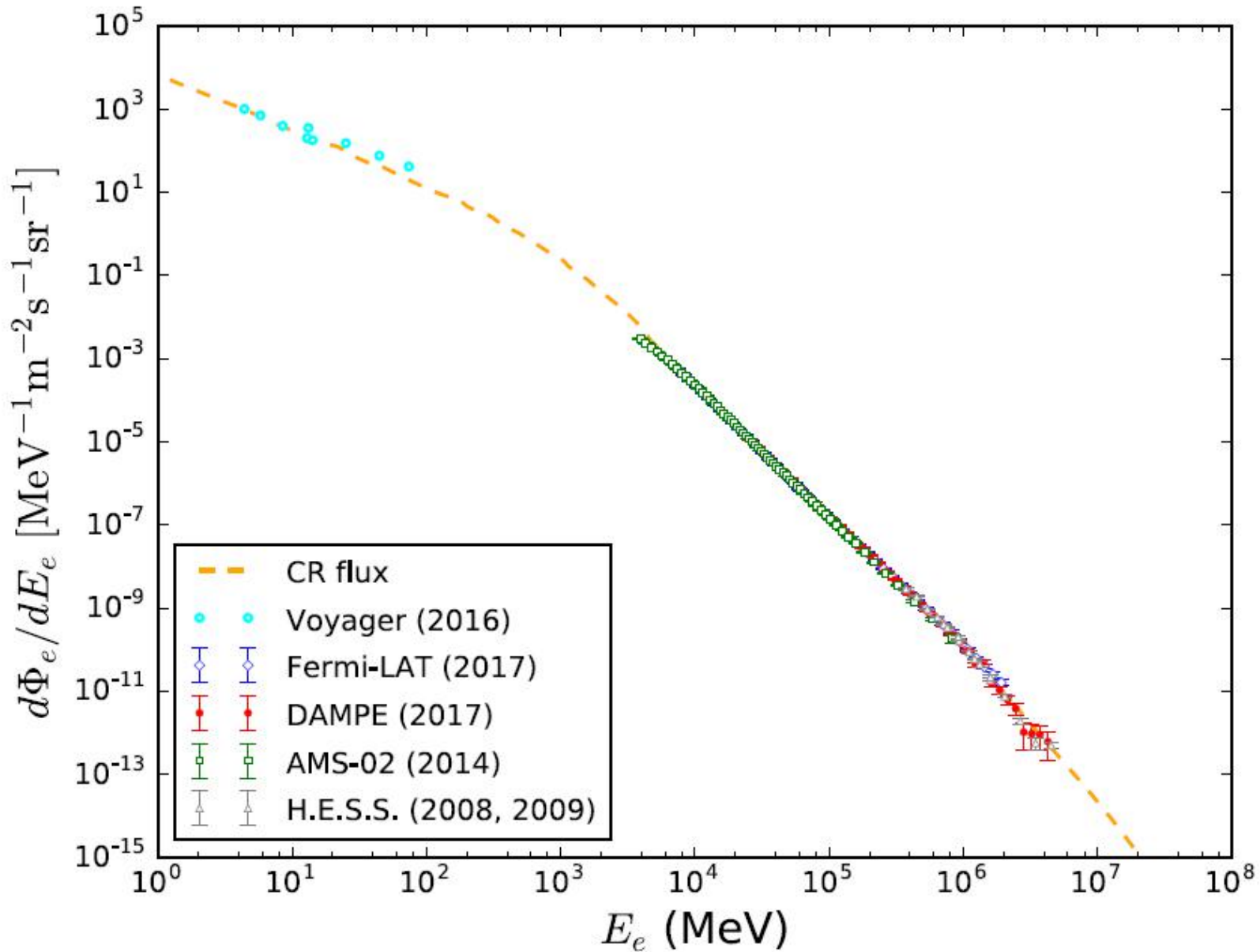
C. Cappiello and J. F. Beacom, arXiv:1906.11283



Calculation framework



Flux of CR electrons

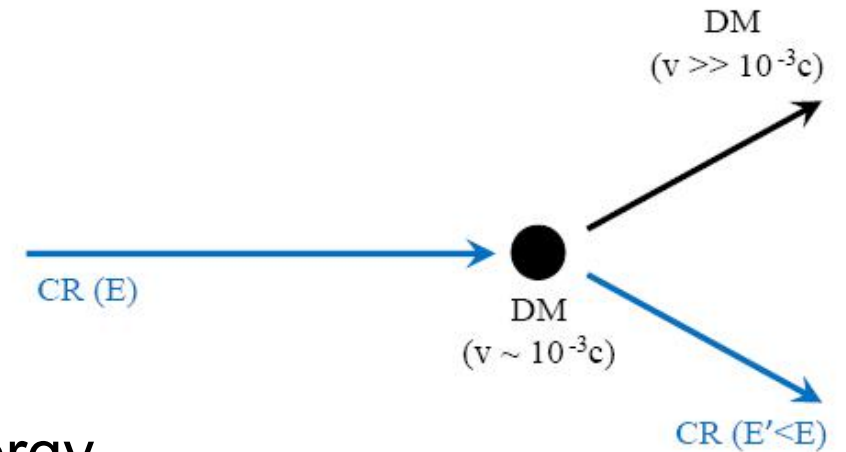


Flux of DM boosted by CR electrons

- Kinematics of DM-CR scattering

$$T_{\chi}^{\max} = \frac{2m_{\chi}T_{\text{CR}}(T_{\text{CR}} + 2m_{\text{CR}})}{(m_{\text{CR}} + m_{\chi})^2 + 2T_{\text{CR}}m_{\chi}}$$

$$T_{\text{CR}}^{\min} = \left(\frac{T_{\chi}}{2} - m_{\text{CR}} \right) \left(1 \pm \sqrt{1 + \frac{2T_{\chi}}{m_{\chi}} \frac{(m_{\text{CR}} + m_{\chi})^2}{(2m_{\text{CR}} - T_{\chi})^2}} \right)$$



- Differential flux at Earth in terms of the CR energy

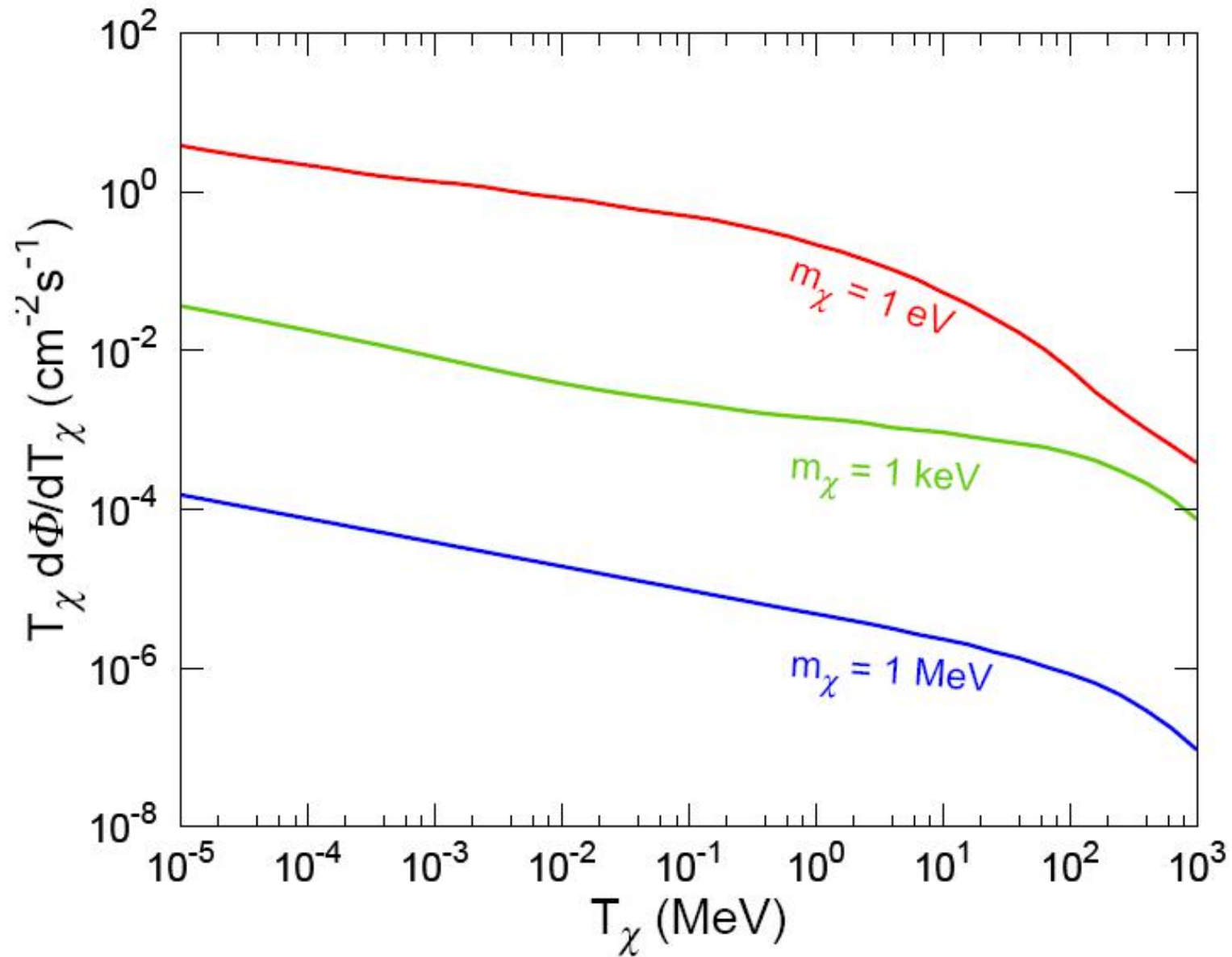
$$\frac{d\Phi_{\chi}}{dT_i} = \int \frac{d\Omega}{4\pi} \int_{l.o.s} dl \sigma_{\chi i} \frac{\rho_{\chi}}{m_{\chi}} \frac{d\Phi_i}{dT_i} \equiv \sigma_{\chi i} \frac{\rho_{\chi}}{m_{\chi}} \frac{d\Phi_i}{dT_i} D_{\text{eff}}$$

line-of-sight integration out to 10 kpc $D_{\text{eff}} = 8.02$ kpc

- Convert into a DM energy spectrum

$$\frac{d\Phi_{\chi}}{dT_{\chi}} = \int_0^{\infty} dT_i \frac{d\Phi_i}{dT_i} \frac{1}{T_{\chi}^{\max}(T_i)} \Theta [T_{\chi}^{\max}(T_i) - T_{\chi}]$$

Flux of DM boosted by CR electrons



Spectrum of electron recoil energy in xenon

- Differential scattering rate

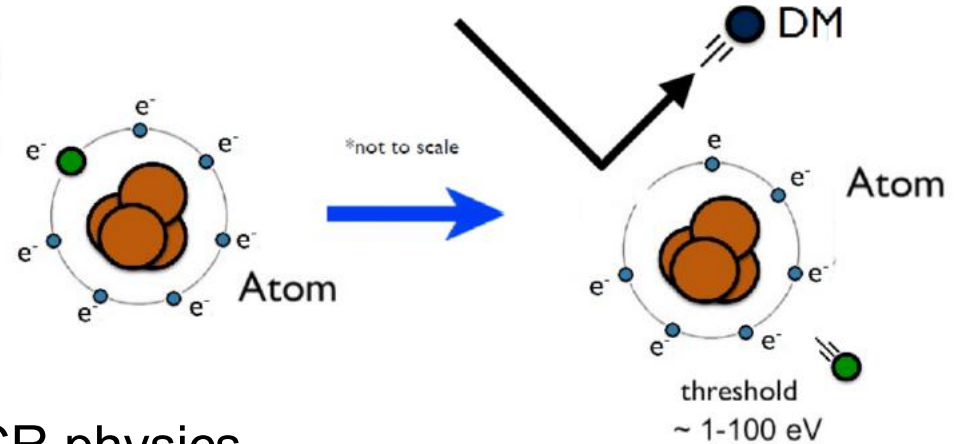
R. Essig, J. Mardon and T. Volansky, arXiv:1108.5383

DM ●

$$\chi(p) + e(k) \rightarrow \chi(p') + e(k')$$

particle physics

$$\frac{d\langle\sigma_{ion}^{nl}v\rangle}{d\ln E_R} = \frac{\sigma_e}{8\mu_{\chi e}^2} \int q dq |F_{DM}(q)|^2 |f_{ion}^{nl}(k', q)|^2 \eta(E_{\min})$$



CR physics

solid state physics

$$\eta(E_{\min}) = \int_{DM} \frac{m_{\chi}^2}{p_{\chi} E_{\chi}} = \frac{1}{\rho_{\chi}} \int_{E_{\min}} dE_{\chi} \frac{m_{\chi}^2}{p_{\chi} E_{\chi}} \frac{d\Phi_{\chi}}{dE_{\chi}}$$

In non-relativistic limit $\eta(v_{\min}) = \int_{DM} \frac{m_{\chi}^2}{m_{\chi}^2 v} = \int_{DM} \frac{1}{v}$

boosted DM energy spectrum

$$p_{\chi}^{\min} = \frac{q}{2(1 - \Delta E_e^2/q^2)} \left(1 - \frac{\Delta E_e^2}{q^2} \pm \frac{\Delta E_e}{q} \sqrt{\left(1 - \frac{\Delta E_e^2}{q^2}\right) \left(1 + \frac{4m_{\chi}^2}{q^2} - \frac{\Delta E_e^2}{q^2}\right)} \right)$$

Spectrum of electron recoil energy in xenon

- Ionization from factor

- ◆ outgoing electron as a free plane wave.

$$|f_{ion}^{nl}(k', q)|^2 = \frac{2k'^3}{(2\pi)^3} \sum_{\substack{nlm, \\ n'l'm'}} |\langle n'l'm' | e^{iqx} | nlm \rangle|^2$$

J. Kopp, V. Niro, T. Schwetz and J. Zupan, arXiv:0907.3159

$$|f_{ion}^{nl}(k', q)|^2 = \frac{2k'^3}{(2\pi)^3} \sum_{\text{deg}} |f_{nl}(q)|^2 = \frac{(2l+1)k'^2}{4\pi^3 q} \int_{|k'-q|}^{|k'+q|} k dk \sum_{m=-l}^l |\chi_{nl}(k)|^2$$

- Roothaan Hartree Fock (RHF) method

$$\chi_{nl}(p) = 4\pi i^l \int dr r^2 R_{nl}(r) j_l(pr)$$

- ◆ Radial wave functions: linear combination of Slater type orbitals

C. F. Bunge, J. A. Barrientos and A. V. Bunge (1993)

$$R_{nl}(r) = \sum_k C_{nlk} \frac{(2Z_{lk})^{n_{lk}+1/2}}{a_0^{3/2} \sqrt{(2n_{lk})!}} (r/a_0)^{n_{lk}-1} \exp\left(-\frac{Z_{lk}r}{a_0}\right)$$

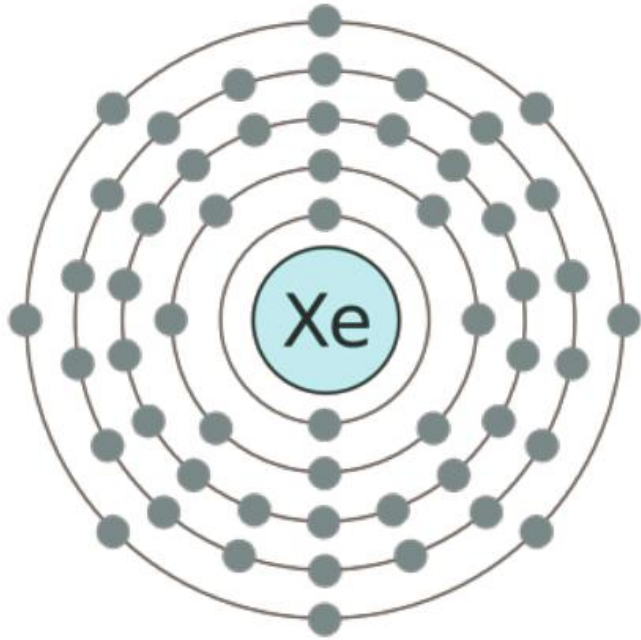
J. Kopp, V. Niro, T. Schwetz and J. Zupan, arXiv:0907.3159

- ◆ Analytical expression

$$\chi_{nl}(p) = \sum_k C_{nlk} 2^{n_{lk}-1} \left(\frac{2\pi a_0}{Z_{lk}}\right)^{3/2} \left(\frac{ipa_0}{Z_{lk}}\right)^l \frac{\Gamma(n_{lk} + l + 2)}{\Gamma(l + \frac{3}{2}) \sqrt{(2n_{lk})!}} \times {}_2F_1 \left[\frac{1}{2}(n_{lk} + l + 2), \frac{1}{2}(n_{lk} + l + 3), l + \frac{3}{2}, -\left(\frac{pa_0}{Z_{lk}}\right)^2 \right]$$

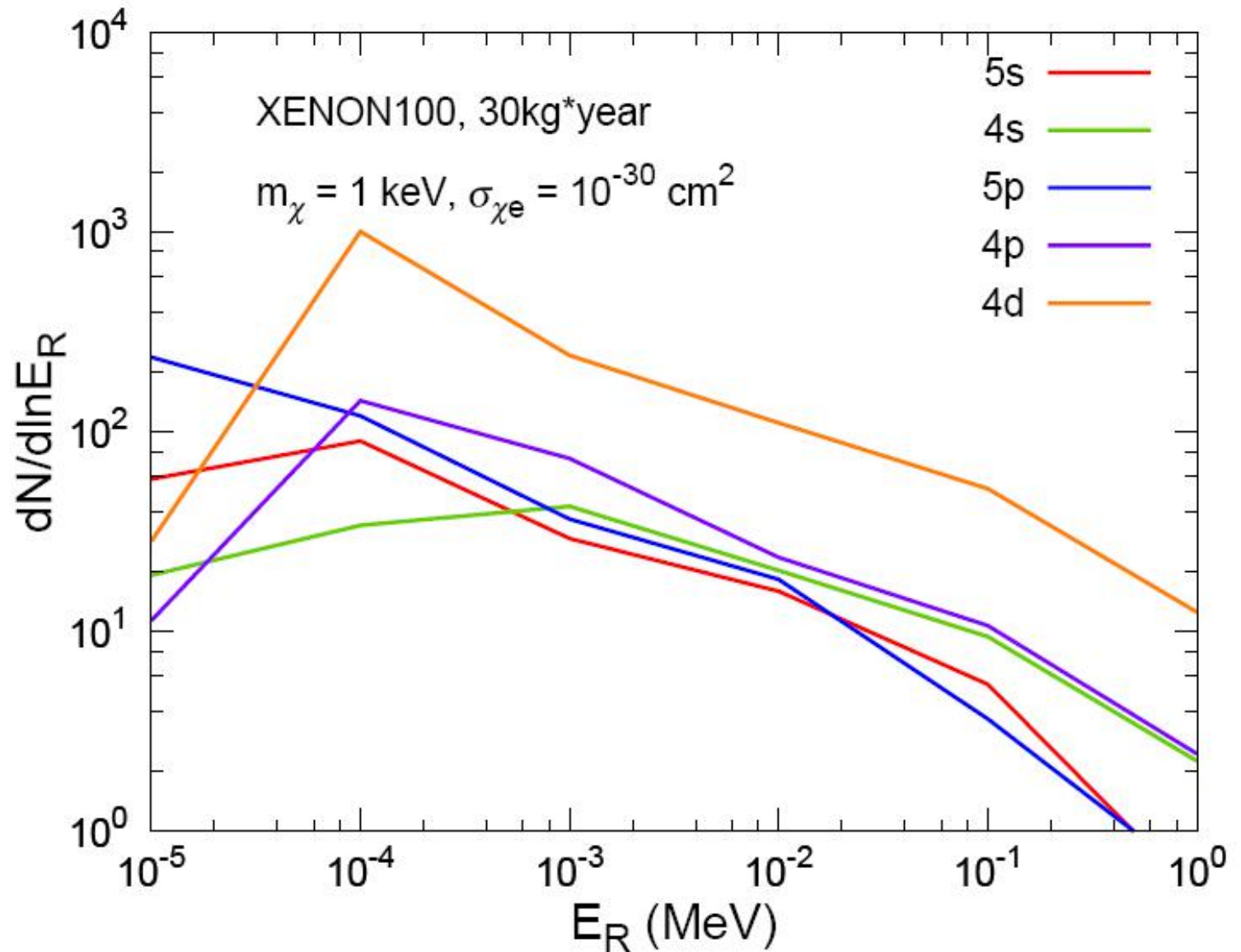
Spectrum of electron recoil energy in xenon

$$\frac{dR_{ion}}{d \ln E_R} = N_T \frac{\rho_\chi}{m_\chi} \frac{d \langle \sigma_{ion}^{nl} v \rangle}{d \ln E_R}$$

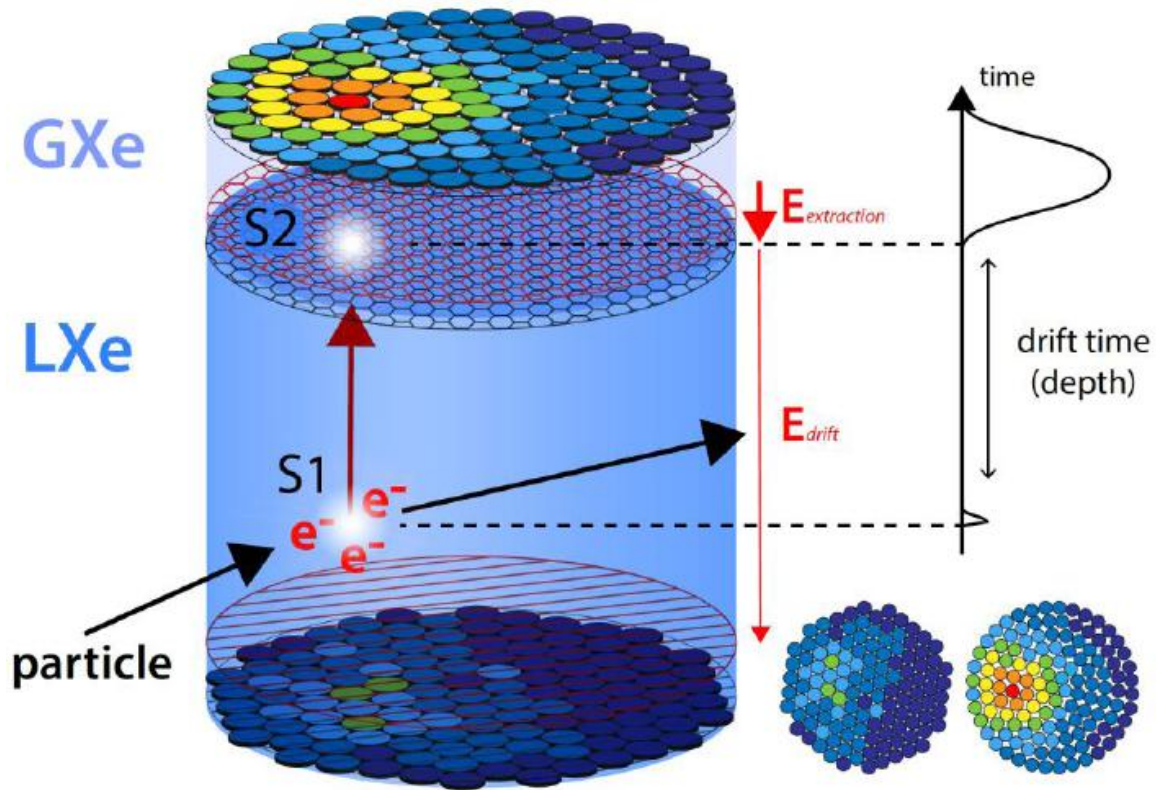


Electron configuration:

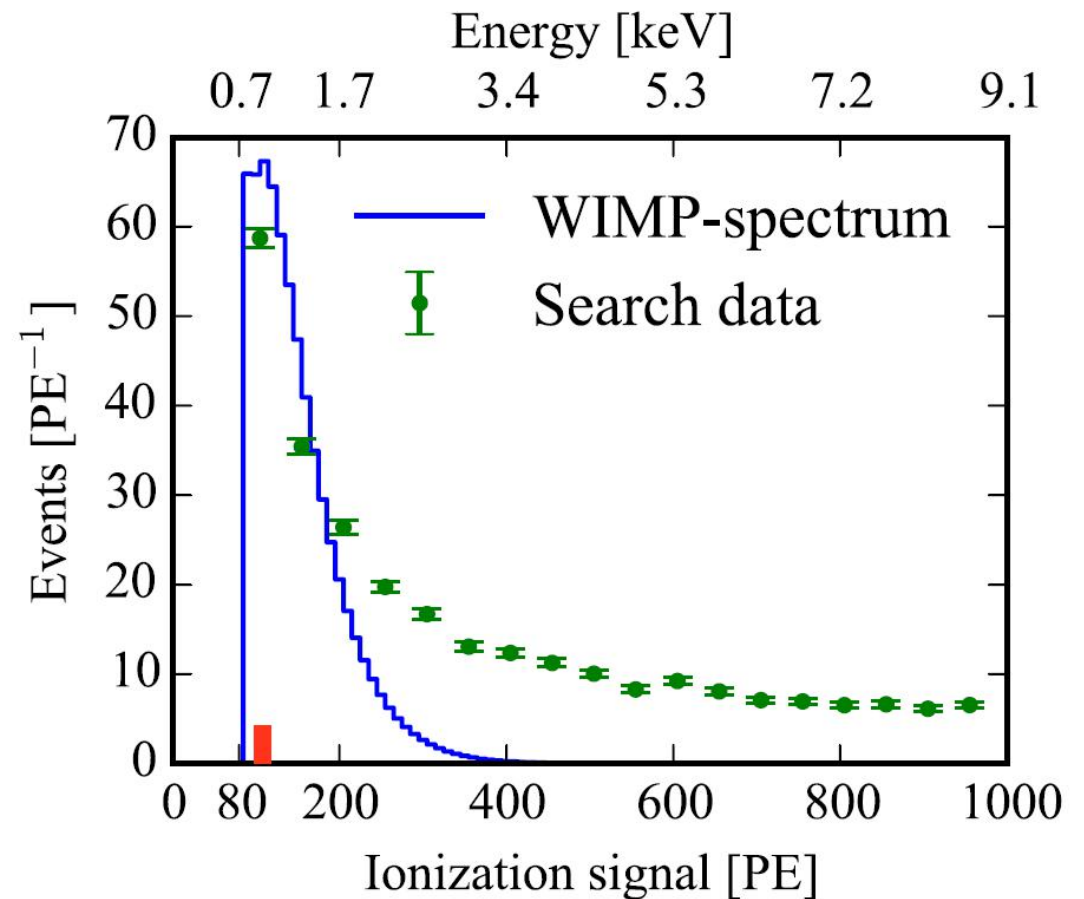
[Ar] 3d¹⁰ 4s² 4p⁶ 4d¹⁰ 5s² 5p⁶



Spectrum of photoelectrons (PE) in xenon



- S1 light signal: Prompt scintillation photons
- S2 Charge signal: Secondary scintillation photons from electroluminescence in Gxe due to drifted electrons



Spectrum of photoelectrons (PE) in xenon

- Event rate for Si responses

$$\frac{dR}{dS_i} = \int dE_R P(S_i|E_R) \frac{dR_{ion}}{d \ln E_R}$$

- Probability to produce S2 for given recoil energy

$$P(n_e|\langle n_e \rangle) = \text{binom}(n_e|N_Q, f_e)$$

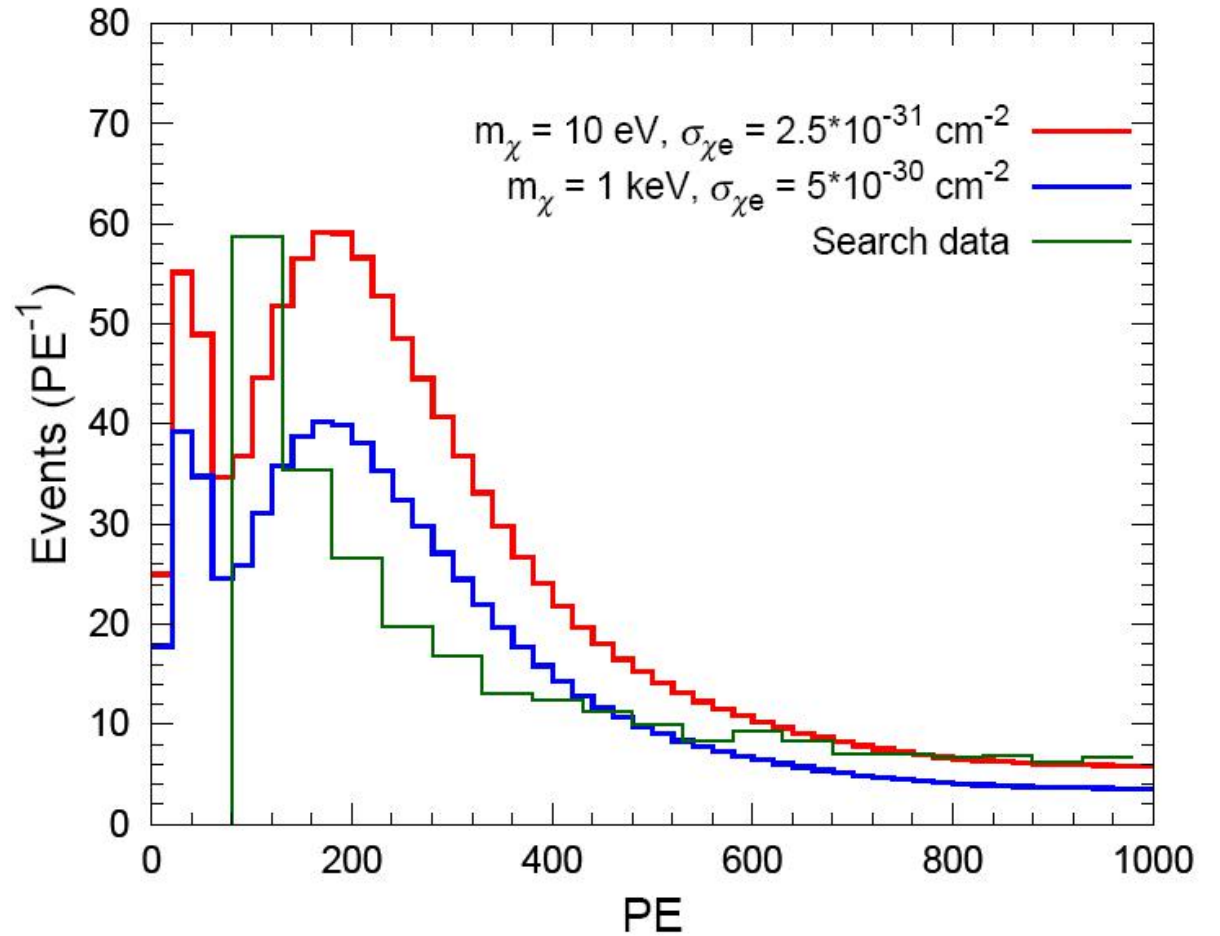
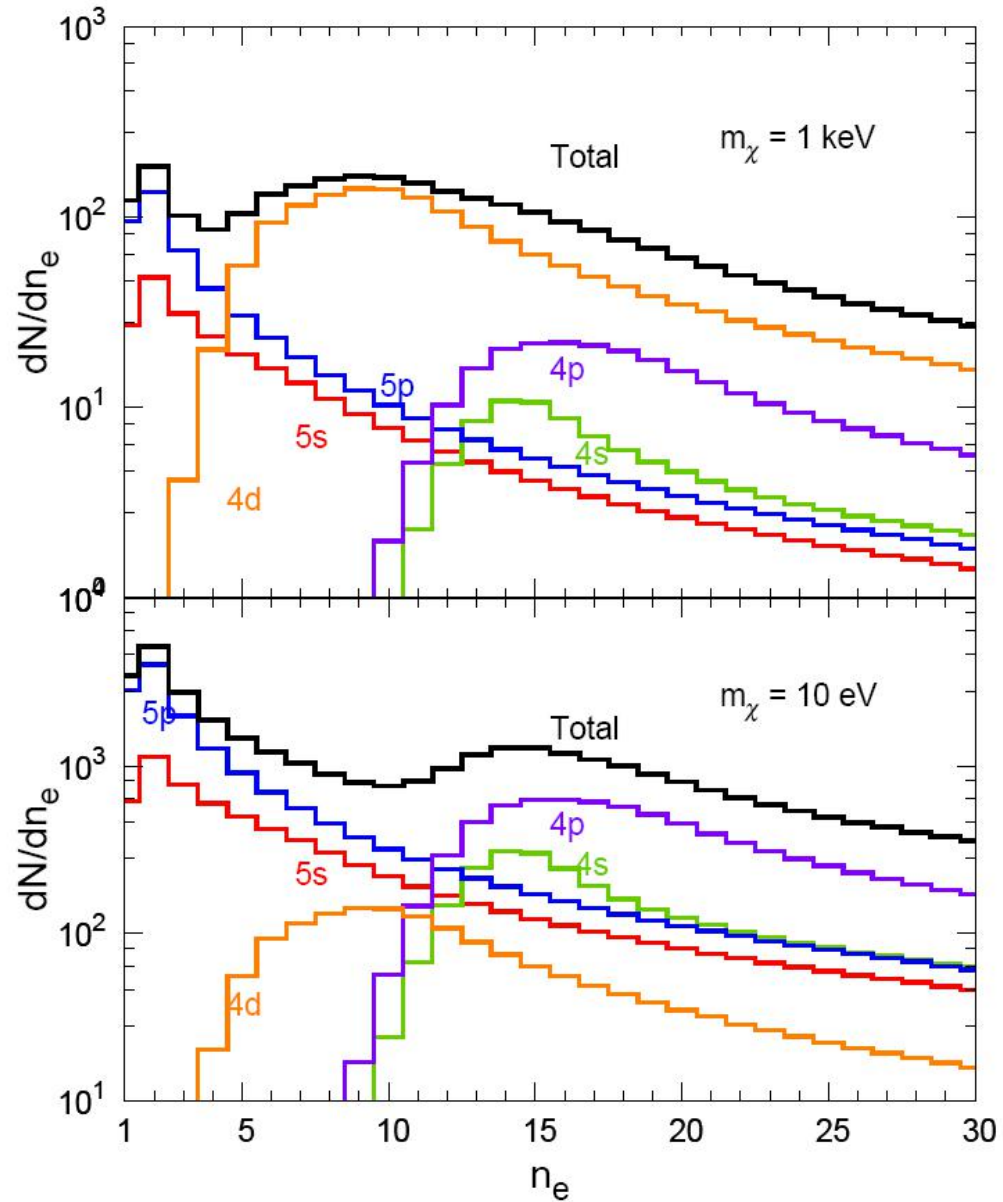
$$N_Q = n_e + n_\gamma = N_i + N_{ex}$$

$$f_e = (1 - f_R)(1 + N_{ex}/N_i)^{-1} \approx 0.83$$

$$P(n_e^{\text{surv}}|n_e) \approx 1, \quad P(S2|n_e^{\text{surv}}) = \text{gauss}(S2|n_e^{\text{surv}}, \sigma_{S2})$$

$$P(S2|E_R) = \sum_{n_e^{\text{surv}}} \sum_{n_e} P(S2|n_e^{\text{surv}}) P(n_e^{\text{surv}}|n_e) P(n_e|\langle n_e \rangle)$$

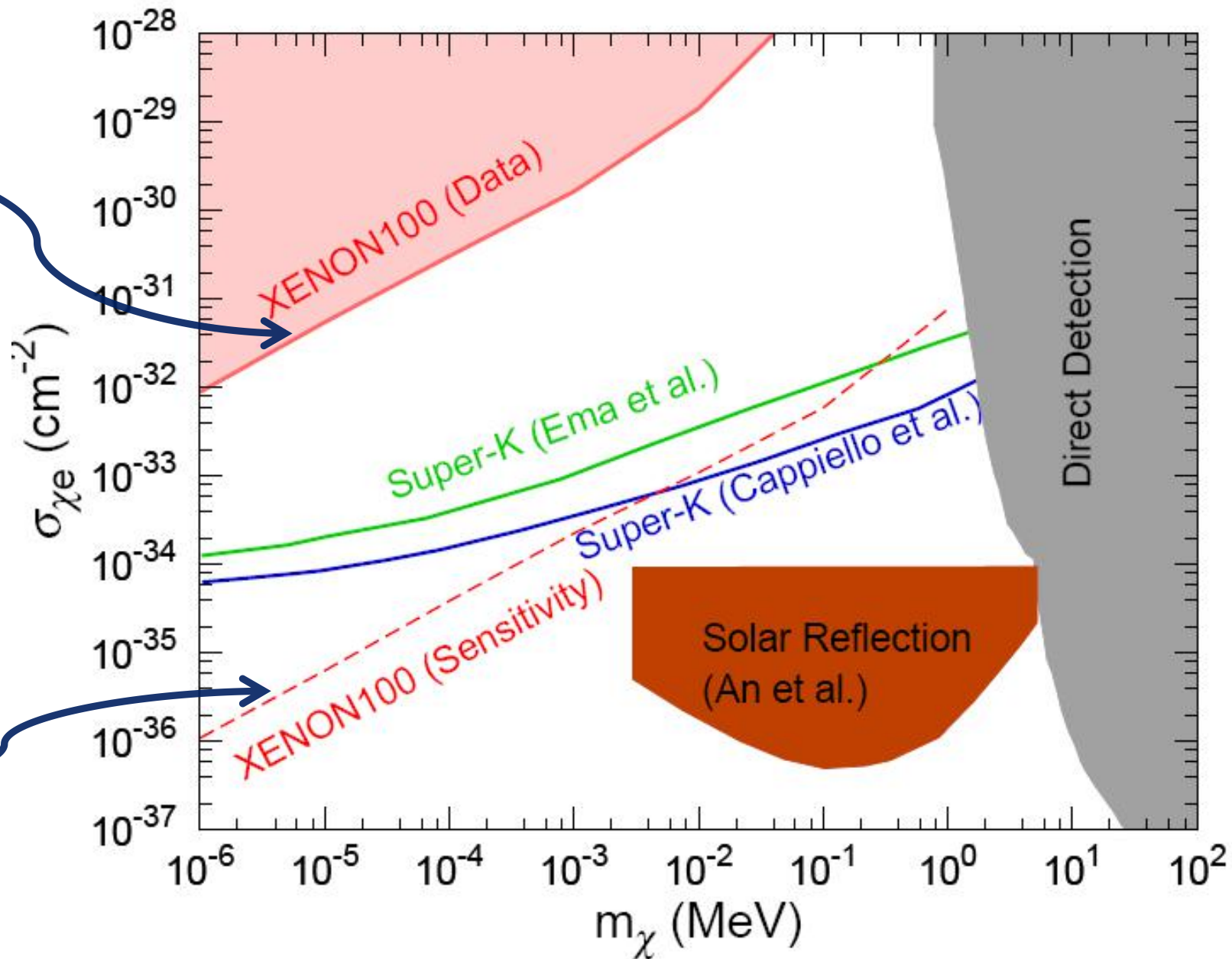
Spectrum of photoelectrons (PE) in xenon



Limit on the DM-electron scattering cross section

$$\frac{dR}{dS_i} = \int dE_R P(S_i|E_R) \frac{dR_{ion}}{d \ln E_R}$$

$$\left. \frac{dR}{dS_i} \right|_{\text{sat}} \simeq \int_{E_R > \text{KeV}} dE_R \frac{dR_{ion}}{d \ln E_R}$$



谢谢大家