

Coherent Axion Production through Laser Crystal Interaction

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OUTLINE

1. Introduction
2. Coherent Axion Production
3. Layer Stacking and Conversion Probability
4. Experimental Design and Result Estimation
5. Summary and Discussion

1.1 Axion and Axion-Like-Particles (ALP)

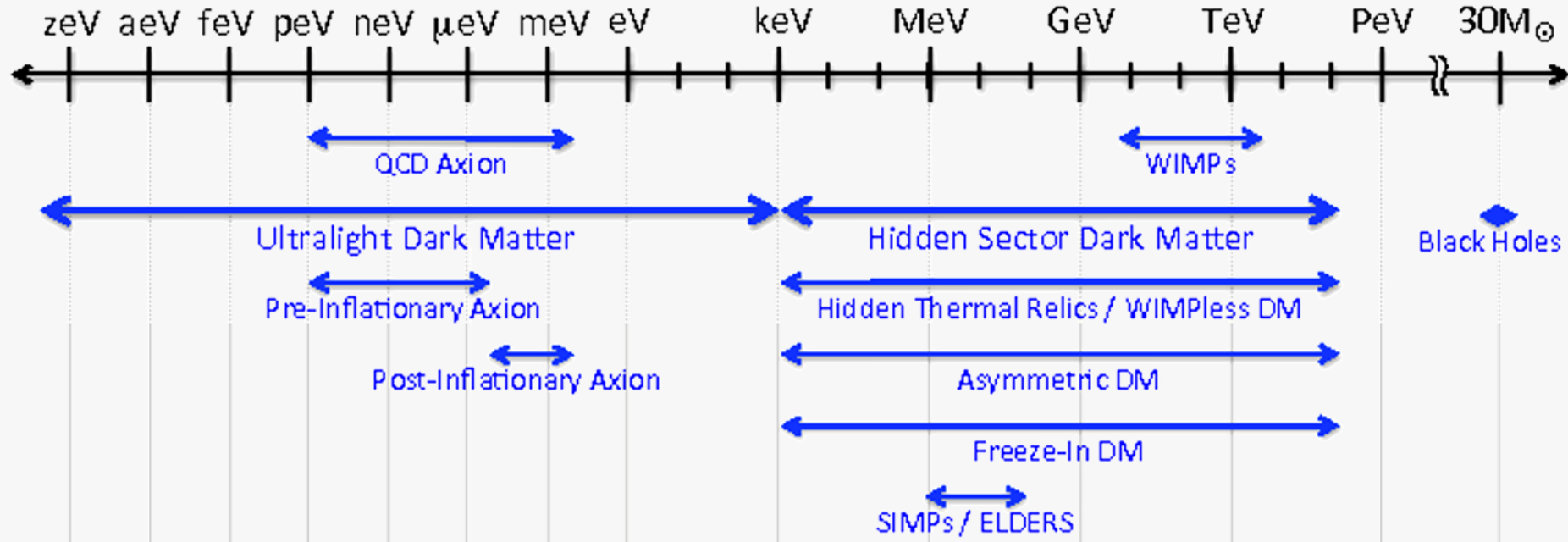


Fig: Some dark matter candidates and their mass range.

- Dark matter explains a lot of problems in cosmology and astrology, but we don't know its nature yet.
- **Axion** solve strong CP problem and dark matter problem simultaneously, making it a hot candidate.
- In researches, axion-like-particles (ALP) are often considered, which have wider parameter space.

1.1 Axion and Axion-Like-Particles (ALP)

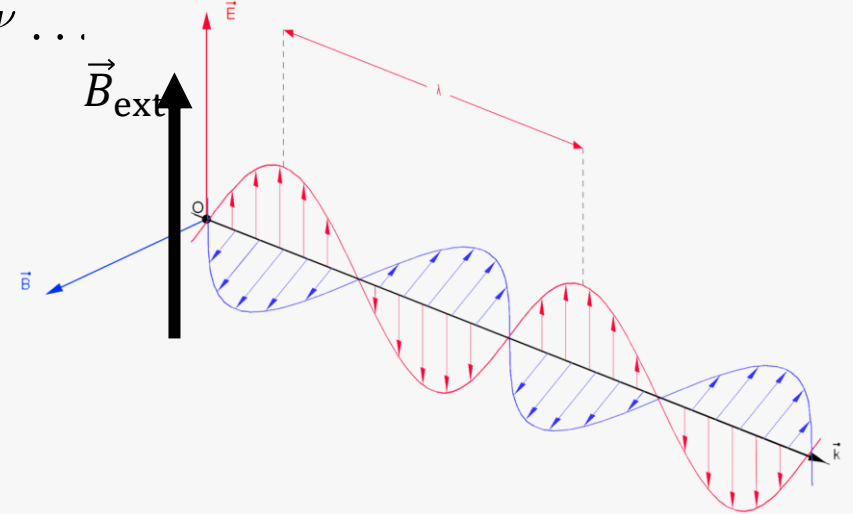
- Axion couples to ordinary matter. We consider its coupling to photon here.

$$\mathcal{L}_{\text{axion}} \supset \underbrace{\frac{1}{4}g_{a\gamma\gamma}aF_{\mu\nu}\tilde{F}^{\mu\nu}}_{\text{photon channel}} + \underbrace{g_{aee}a\bar{e}i\gamma_5 e}_{\text{electron channel}} + \underbrace{4\pi g_{aGG}aG_{\mu\nu}\tilde{G}^{\mu\nu}}_{\text{gluon channel}} \dots$$

photon
channel

electron
channel

gluon
channel



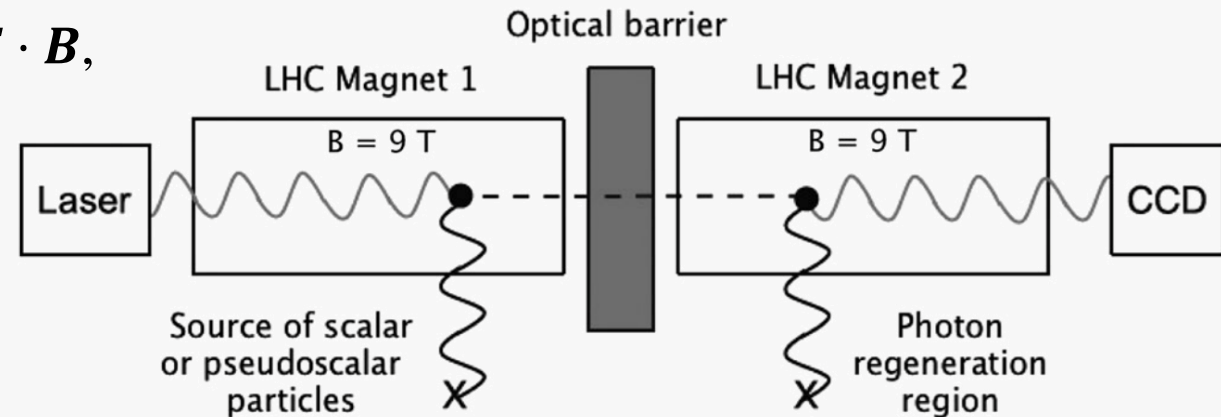
- Field equation for axion:

$$(\partial_t^2 - \nabla^2 + m_a^2) a = g_{a\gamma\gamma} \mathbf{E} \cdot \mathbf{B}$$

We can use laser and magnetic field to get nonzero $\mathbf{E} \cdot \mathbf{B}$,
and generate axion.

(Light-Shining-Through-Wall experiment, LSW)

- Conversion probability: $P_{a \leftrightarrow \gamma} \approx \frac{1}{4} g_{a\gamma\gamma}^2 (BL)^2$



1.2.1 CERN Axion Solar Telescope (CAST)

- In sun, photon interact with nucleus and generate axion, whose energy corresponds to X-ray.
- Axion arrive on earth, traverse magnetic field and converts to photon, and then be detected.
- Same mechanism as light-shining-through-wall experiments.

● How to increase magnetic field?

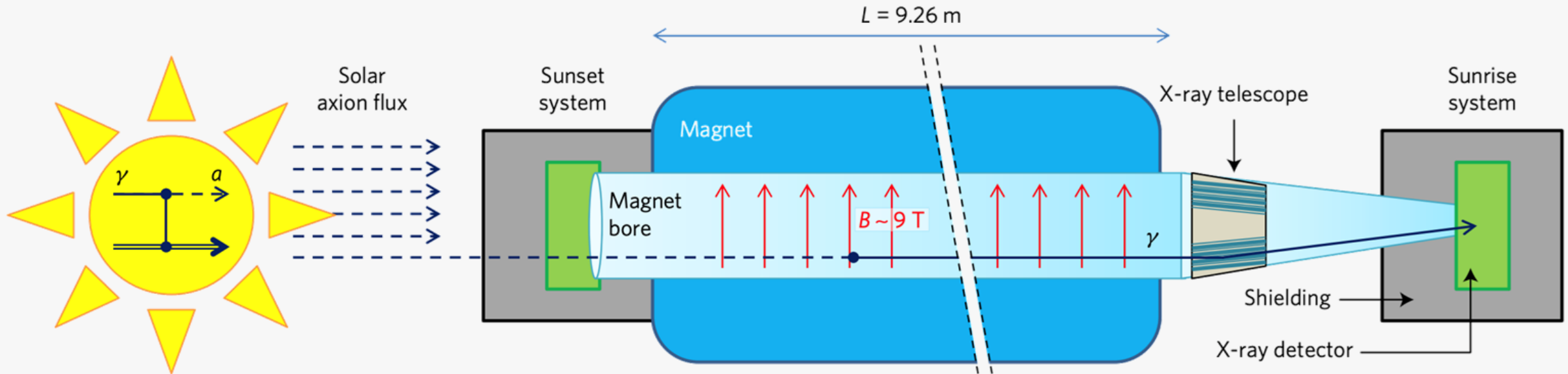


Fig: Experiment design for CAST. Taken from [3].

1.2.2 Cryogenic Dark Matter Search, CDMS

- Strong electric field in crystal can be used for axion detection.
- When axion wavelength match the crystal lattice constant, there is coherent enhancement✧.
- Solar axion interacts with crystal for detection.

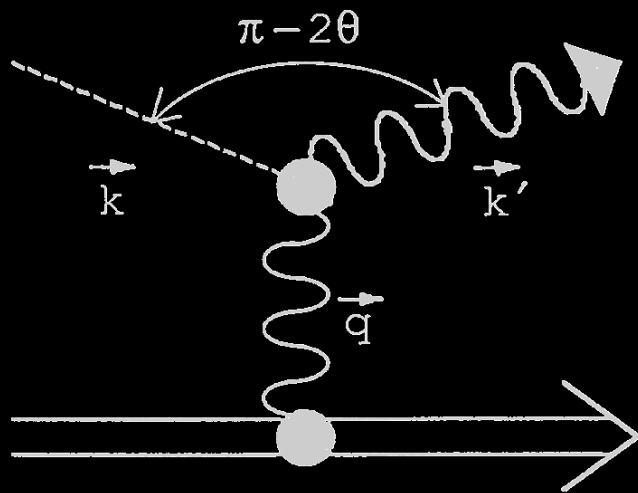


Fig: Axion interact with nuclei[4].

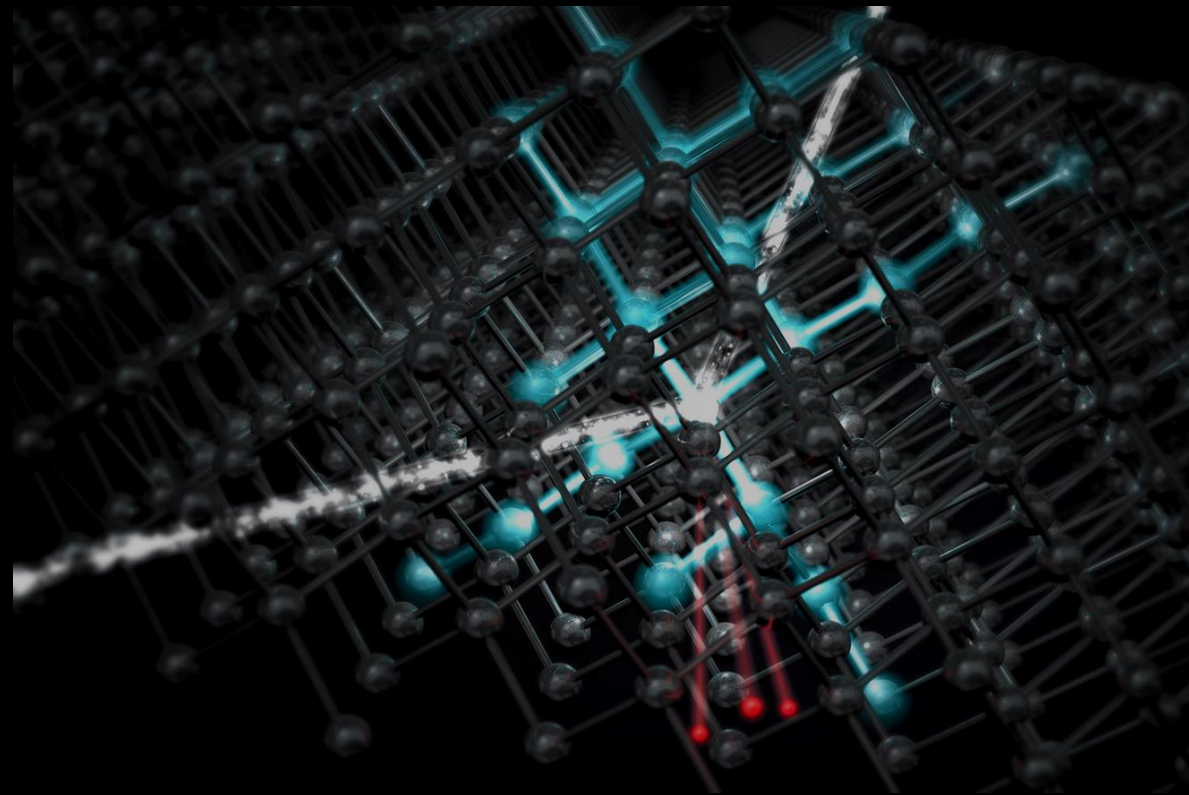


Fig: Conceptual figure for CDMS

[4] R. J. Creswick, et al., PLB 427, 235(1998)

✧ Detailed explanation later

1.3 Axion Coupling Constraints

- Indirect search, direct space axion search and direct lab axion search are complement to each other.
- Currently, terrestrial experiments give much **weaker** constraints than other two approaches.
- We want to propose new schemes for axion production.

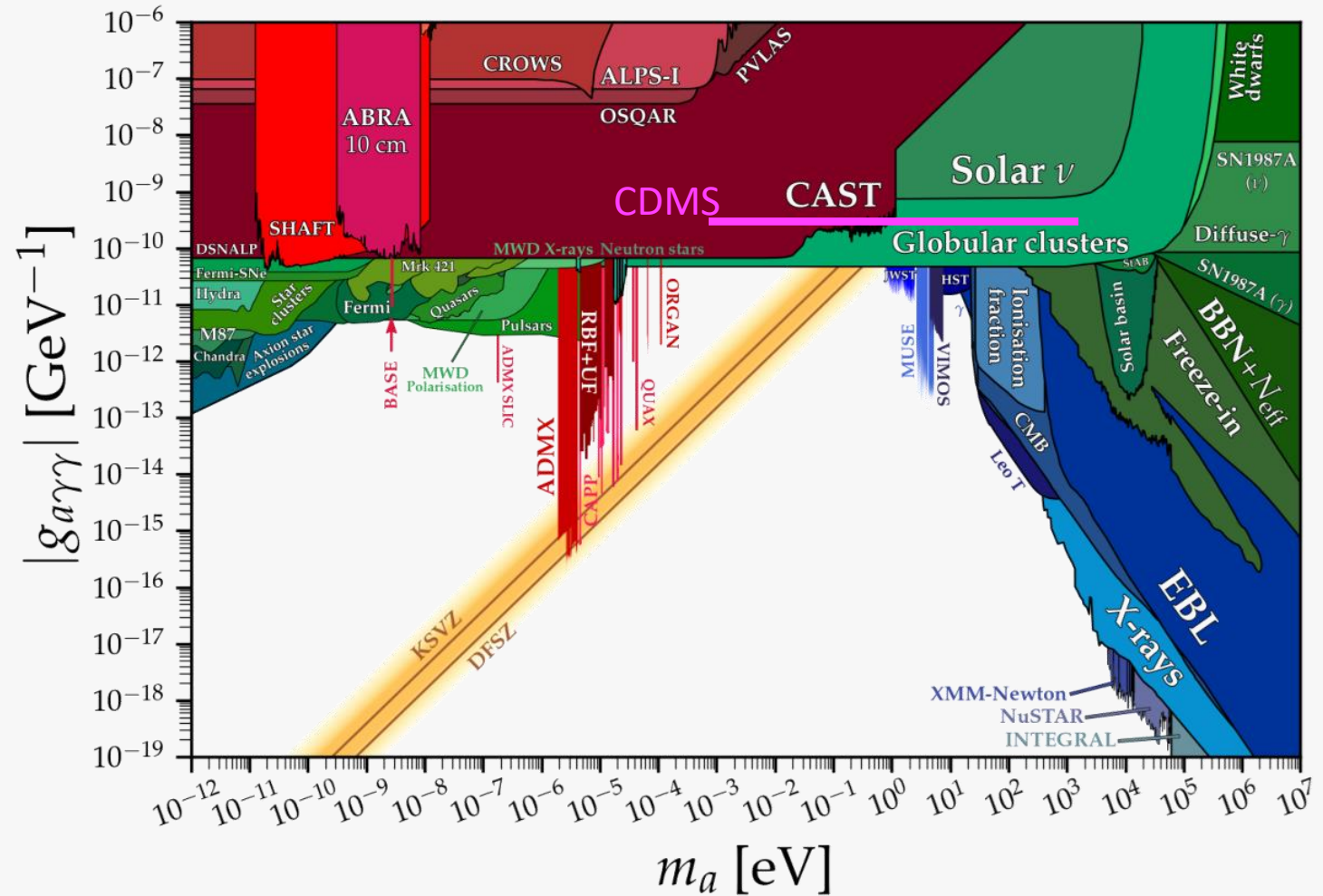


Fig: Constraints for axion parameters from different studies.
Taken from [5].

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2. Coherent Axion Production

2.1 Bragg Scattering

- X-ray can be scattered by crystal lattice. When Bragg condition, $n\lambda = 2d\sin\theta$, is satisfied, the scattered waves from different lattice undergo coherent superposition.

- A more general condition is:

$$\mathbf{k}_{\text{out}} - \mathbf{k}_{\text{in}} = \mathbf{G} \equiv \frac{2\pi}{d}(i_1, i_2, i_3)$$

e.g., if $\mathbf{G} = \frac{2\pi}{d}(0,1,0)$, the phase difference between two adjacent lattices in the y-direction is 2π , while the phases of adjacent lattices in the x- and z-directions are identical.

- If \mathbf{k}_{out} represents the momentum for axion, there can be coherent enhancement, same as Bragg scattering.

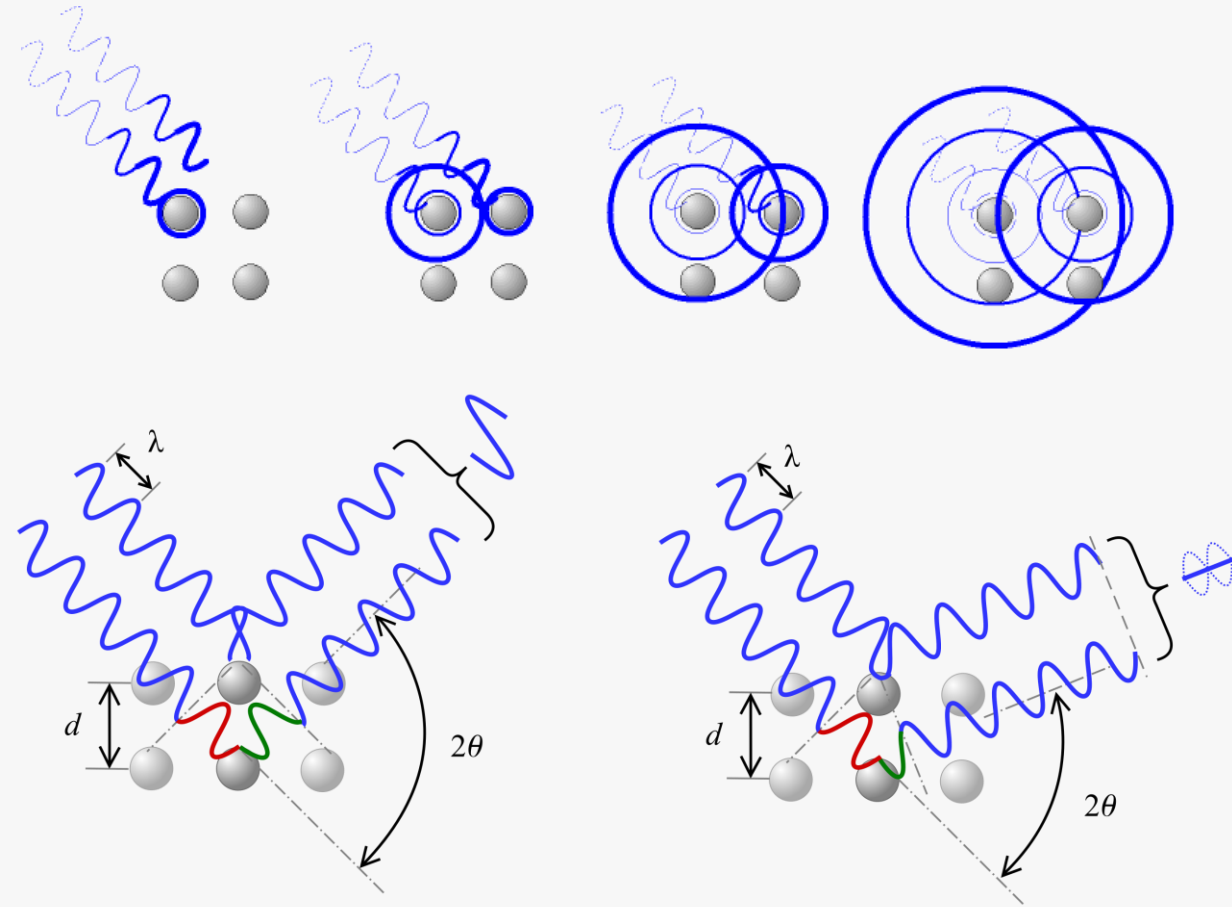
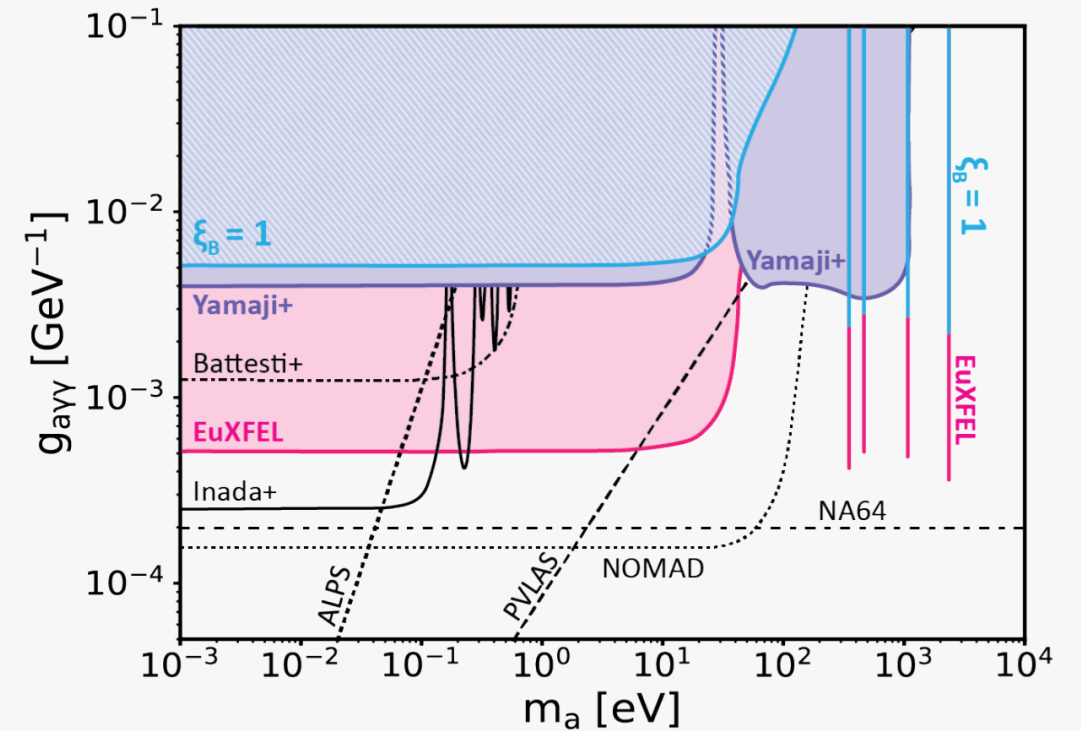
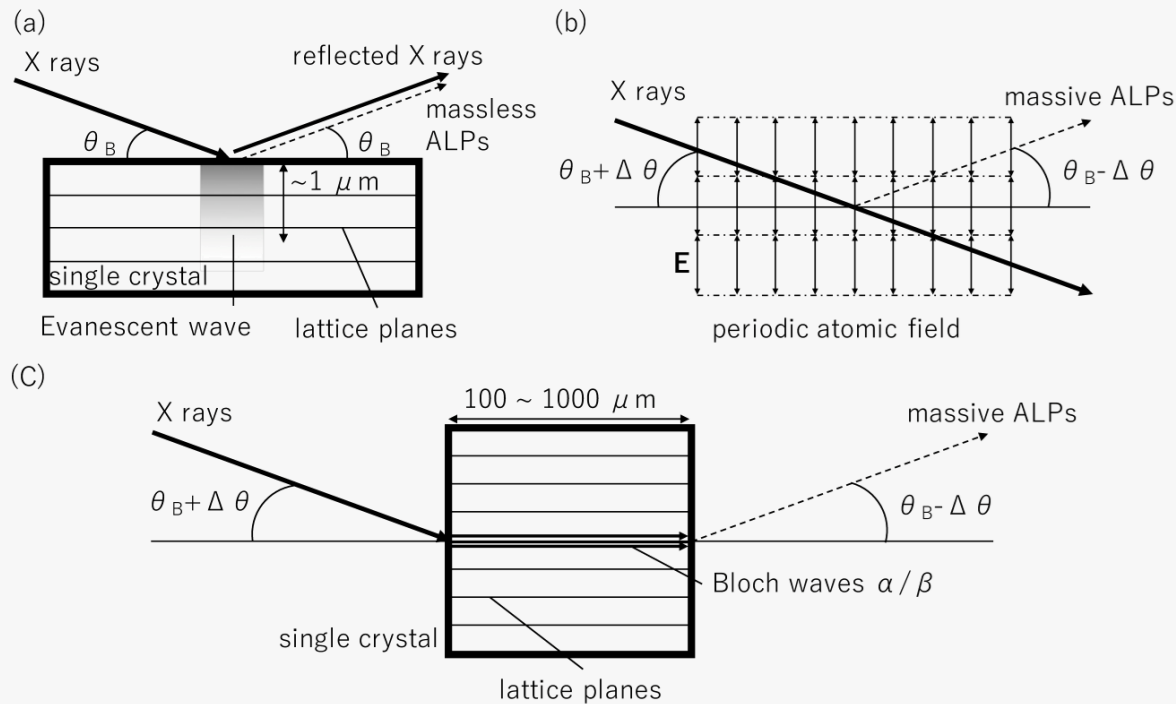


Fig: Schematic figure for Bragg scattering.

2.1 Bragg Scattering—Axion

- When X-ray satisfies Bragg condition, the axion production will enhance coherently.
- Due to reflection or absorption, X-ray only enters crystal for milli-meter or micro-meter, much shorter than typical LSW experiments ($O(10\text{m})$)
- This approach give weaker constraints on axion parameters.



Left: Schemes for axion production with X-ray. Taken from PRD96, 115001(2017);
 Right: Estimated exclusion line with EuXFEL parameters. Taken from PRL134, 17333(2025).

2.2 Optical Light and Ionic Crystal

- Due to reflection or absorption, X-ray only enters crystal for milli-meter or micro-meter
- For larger interaction length, we consider optical light and transparent crystal. Optical laser also has much larger photon number.
- Optical light cannot identify electric field inside atom, we choose ionic crystal* (CaF₂ as example) .
- Optical light energy is low, we can only choose:

$$\mathbf{k}_a - \mathbf{k}_L = \mathbf{G} \equiv \frac{2\pi}{d} (0, \epsilon, 0)$$

i.e., the phases of two adjacent lattices in the x- and z-directions are identical, while those in the y-direction are approximately identical.

* Ionic crystal is still charge neutral, optical light can only faintly resolve its electric dipole moment.

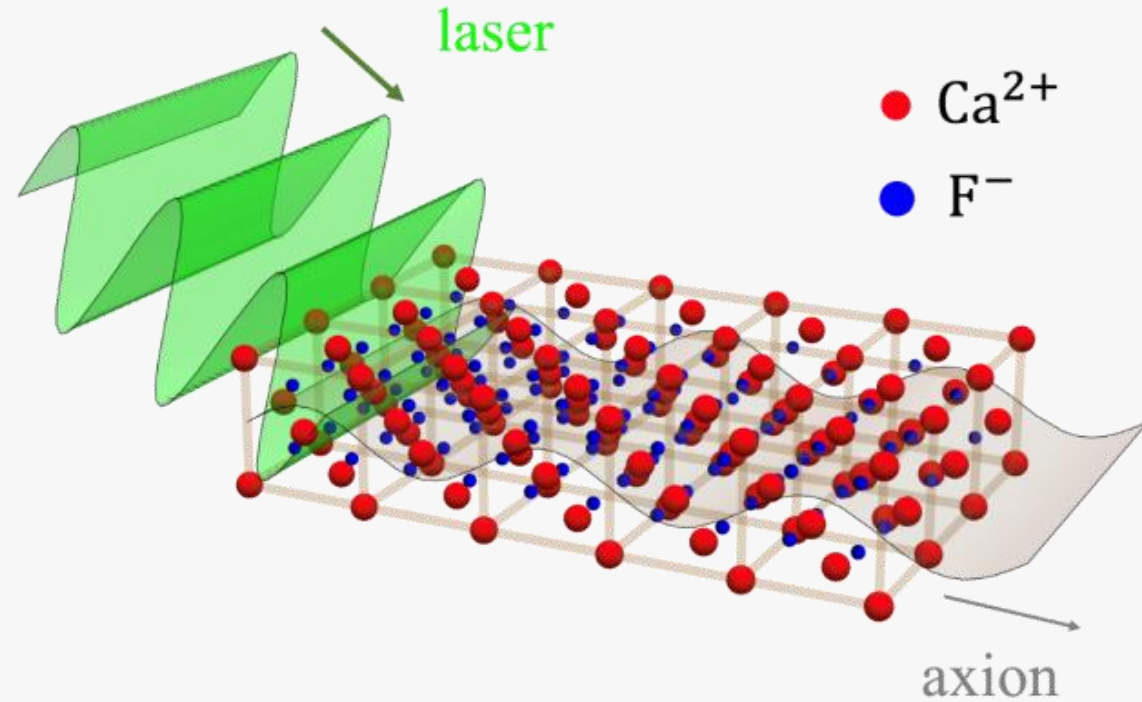


Fig: Axion production through laser-axion interaction

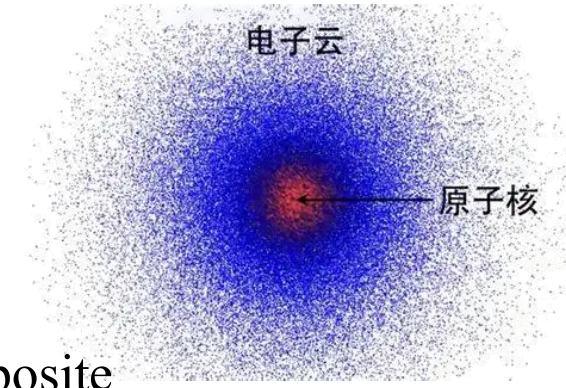
	Optical	X-ray
Interaction length	✓	×
Photon number	✓	×
Spatial resolution	×	✓

TAB: pros and cons for optical light and X-ray

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3.0* Damping of Long Wavelength



- Axion wave function is proportional to $a(t, x) \propto \sum_s Q_s \exp[-i\Delta \mathbf{k} \cdot \mathbf{r}_s]$
- The wavelength λ of optical light is long and it's difficult to distinguish two close opposite charges. The wave function will be damped by order $\mathcal{O}(d/\lambda)$, where d is the distance between the two charges.
- For charge neutral atom, it can be regarded as point-like nucleus and spherical symmetric electric cloud. The axion wave function will be damped by order $\mathcal{O}(R^2/\lambda^2)$, where r is the atom radius.

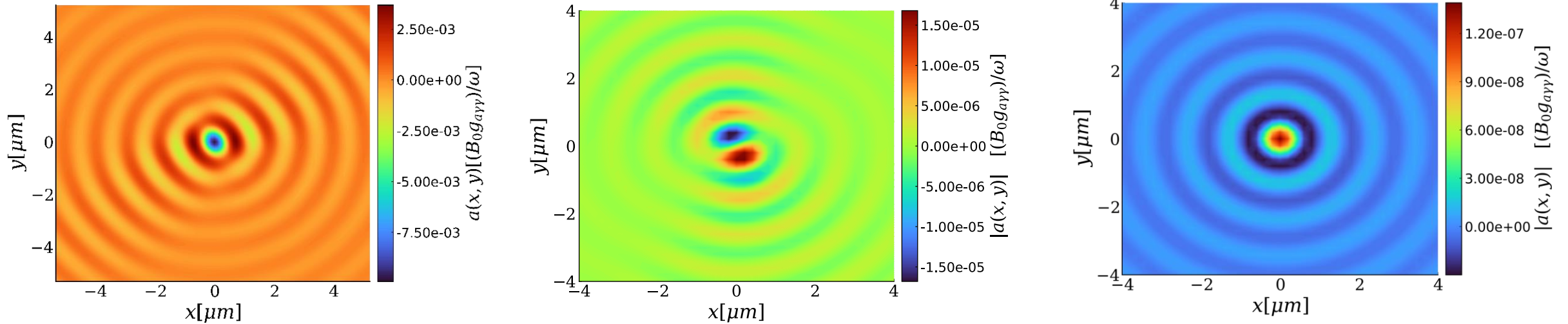


Fig: Axion wave function for laser interacting with ion and atom. Left: single positive charge; Middle: two closely positioned opposite charges; Right: charge-neutral atom.

3.1 Phase Mismatch

- The axion waves from each crystal cell should be coherently superposed.
- The phase of axion waves will shift transversally.
- When the phase shift from two cells reaches π , the axion wave will be coherently destructed.
- The transversal length for π in phase corresponds to half wavelength, which is $\sim 0.5\mu m$ for optical light, which is a very thin film.

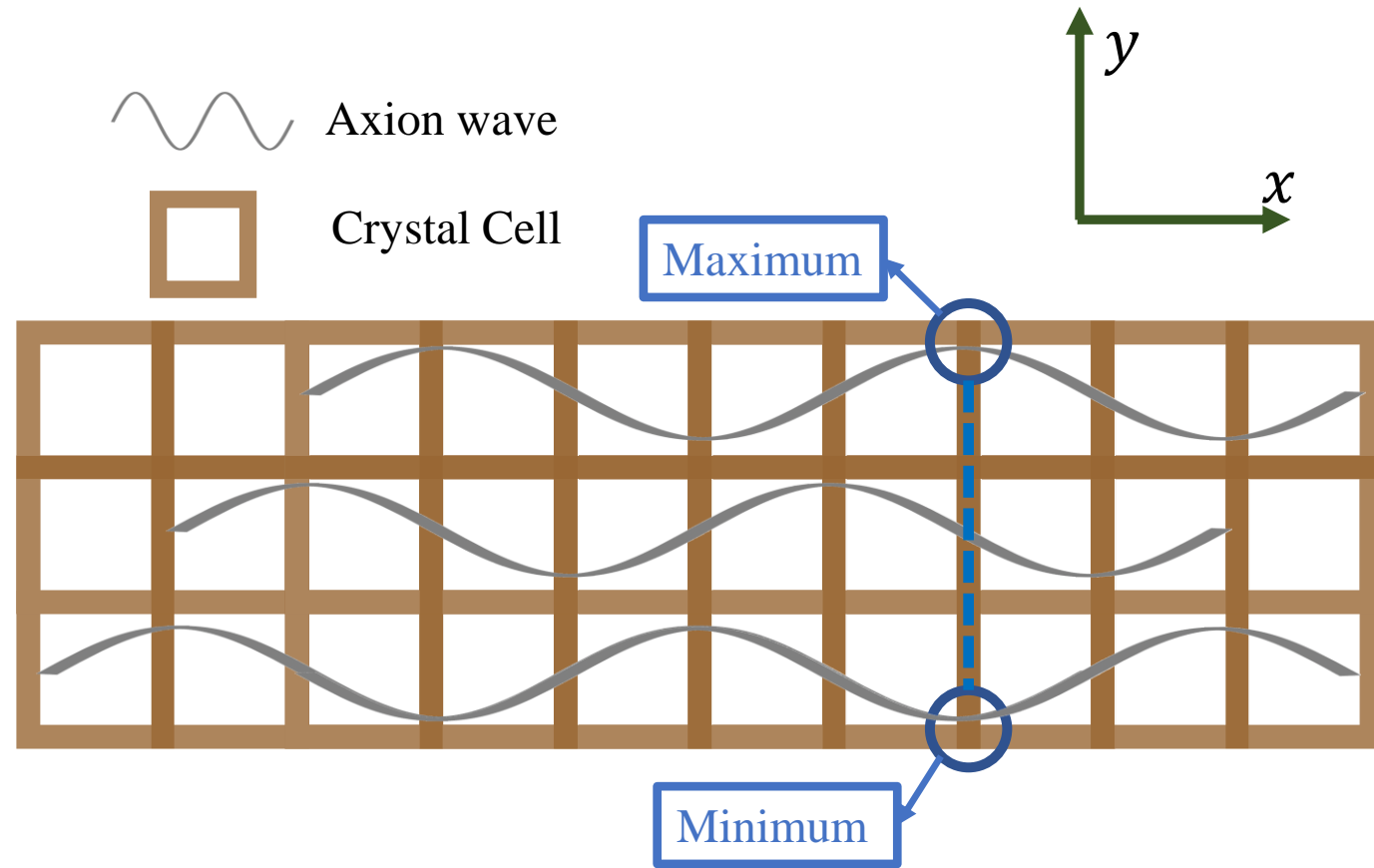


Fig: Schematic figure for coherent destruction on transverse direction.

3.2 Thin Film Stacking

- Stacking films along y-axis. The thickness of each layer is π , and the phase difference between two layers is 2π .
- The laser keep reflecting while propagating along x-axis. This is equivalent to rectangular waveguide in TE mode. The transition probability scales as $P \propto L_x^2$.
- If there are thousands of layers, and if $L_x = 1\text{m}$, $P_{\text{laser} \rightarrow a}^{1\text{m}} = 7.14 \times 10^{-11}$.
If $L_x = 10\text{m}$, $P_{\text{laser} \rightarrow a}^{10\text{m}} = 7.10 \times 10^{-9}$.
- For comparison, in LSW experiment[2], $B = 9\text{T}$, $L = 14.3\text{m}$, and $P^{\text{LSW}} = 4.06 \times 10^{-11}$

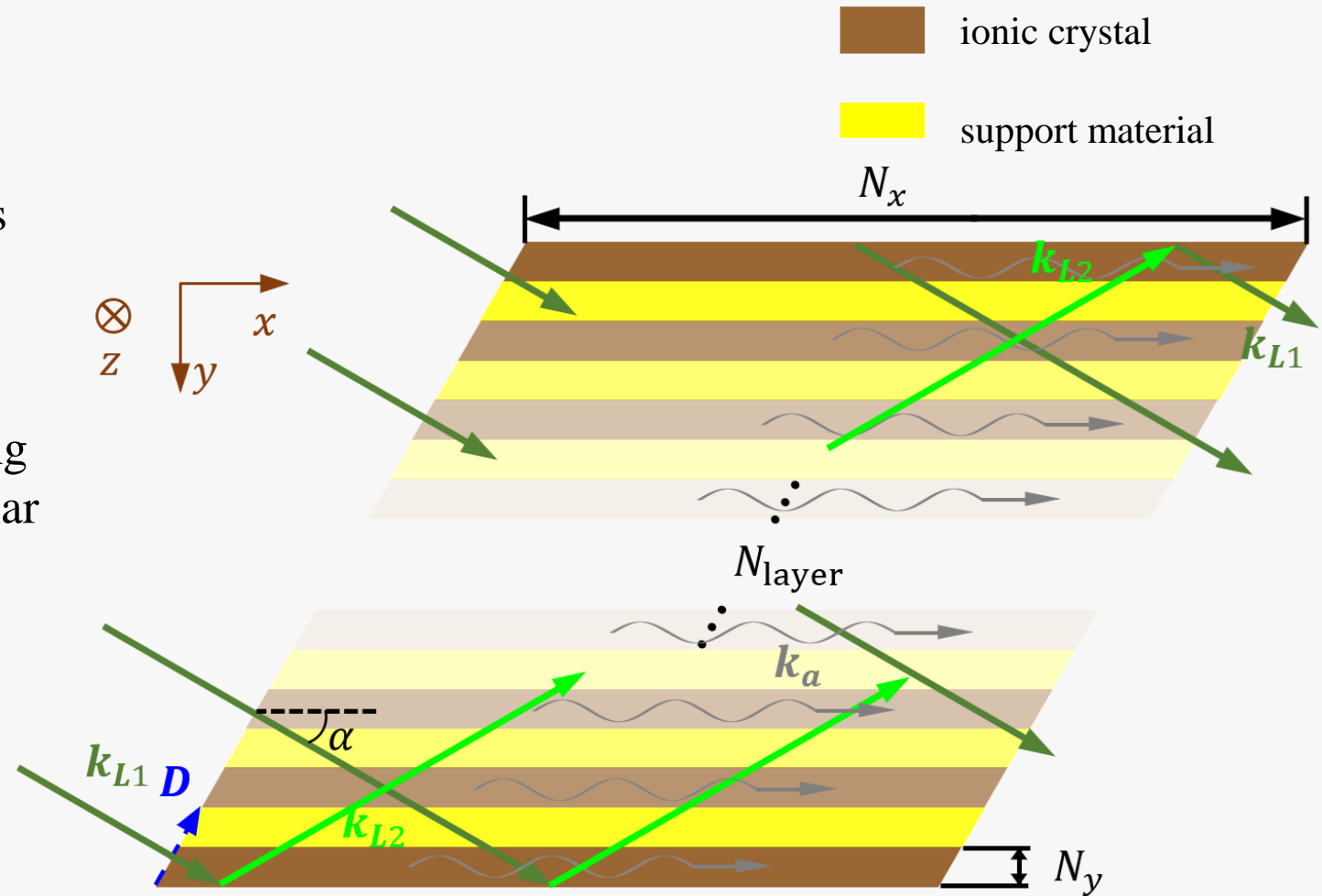
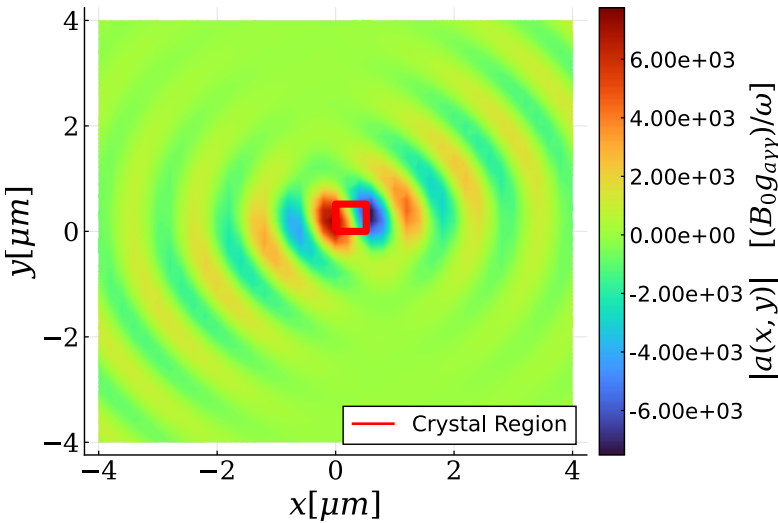
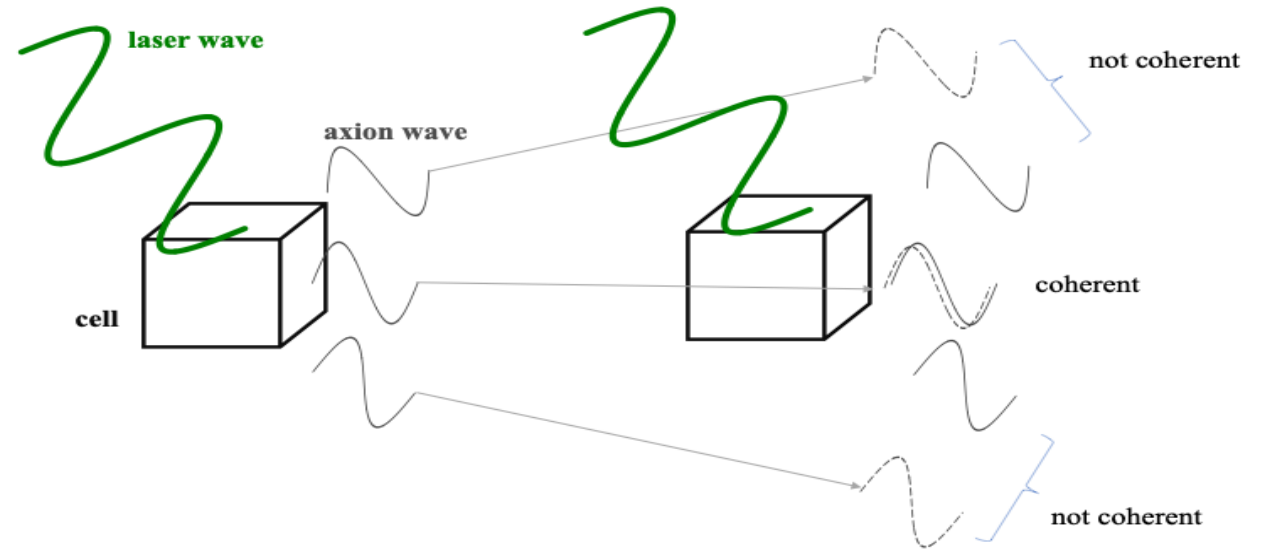


Fig: Schematic figure for thin film stacking. The brown areas are the ionic crystal, while the yellow areas are supporting materials.

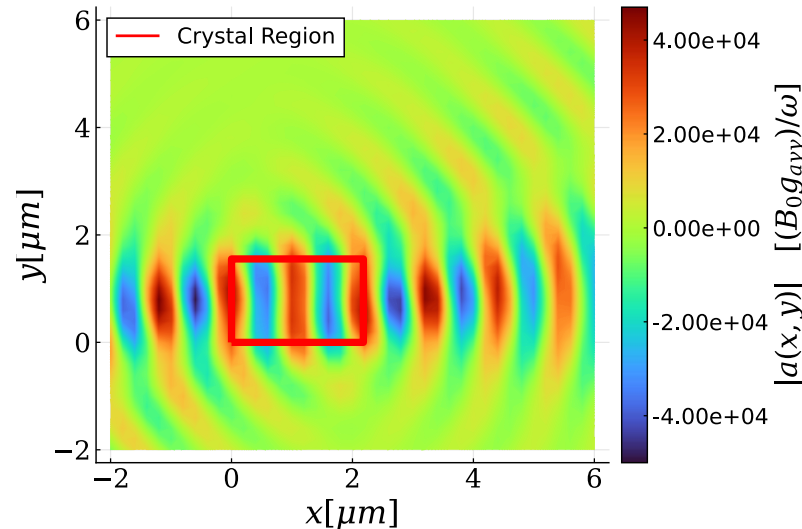
3.3*Longitudinal Coherent

● With the increase of transverse size of crystal, the axion wave approaches plane wave.

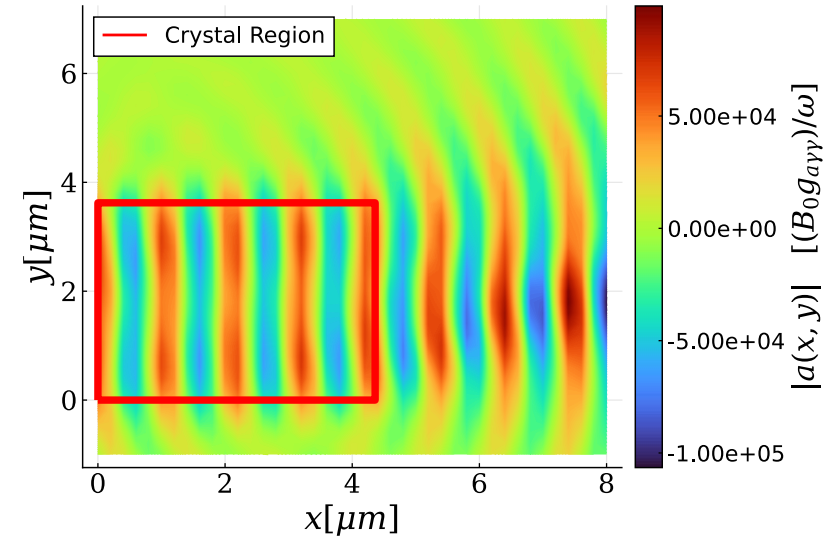
● The axion wave produced at one cell, when propagating to the next cell, if the phase matches, will coherently superposed with axion produced at next cell.



$$N_x = N_y = N_z = 950$$



$$N_x = N_z = 4000, N_{\text{layer}} = 2$$



$$N_x = N_z = 8000, N_{\text{layer}} = 4$$

3.4* Random Distributed Charge

- In crystal, the charges are regularly distributed, while in plasma, the charges are randomly distributed.
- The phase of random charge is undetermined, and the axion wave cannot coherently superposed.

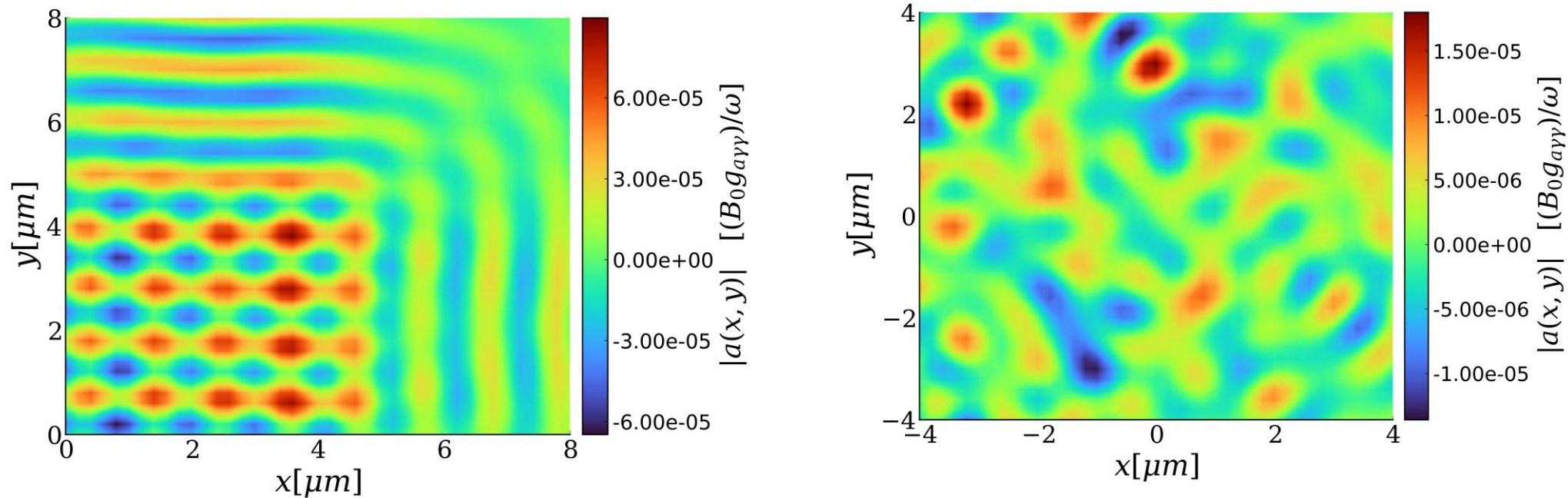


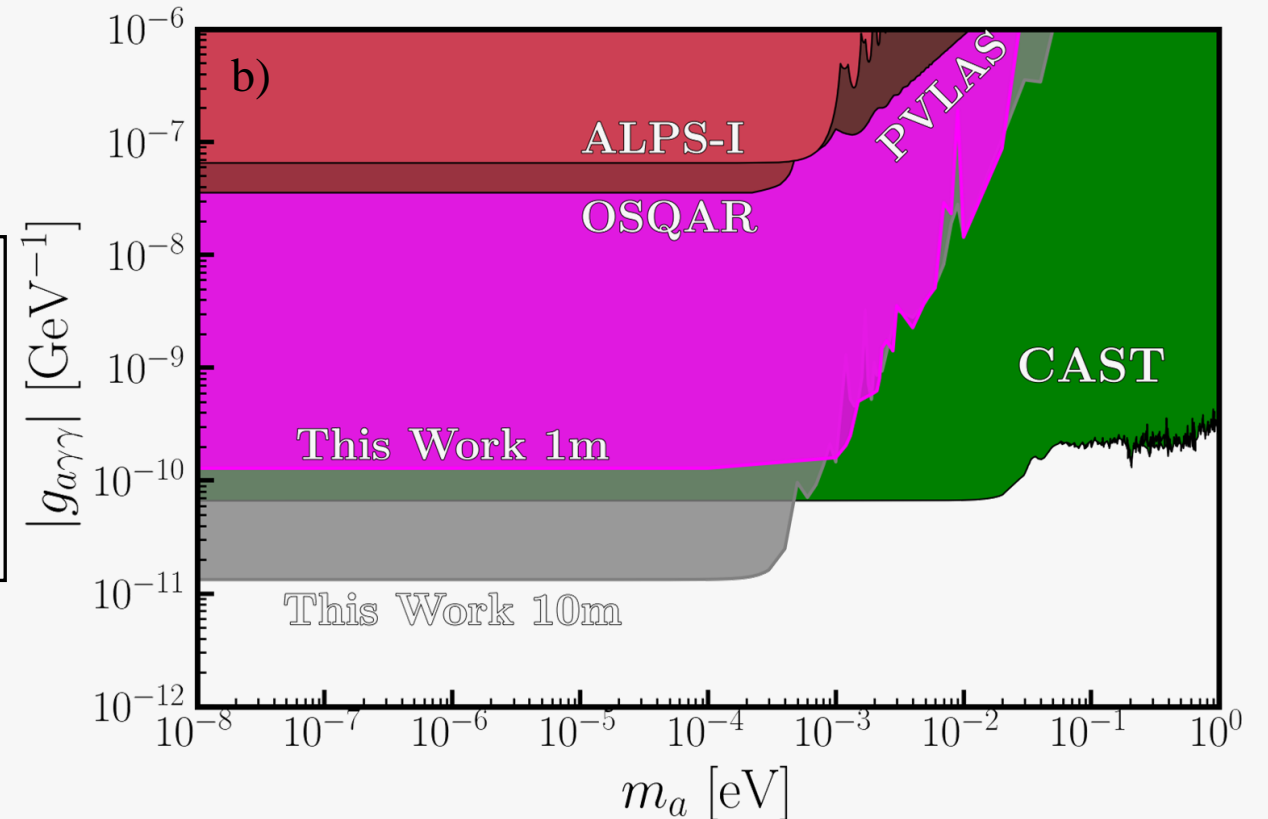
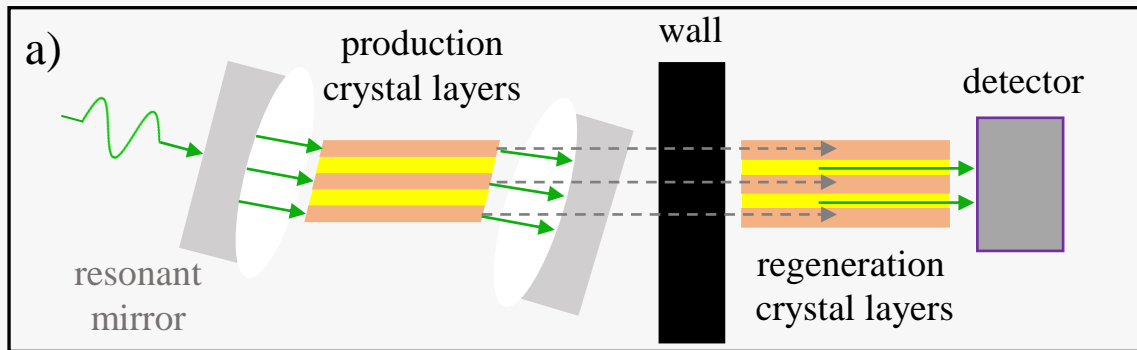
Fig: Axion wave function generated by laser and charge. Left: regularly positioned charge; Right: randomly distributed charge.

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4. Experimental Design

- Our design is similar to LSW: generate axion \rightarrow using wall to block light \rightarrow regenerate light and detect.
- Laser wavelength 1064nm, power 150kW, running for one year 3×10^7 s.
- If the crystal length is 10m, we can constraint the coupling constant to $g_{a\gamma\gamma} \gtrsim 1.32 \times 10^{-11} \text{GeV}^{-1}$, better than CAST.



Left: Experimental design. Right: Estimated exclusion line based on parameters.

5. Summary

- We propose a novel way to generate axion using optical light and ionic crystal.
- We find that, if the laser is injected into the crystal at some specific angle, there can be a coherent enhancement for axion production.
- By stacking thin film layers with specific distance, the coherence can be further enhanced.
- In this way, if the size of the crystal is as large as the magnetic field in LSW experiments, the conversion probability can be 2-order-of-magnitude larger.
- For crystal of 10 meters long, and after one year of data collection, we can obtain an exclusion line better than all current terrestrial experiments.

Appendix. Coherent Enhancement

- The transition probability is:

$$P_{\text{laser} \rightarrow a} \propto \frac{g_{a\gamma\gamma}^2}{S} \int_0^\pi d\theta \int_0^{2\pi} d\phi P_{\text{single}} \times \underbrace{|\mathcal{T}_{\text{lat}}(\Delta \mathbf{k}; \{\mathbf{r}_s\})|^2}_{\text{Multiple cell superposition} \star}$$

Single cell contribution

$$\mathcal{T}_{\text{lat}} = \sum_l \exp(-i\Delta \mathbf{k} \cdot \mathbf{r}_l)$$

- When $\Delta \mathbf{k} \parallel \mathbf{y}$, there are coherent enhancement.
- At the peak of the coherence:

$$\mathcal{T}_{\text{lat}}^{\text{max}} = \frac{1 - \exp[iN_y \omega d \cot \alpha]}{1 - \exp[i\omega d \cot \alpha]} N_x N_z$$

- when $N_y \omega d \cot \alpha = \pi$, the coherence is maximum. For CaF_2 , this corresponds to $L_y = N_y d = 518\text{nm}$. This is a very **thin layer**.

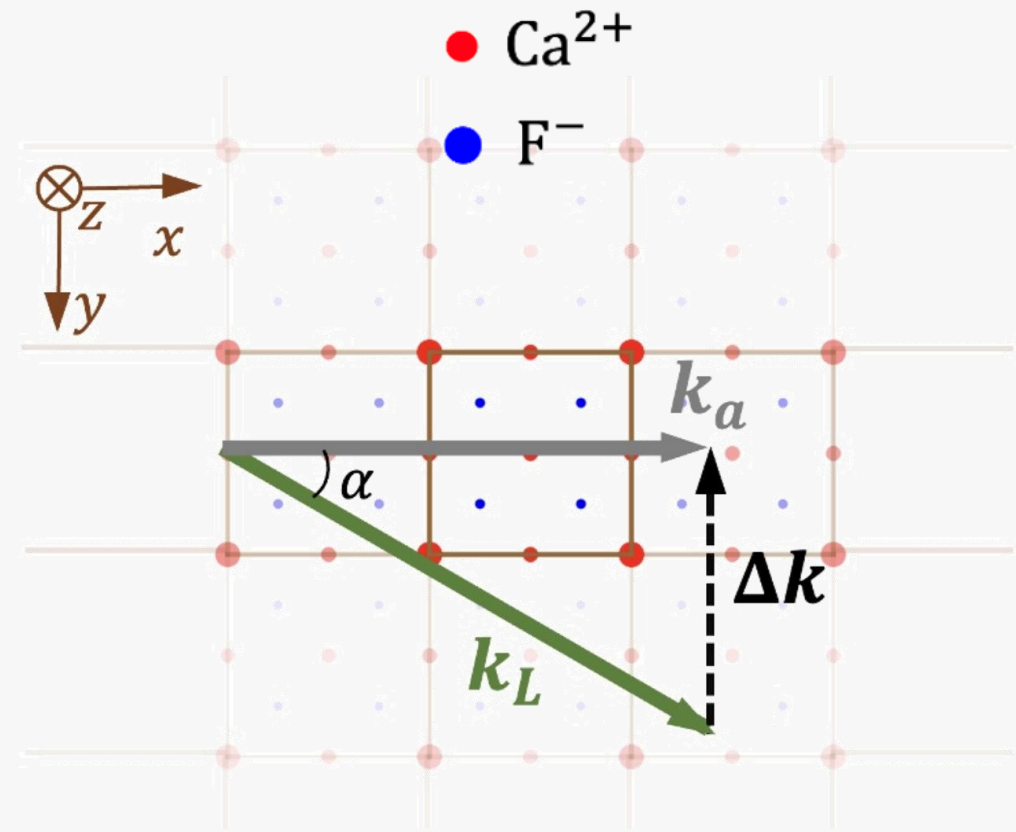


Fig: Schematic figure for coherent condition.