

neutron EDM with SuperFluid ^4He (and ^3He) Experiment

Kent Leung (on behalf of the nEDMSF collaboration)

“Electric Dipole Moments: Experimental and Theoretical Horizons”

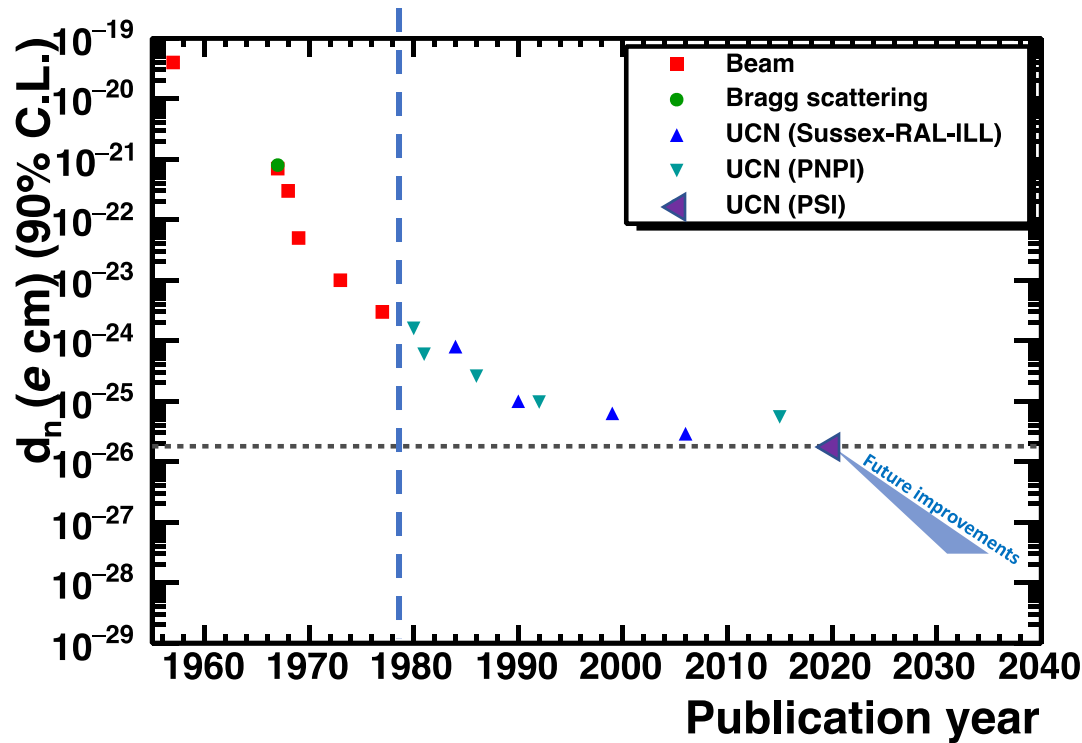
Caltech, May 12-14, 2025

Organized by Caltech and the T.D. Lee Institute, Shanghai



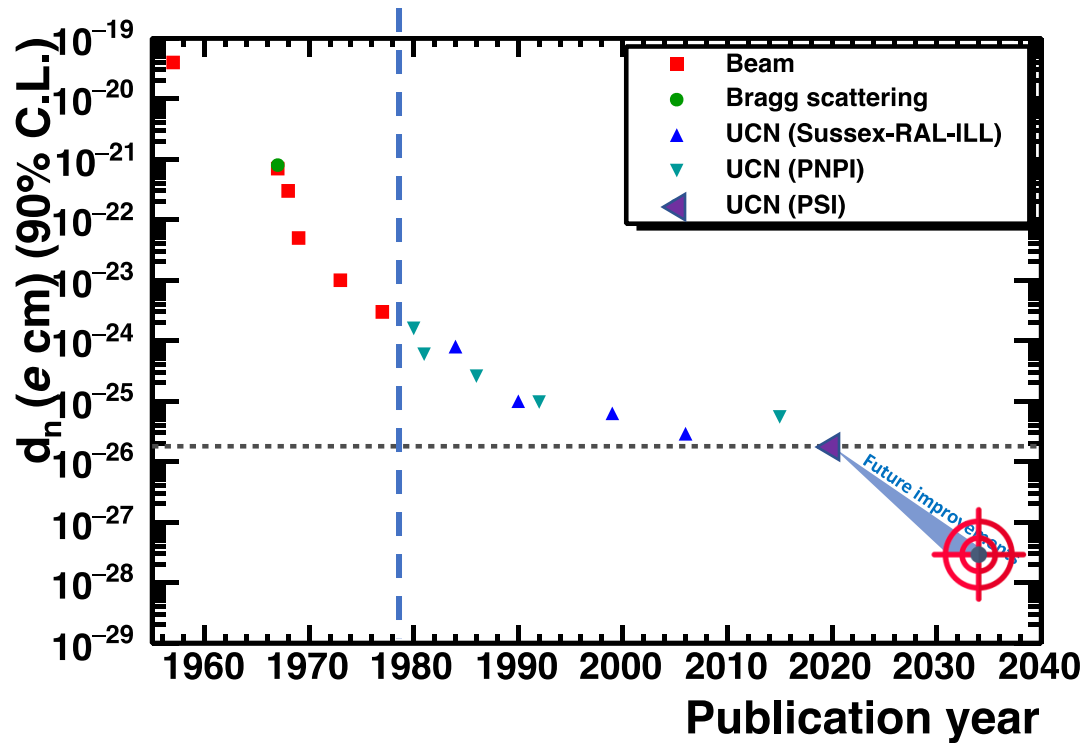
**MONTCLAIR STATE
UNIVERSITY**

Neutron EDM sensitivity



From: “Electric dipole moments and the search for new physics” (2022 community white paper) arxiv:2203.08103

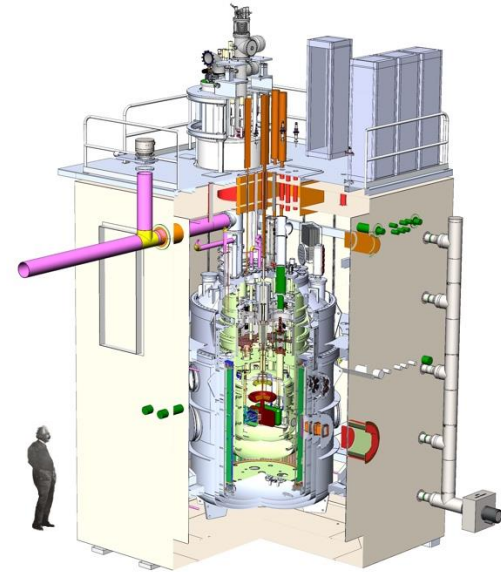
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nEDMSF's new approach

- Funding for the “legacy” nEDM@SNS was cancelled by US Department of Energy (Office of Science) in Nov. 2023
 - History: collaboration was formed in ~ 2002
 - First funding from DOE/NSF in ~2012-14
 - Total was ~ \$58M by DOE count
 - About half the apparatus has been constructed
 - Cancellation letter noted that the experiment had overcome “*a number of impressive technical challenges ... and in some cases developing new technical solutions that have value on their own*” but also that there remained a “*low state of technical readiness for critical subsystems*”
 - Had a significant R&D component remaining
- Based on the investment to date, we start building up the measurement technique in a **staged approach**
 - **New international nEDMSF collaboration formed**
 - Neutron sources: **Institut Laue-Langevin (France) & European Spallation Source (Sweden)**
 - Incorporate what we learn AND new ideas into subsequent steps



Institut Laue-Langevin (France)



European Spallation Source (Sweden)



Dr. Paul Mantica
Facilities and Project Management Division
Office of Nuclear Physics (NP)
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19901 Germantown Rd.
Germantown, MD 20874

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Program Director, Nuclear Physics - Experiment
National Science Foundation
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Alexandria, VA 22314

Dr. Timothy Hallman
Associate Director of Science for Nuclear Physics
US Department of Energy
1000 Independence Avenue, SW
Washington, DC 20585-1290

We [...] suggest an **evolution of the project’s achievements towards a larger international collaboration** for a cryogenic neutron EDM measurement **hosted at the best possible neutron source in the world.**





[...] the **U.S. contribution to the new collaboration** would therefore be, in first order, the **in-kind transfer of equipment as well as operational support for the U.S. groups to participate.** [...] the **US would remain the major player within the collaboration** with corresponding weight in key decisions regarding project evolution.

We [...] give our **strongest support for the future nEDMSF project [...].** Beam time proposals of this collaboration were recently **evaluated by ILL’s scientific subcommittee and the scientific excellence was clearly recognized.** The subcommittee highlighted the **strategical importance of this future-of-the-field-defining experiment** and **recommended ILL to support this program.** The critical importance of such a global nEDM program is also **endorsed by nuclear and particle physics communities** within the **NUPECC Long range plan and the European Strategy for Particle Physics.**

ILL’s management and subcommittee **welcome the proposed staged approach [...].**

Yours sincerely,

Dr. Ken Andersen
Director
Institut Laue-Langevin
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Ken Andersen ILL Director	Jacques Jestin ILL Science Director	Helmut Schober ESS Director	Giovanna Fragneto ESS Science Director
			

Staged Approach to Experiment Realization I

- A proof-of-principle for 3 key techniques in the experiment

Phase 1 (Institut Laue-Langevin's PF1B cold neutron beam)

- Ultracold Neutron production in superfluid & detection of $n + {}^3\text{He}$ capture light in a (legacy) measurement cell (largely existing apparatus)

Phase 2 (Institut Laue-Langevin's PF1B cold neutron beam)

- UCN production, detection of capture light & *polarization* (largely existing apparatus)

Phase 3

- Superfluid HV test with Cavallo multiplier, realistic electrodes, measurement cell (significant existing apparatus)



International Collaboration for Phase I experiments at the Institut Laue-Langevin

1	Brad Filippone	<i>Caltech</i>	16	Weijun Yao	<i>Oak Ridge National Laboratory</i>
2	Skyler Degenkolb	<i>Heidelberg University</i>	17	Maurits van der Grinten	<i>Rutherford Appleton Laboratory</i>
3	Hanno Filter	<i>Institut Laue Langevin</i>	18	David Milstead	<i>Stockholm University</i>
4	Tobias Jenke	<i>Institut Laue Langevin</i>	19	Peter Fierlinger	<i>Technical University of Munich</i>
5	Michael Jentschel	<i>Institut Laue Langevin</i>	20	Robert Georgi	<i>Technical University of Munich/FRM II</i>
6	Oliver Zimmer	<i>Institut Laue Langevin</i>	21	Thomas Neulinger	<i>Technical University of Munich/FRM II</i>
7	Benoit Clément	<i>Laboratoire de Physique Subatomique & Cosmologie</i>	22	Florian Piegsa	<i>University of Bern</i>
8	Kent Leung	<i>Montclair State University</i>	23	Doug Beck	<i>University of Illinois at Urbana-Champaign</i>
9	Robert Golub	<i>North Carolina State University</i>	24	Chen-Yu Liu	<i>University of Illinois at Urbana-Champaign</i>
10	Paul Huffman	<i>North Carolina State University</i>	25	Chris Crawford	<i>University of Kentucky</i>
11	Ekaterina Korobkina	<i>North Carolina State University</i>	26	Wolfgang Korsch	<i>University of Kentucky</i>
12	Vince Cianciolo	<i>Oak Ridge National Laboratory</i>	27	Brad Plaster	<i>University of Kentucky</i>
13	Paul Mueller	<i>Oak Ridge National Laboratory</i>	28	Valentina Santoro	<i>University of Lund</i>
14	John Ramsey	<i>Oak Ridge National Laboratory</i>	29	Clark Griffith	<i>University of Sussex</i>
15	Andy Saunders	<i>Oak Ridge National Laboratory</i>	30	Steve Lamoreaux	<i>Yale University</i>

Staged Approach to Experiment Realization II

- Three major areas requiring further development

High voltage

- Complete development of HV delivery, robust electrode material
 - LN2 testing of Cavallo multiplier at LANL ongoing

Composite vessels

- Vessel holding ~1400 L of superfluid helium in AC magnetic field (spin dressing)
 - Have started working with Rutherford Appleton Laboratory/Sussex to continue existing development

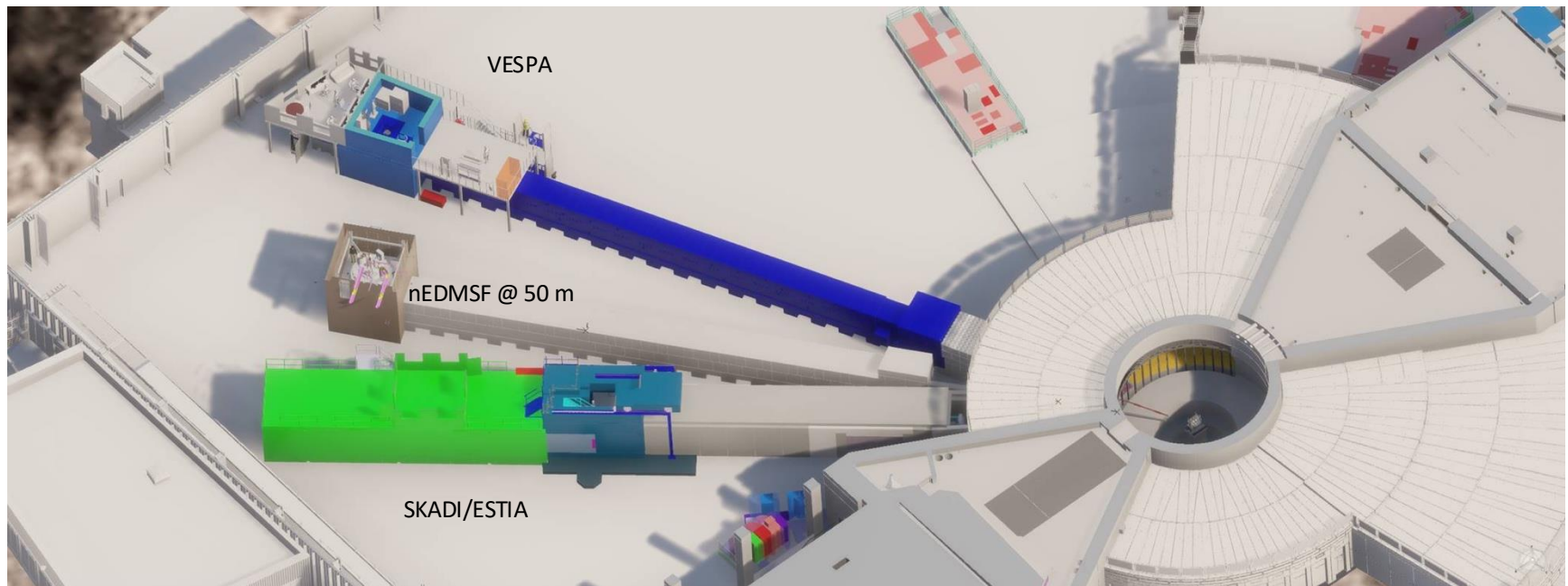
Helium-3 services

- Completion of system to add polarized ^3He , then remove it after measurement

Many opportunities for new contributions!

Staged Approach to Experiment Realization III

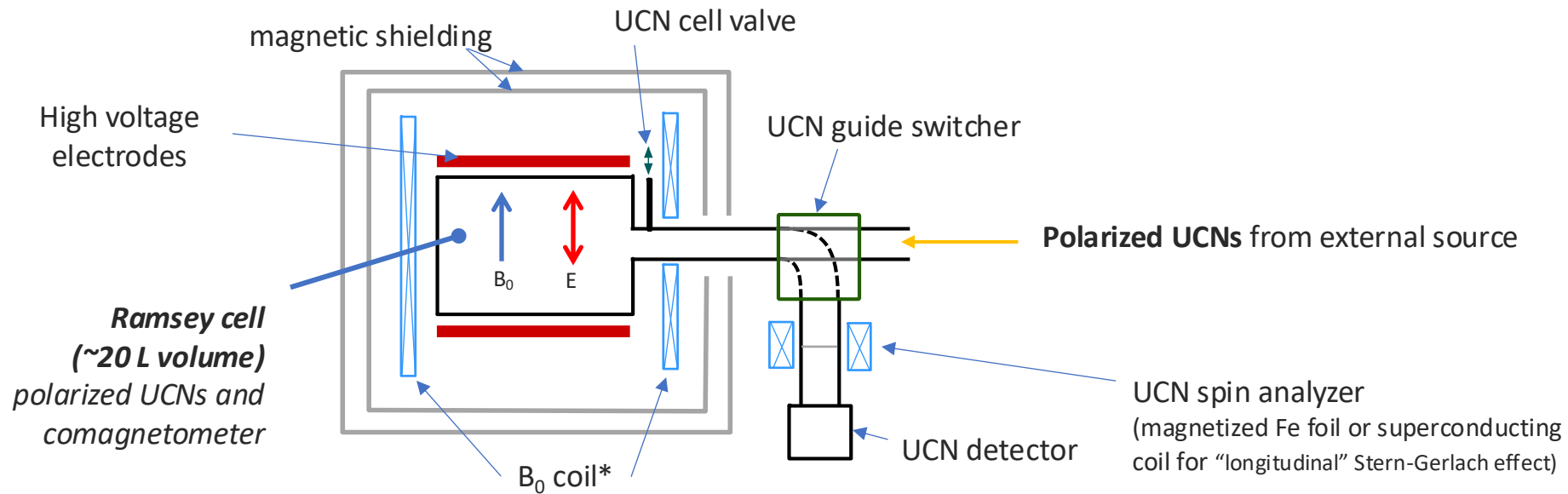
- The physical scale of a superfluid experiment is difficult to accommodate at ILL
- Have started working with ESS collaborators on possible beamline design
 - Published ANNI beamline design has $\sim 10 \times$ flux of at SNS*
- Engineering on beamline @ ESS started (March 2025)
- Expect higher cold neutron flux (at 1 meV) c.f. SNS



* new, “existing” guide in monolith will reduce the flux by factor ≤ 2

Experimental Overview

“Traditional” Ramsey nEDM experiment with UCNs



At the **end of free precession** (i.e. at a single time), neutrons’ **final phase** measured by counting number of spin down N_{down} .

(which is done **ex-situ** for UCNs measurements)

(modern experiments also measure N_{up} for normalization & have two cells with opposite E-field)

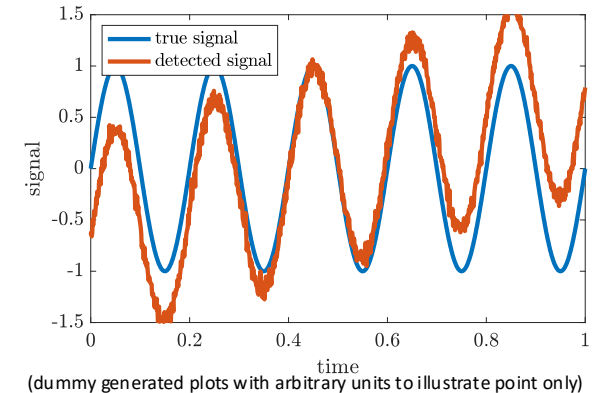
“Always measure frequency...” (Rabi? Ramsey? Wieman?)

- Control and measurements of time/clocks can be done to a high precision.
 - ✓ used in Ramsey technique
- **BUT** the power of a frequency measurement is needing only the *relative change* of a signal **over short time scales**
 - Drifts in system on time scales greater than a few oscillations are suppressed



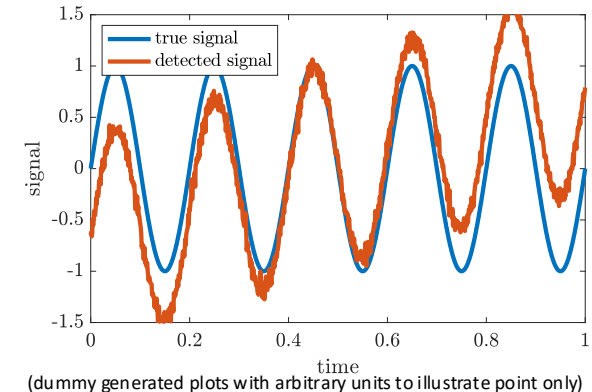
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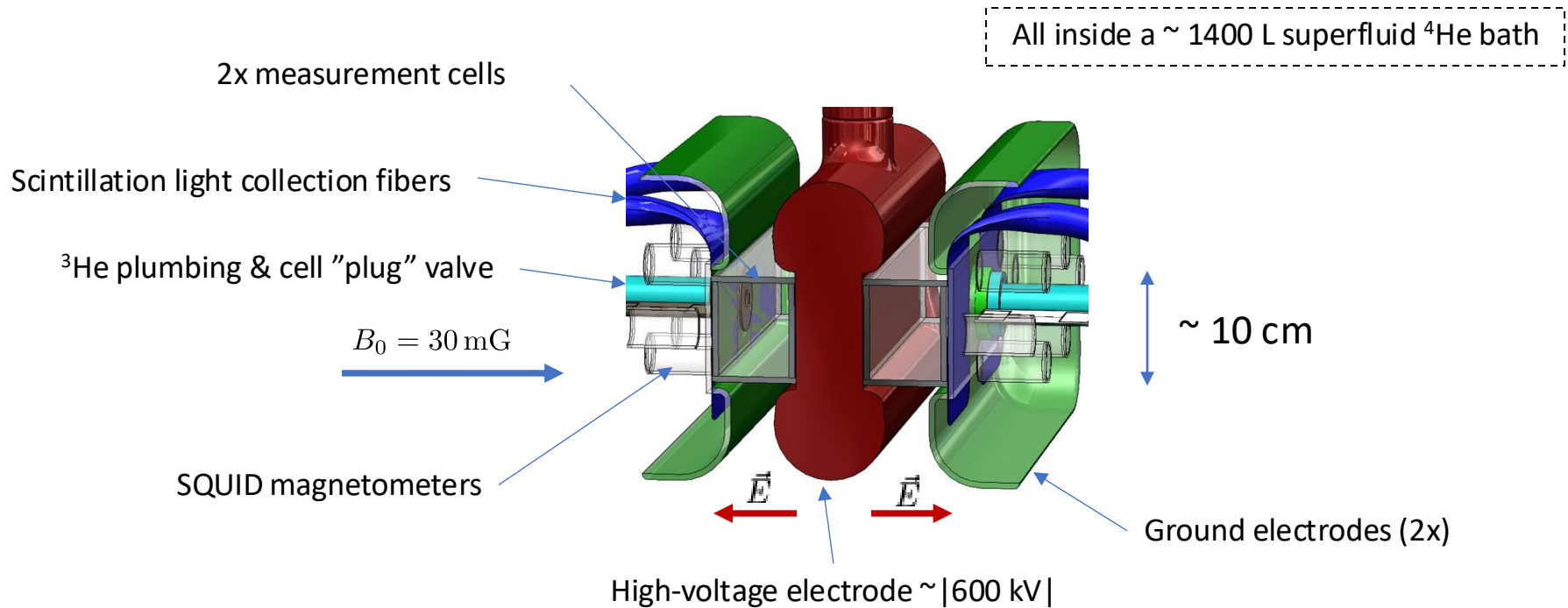
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- Recall: **Ramsey’s technique** only measures **final neutron phase via. counting neutrons**. (Can’t know the frequency at all times.)
- **Need two** (possibly four) **separate repeated Ramsey cycles** (with different clock frequencies within the same cell) to **determine a single EDM value**.
- Each cycle takes ~100-200s. **A lot can happen in this time:**
- Detector efficiency drifts, magnetic field drifts, $\pi/2$ efficiency and other polarization drifts, UCN source intensity drifts, etc..
- Can be corrected but introduces additional statistics & systematics which becomes more difficult **for 10^{-28} e.cm level**



nEDMSF: UCNs + superfluid ^4He + polarized ^3He

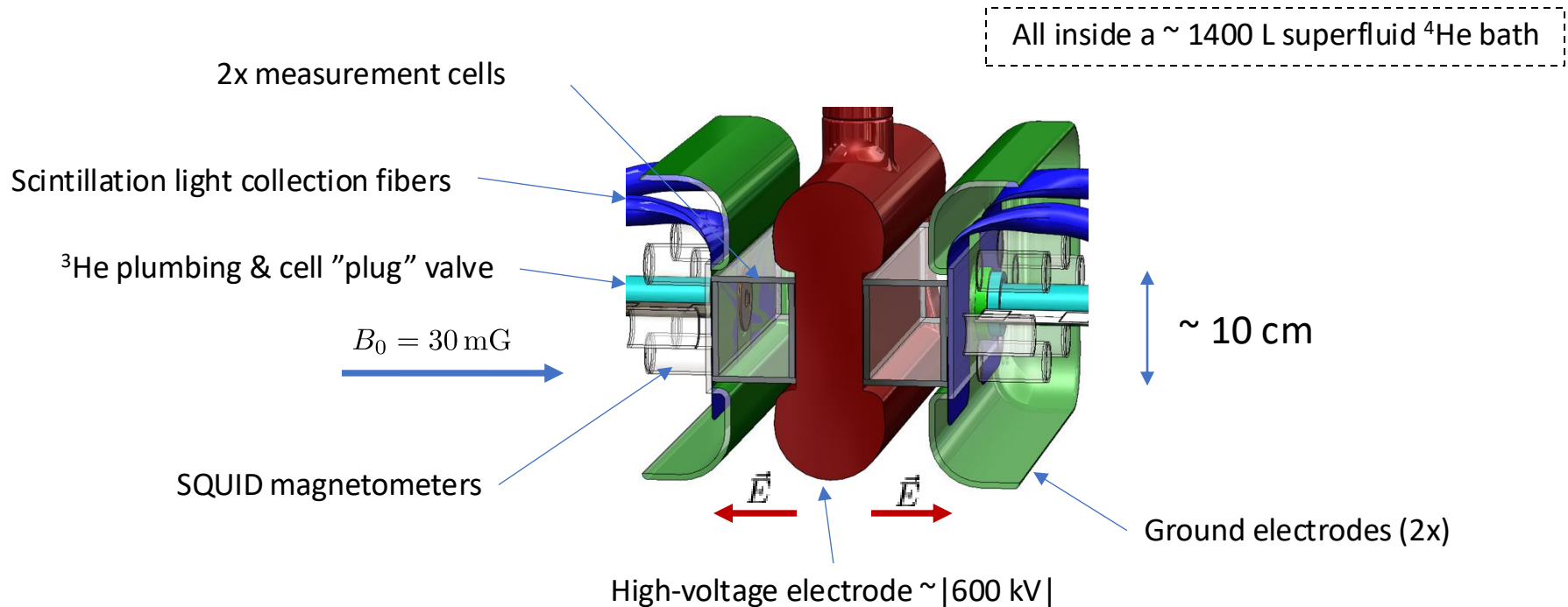
Based on Golub & Lamoreaux, Phys. Rep. (1994)



- **Double cell** setup with E-field relative to B-field opposite in each cell
- **>95% polarized ^3He** loaded into cell (dissolved in isotopically-pure superfluid ^4He at 0.4 K).
- **^3He serves as a comagnetometer and UCN spin analyzer** (see later)

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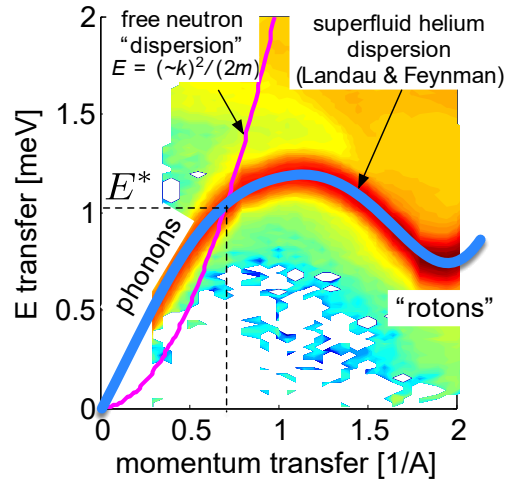


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- **^3He serves as a comagnetometer and UCN spin analyzer** (see later)
- **In-situ super-thermal UCN (~ 100 neV) production** and accumulation in **superfluid ^4He with polarized 1 meV cold neutron beam** (direction into the page).
- Cold neutron beam $\sim 100\%$ polarized \rightarrow produce $\sim 100\%$ UCNs polarized.

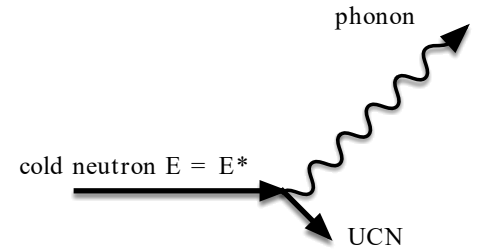
Roles of the superfluid ^4He

- 1 meV cold neutrons (11 K) scatter off phonons in superfluid ^4He to become UCNs (< 160 neV, “2 mK”)

Log contour plot of dynamic structure factor of He-II @ 1.2K
from [Andersen et al. J. Phys. Condens. Matter (1994)]



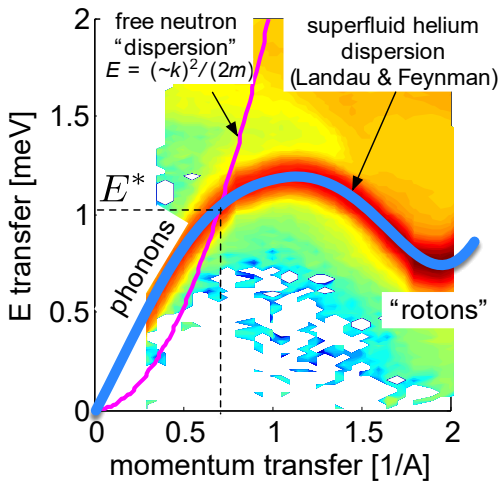
Down-scattering from single phonon
($\sim T$ independent):



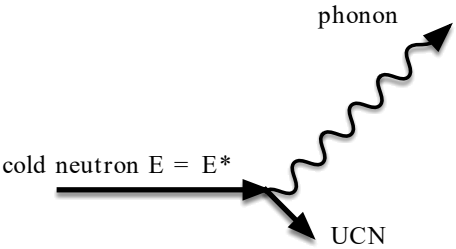
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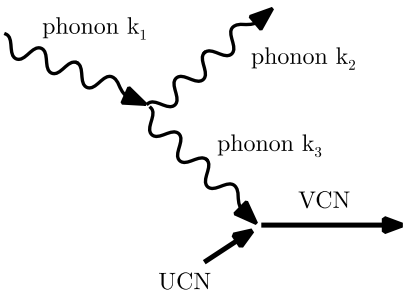
Down-scattering from single phonon
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Dominant UCN loss is two-phonon scattering:

$$\tau_{\text{up},2\text{-phonon}} = (100 \text{ s K}^7) T^{-7}$$

[Golub & Pendlebury, Physics Letters A (1977)]



Experimental studies & demonstrations:

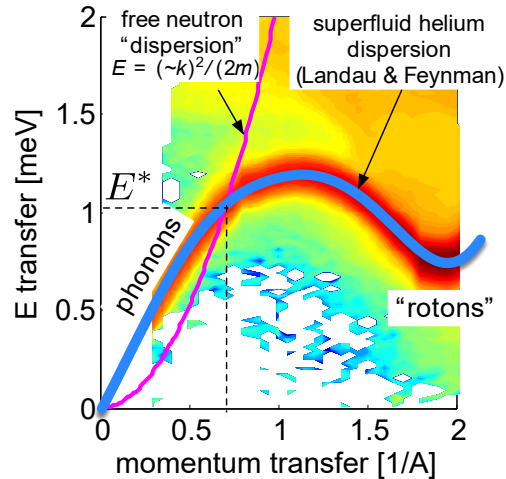
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- At $T = 0.4\text{K}$, up-scattering (or “thermalization”) $\tau_{\text{up}} \approx 20$ hours. If thermalized, UCNs are lost.
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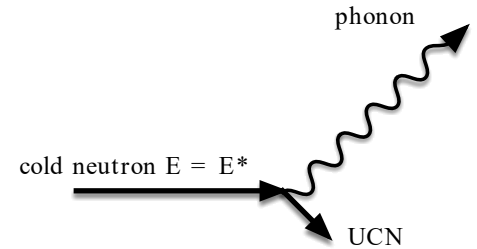
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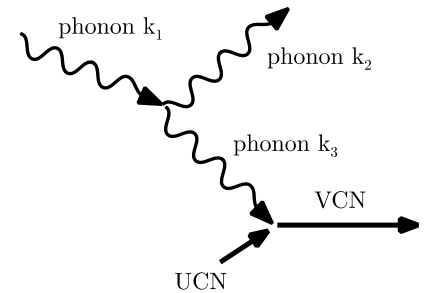
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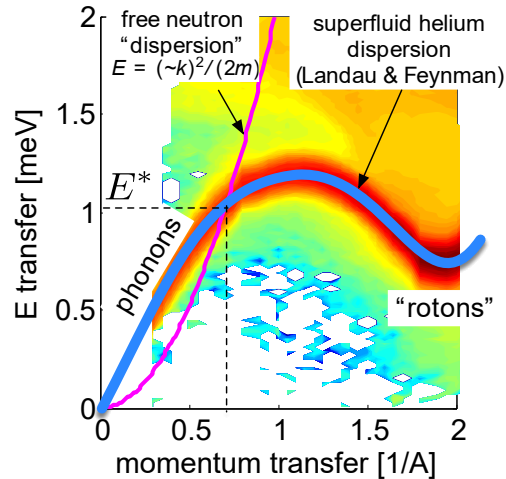
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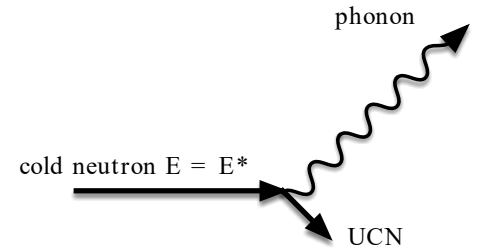
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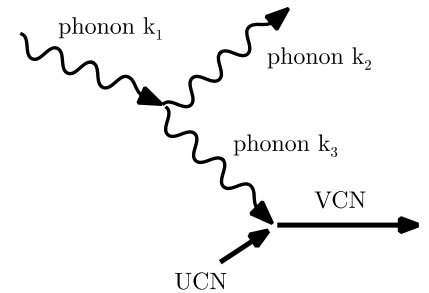
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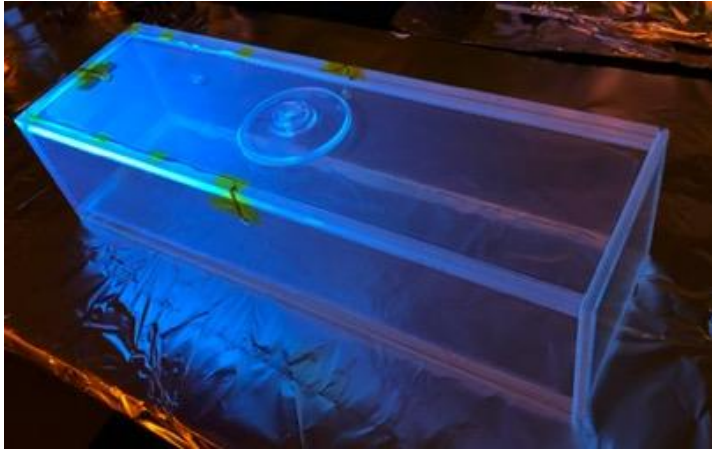
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- Superfluid also **scintillates** at ~ 80 nm (EUV) \rightarrow used to detect n- ^3He capture events (later)
- ^3He atoms scatter off phonons \rightarrow mean-free-path $\sim T^{7.5} \rightarrow$ important for key “false EDM” systematic control

Measurement cells



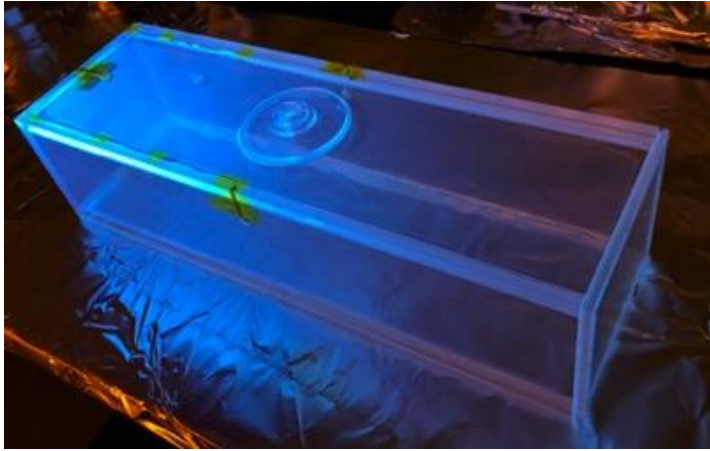
Full-sized prototype measurement cell illuminated by 300 nm UV lamp

- Inside: 7.5 cm (W), 10 cm (H), 40 cm (D) rel. to beam.
- Inner walls: **deuterated PS + d-TPB** (for fluorescent properties to detect 80 nm ^4He scintillation).
- Neutron optical potential = **170 neV**
- Design goal wall loss $\tau_{\text{cell}} = 2,000 \text{ s}$

- **Cryogenic UCN storage cell** => **low wall loss**. Design goal wall loss $\tau_{\text{cell}} = 2,000 \text{ s}$. Due to unavoidable β -decay:

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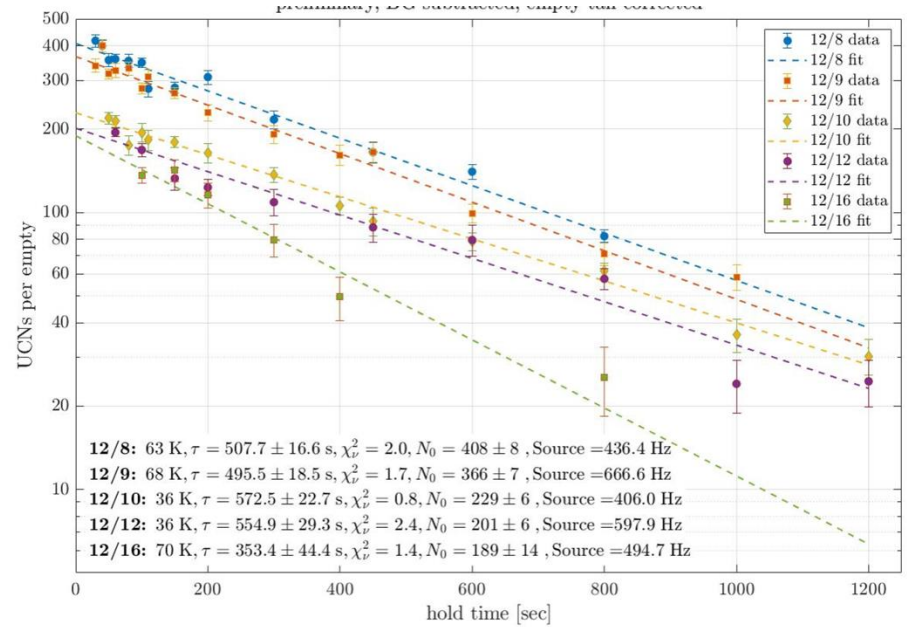
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- Long $\tau_{\beta+\text{cell}}$ allows a large number of UCNs accumulated per fill. **Polarized UCN density $\sim 180\ \text{UCN}/\text{cm}^3$**
=> $\sim 500,000$ UCNs per cell.
- **UCN “usefulness” time restricted by $\tau_{\beta+\text{cell}}$** (transverse spin coherence times are $\sim 20,000\ \text{s}$)
- “Useful” UCN time of **1,000 s** use in a measurement cycle

Measurement cell II

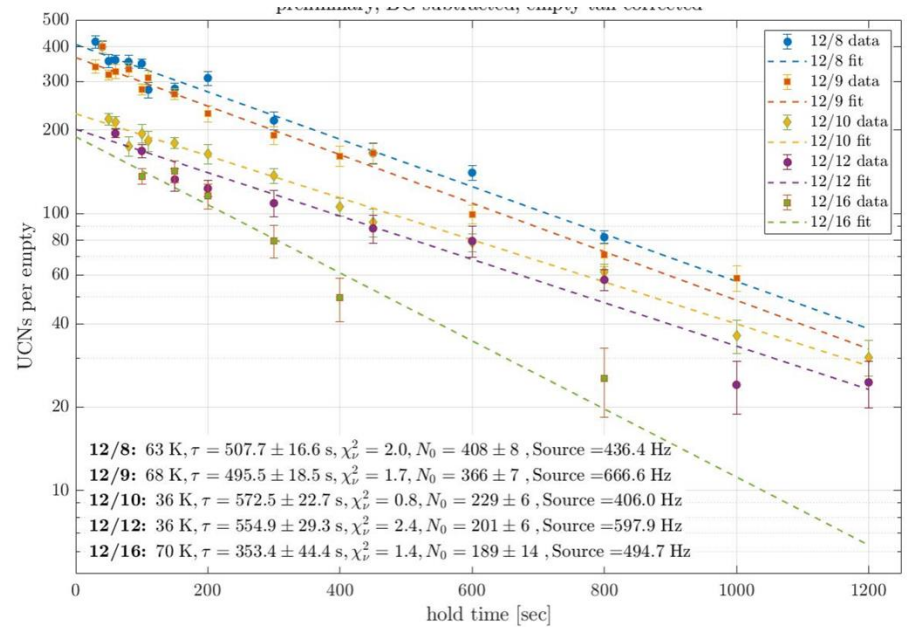
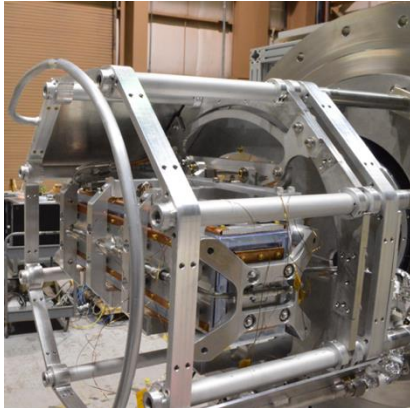
Fill and empty UCN storage measurements at LANL (external UCN source)



- **Latest cell** tested at LANL UCN source had $\tau_{\beta+\text{cell}}$ of 570 ± 20 s. (Single exponential decay only observed)

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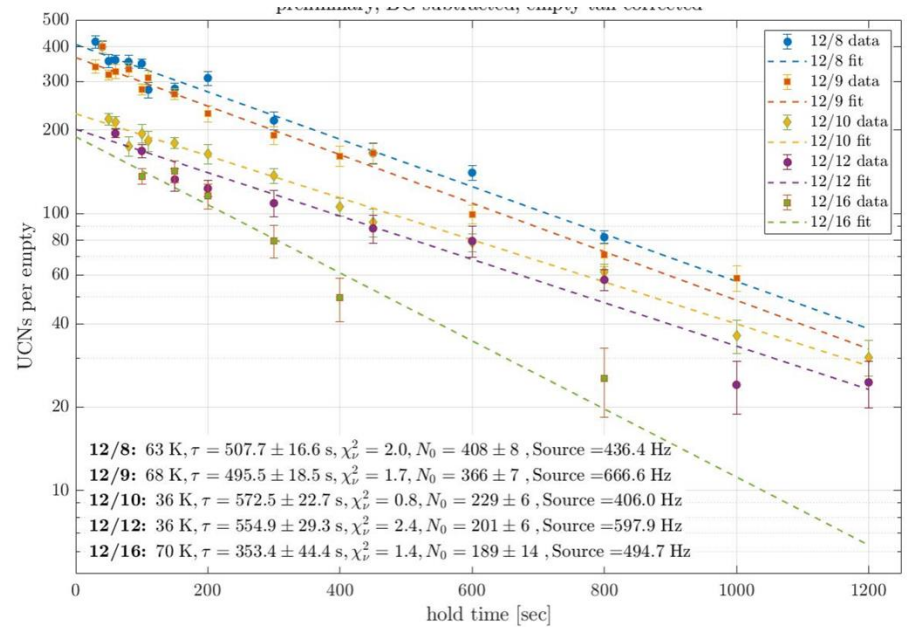
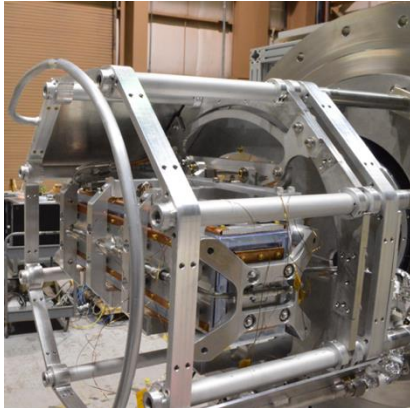


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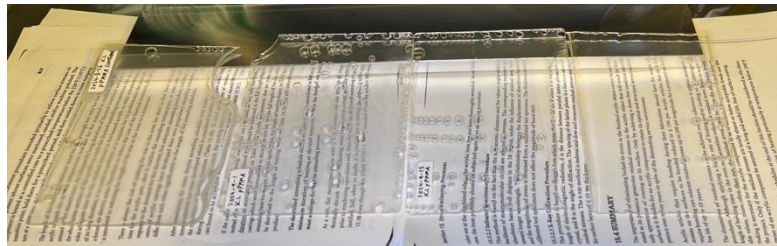


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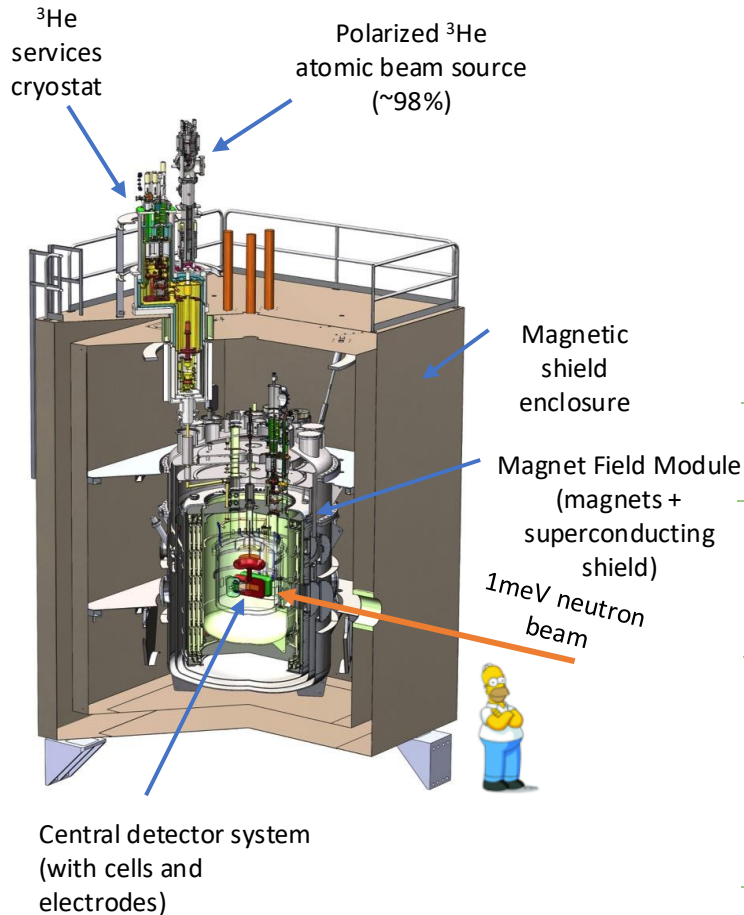
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- UCN spectrum from external source different to superfluid -> building spectrometer to understand
- Disadvantage to in-situ: **cold neutron beam induced activation of cell**
 - End windows d-PMMA purity to ppm level. Recently achievement **in-house polymerized** end windows 150 Hz -> **27 Hz background (10% impact in sensitivity)**



Approximate sensitivity scaling



Statistical “shot noise” figure of merit: $\sigma(d_n) \sim \frac{\hbar}{2\alpha ET\sqrt{N}}$

E = electric field 75 kV/cm in superfluid ⁴He @ ~ 2 atm pressure
vs ~ 10 kV/cm in vacuum

α = polarization contrast (UCN & ³He polarization ~ 98%)

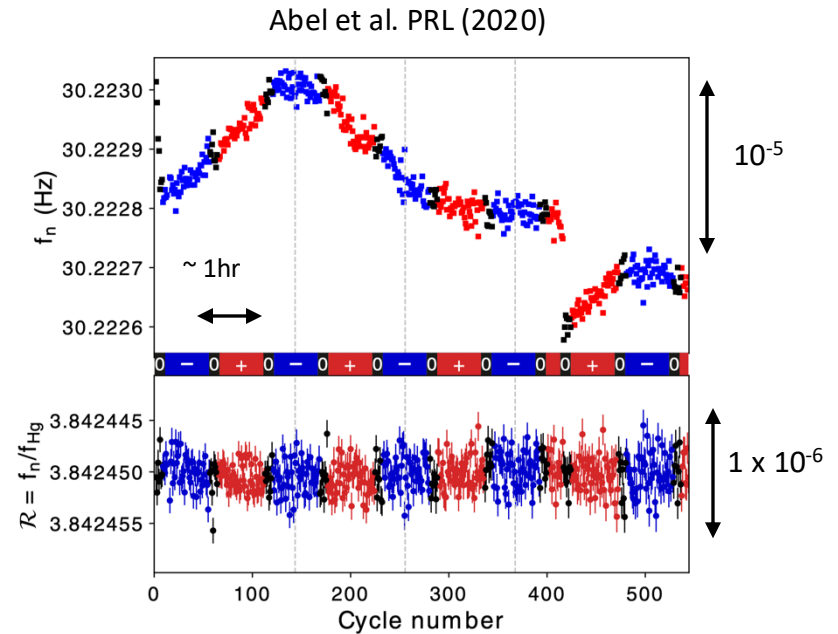
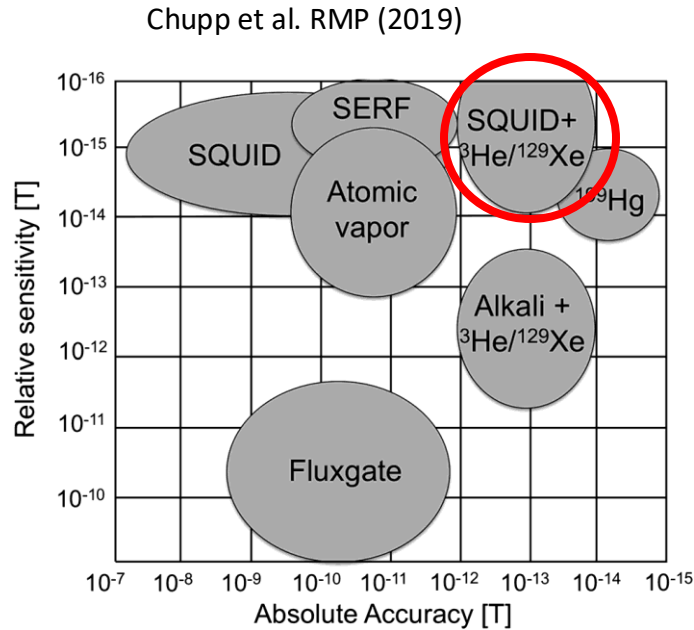
T = precession time (we can use 1000 s vs. ~ 200 s for room temp. experiments)

N = no. detected neutrons 2-3x improvement on running/close-to-running 10^{-27} e.cm experiments.
We have high-density in small cells which reduces systematics (see later)

Systematics: small cell -> smaller region to achieve average field spec < 100 nG/cm.
Superconducting Pb shield -> great for suppressing field drifts

Comagnetometry

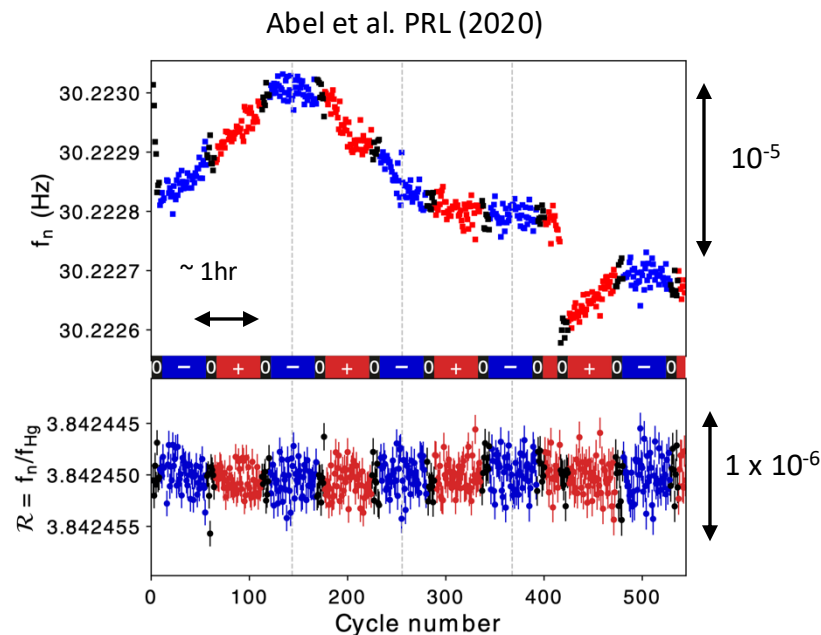
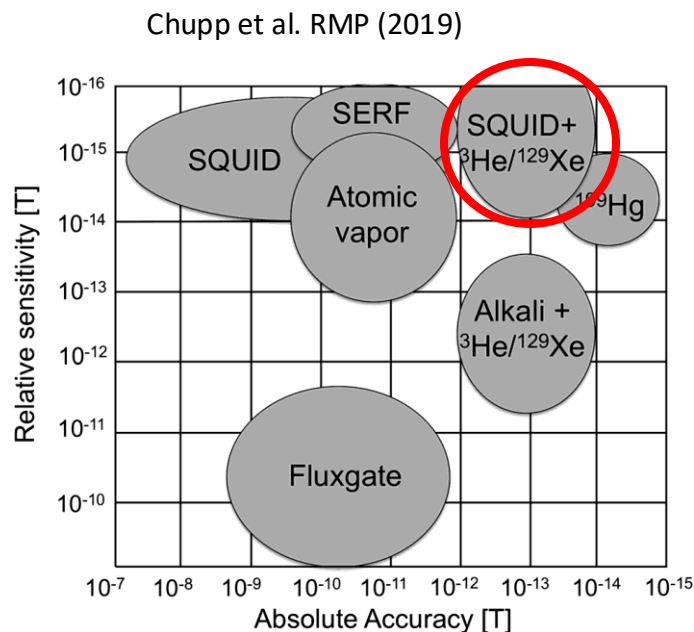
- A magnetometer species (typically a gas) that “cohabits” the same volume as the UCNs experience.
- Need sufficient comagnetometer atoms for precision; too high can cause electric breakdowns & UCN loss ($\sim 10^{-6}$ mbar)



(for scale: 10^{-28} e.cm sensitivity $\rightarrow \sim 10^{-10}$ fractional frequency shift)

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- nEDM experiments ultimately measure relative to comagnetometer’s EDM (usually “Schiff suppressed”)
- However, the two species still experiences **different (time & space) averaged magnetic fields** because:
 - UCNs have such low speeds, they “sag” in gravity a few mm.
 - Relativistic $\vec{E} \times \vec{v}$ motional-field related effects are different. One produces a *false* EDM (most serious systematic error)

Polarized ^3He as live and in-situ UCN spin analyzer

- Polarized ^3He gas cells widely used as neutron beam spin analyzers (count survivors)
- Strong **spin-dependent capture cross-section**: $^3\text{He} + \vec{n} \rightarrow \text{p} + ^3\text{H} + 764 \text{ keV}$

Anti-parallel spins: $\sigma_{\downarrow\uparrow,\text{thermal}} \approx 11 \text{ kb}$

Parallel spins: $\sigma_{\uparrow\uparrow,\text{thermal}} \lesssim 0.1 \text{ kb}$

- Capture rate** for polarized UCNs and ^3He in same volume :

$$\dot{N}_3 = N_{\text{UCN}} \bar{\tau}_3^{-1} (1 - P_n P_3 \cos \theta_{n3})$$

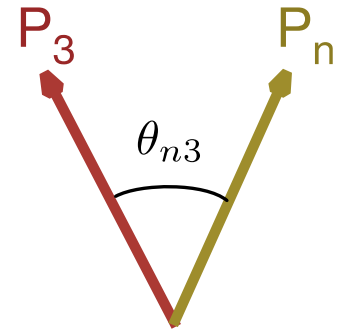
number of UCNs

polarizations

angle between spins

where $\bar{\tau}_3^{-1} \approx [n_3 \sigma_{\downarrow\uparrow,\text{thermal}} \times (2200 \text{ m/s})]/2$

^3He number density



Polarized ^3He as live and in-situ UCN spin analyzer

- Polarized ^3He gas cells widely used as neutron beam spin analyzers (count survivors)
- Strong **spin-dependent capture cross-section**: $^3\text{He} + \vec{n} \rightarrow \text{p} + ^3\text{H} + 764 \text{ keV}$

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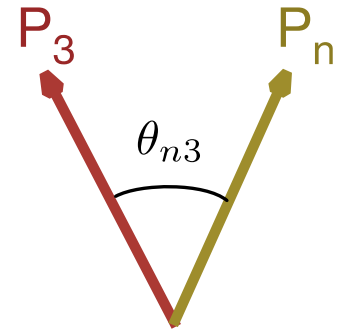
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^3He number density



or “continuous”

- Detect the 760 keV scintillation as generated \rightarrow in-situ and live UCN spin analyzer
 - 80 nm scintillation light \rightarrow blue photons (walls of cells) \rightarrow wavelength shifting fibers \rightarrow optical fibers \rightarrow Silicon Photomultipliers

Live & in-situ UCN spin analysis

- Can't look **too intensely** or will kill UCNs too quickly



- Want ^3He -n capture rate to be similar to UCN loss time (cell + β), i.e. $\bar{\tau}_3 \approx 500\text{ s}$

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→ **Achievable with atomic beam source.**

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→ **Achievable with atomic beam source.**
- This n_3 & P_3 of ^3He produces $\sim 6 \text{ fT}$ fields → **detectable by SQUIDs.**

^3He can be used as comagnetometer AND live & in-situ UCN spin analyzer!

False EDM systematic effect

- Recall: nEDM is measured relative to the comagnetometer's EDM. Most effects cancel out with opposite E-field.
 - Comagnetometer's EDM suppressed by Schiff screening but **can experience a false EDM**
 - From *interaction* between the $\vec{E} \times \vec{v}$ motional field and magnetic field gradients (“geometric-phase induced false EDM”)

Radius of cell
 - “Discovered” in the nEDM field, false EDM for comagnetometer: $d_{af} = \frac{J\hbar}{4} \left(\frac{\partial B_{0z}}{\partial z} \right) \frac{\gamma^2 \boxed{R^2}}{c^2} \left[1 - \frac{\omega_0^2}{\omega_r^2} \right]^{-1}$
 Pendlebury et al. Phys. Rev. C (2004)
- (Simplified cylindrical cell but general relationships hold. Rectangular cell work by Swank & Golub [Phys. Rev. A 93, 062703 \(2016\)](#))
- Note:** The transverse spin coherence time scales as R^4 . Want small cells (with high UCN density)

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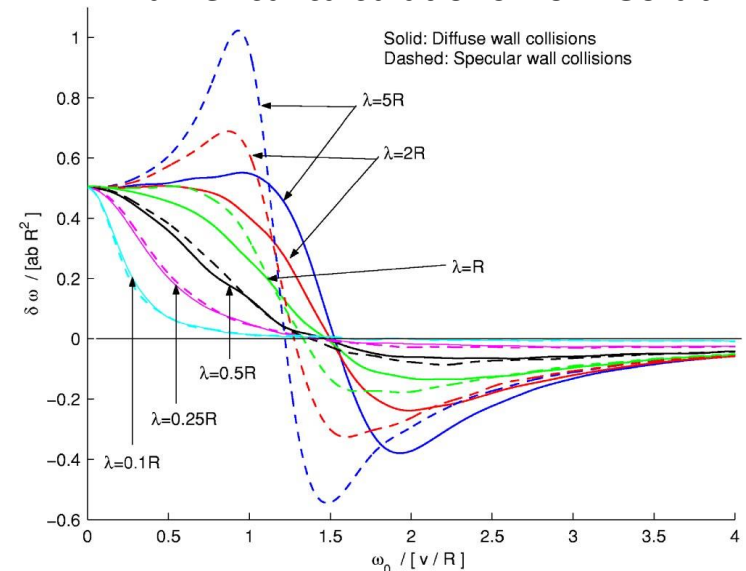
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- Note:** The transverse spin coherence time scales as R^4 . Want small cells (with high UCN density)
- Can change ^3He -phonon scattering mean-free-path by changes in superfluid temperature:

$$\lambda_{^3\text{He}} \approx 0.077 \text{ cm} \times \left(\frac{0.45 \text{ K}}{T} \right)^{15/2}$$

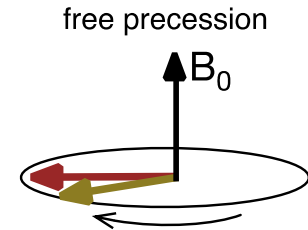
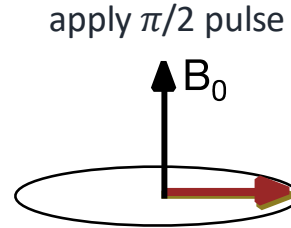
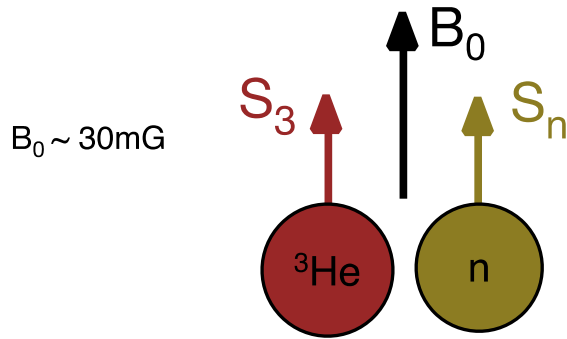
Can tune temperature to make false EDM zero by scanning T!

Numerical calculations from Golub:



**Two measurement modes of nEDMSF:
double free precession & critical dressed spin**

Double free precession mode



- Time evolution of phase: $\theta_{3n}(t) = \theta_3(t) - \theta_n(t) = \left[(\gamma_n - \gamma_3)B_0 \pm \frac{2d_n|E|}{\hbar} \right] t + \phi_0 \equiv \omega_{3n}^\pm t + \phi_0$,
- With $B_0 = 30\text{ mG}$: $\gamma_3 B_0 / (2\pi) \approx 100\text{ Hz}$ $\gamma_3 \approx 1.1 \gamma_n$
- Scintillation light oscillation frequency: $|\gamma_n - \gamma_3| B_0 / (2\pi) \approx 10\text{ Hz}$
- The transverse spin coherence time (wall depolarization + gradient depolarization), $T_2 \sim 10,000\text{ s}$
- Flipping high-voltage electrode often with known sequences to suppress 1st order drifts (e.g. + - - + - + + -) and analysis as a "super-asymmetry"

Free precession scintillation signal

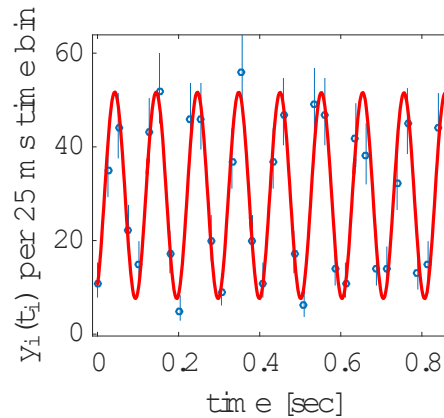
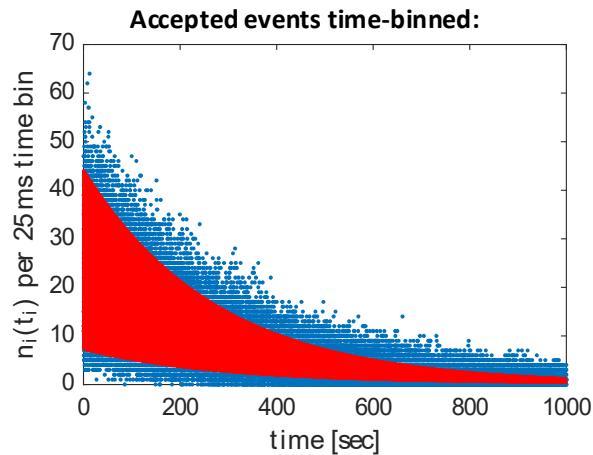
- “accepted” rate of scintillation light events: $\dot{N}_{ac-s}(t) = N_n(t) \left(\frac{\epsilon_\beta}{\tau_\beta} + \frac{\epsilon_3}{\bar{\tau}_3} \{1 - P_3(t)P_n(t) \cos[\theta_{3n}(t) + \phi_{3n0}]\} \right) + R_{BG}$
 - Acceptance probability for β -decay events (~ 0.33)
 - no. UCNs in cell
 - recall: 500 s
 - Acceptance probability for n - ^3He capture (~ 0.93)
 - polarizations
 - Background rate (could be time-dependent)

$$P_3(t)P_n(t) = P_{30}P_{n0} \exp(-t/T_{2,\text{tot}}) \approx (0.98)(0.98) \exp(-t/[10,000 \text{ s}])$$

$$N_n(t) = \int_0^{E_{\text{max}}} dE n_{n0}(E) \exp \left[-\frac{t}{\tau_{\text{cell}}(E)} - \frac{t}{\bar{\tau}_\beta} + \frac{P_n(t)P_3(t)}{\bar{\tau}_3} \int_0^t \cos \phi_{3n}(t') dt' \right]$$

UCN spectrum

Oscillating term due to previous n - ^3He absorption



- Continuously measuring the UCN phase (relative to ^3He) → **continuous frequency measurement** (via derivative)!

Free precession scintillation signal

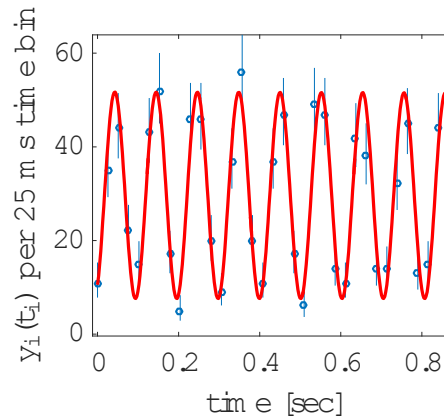
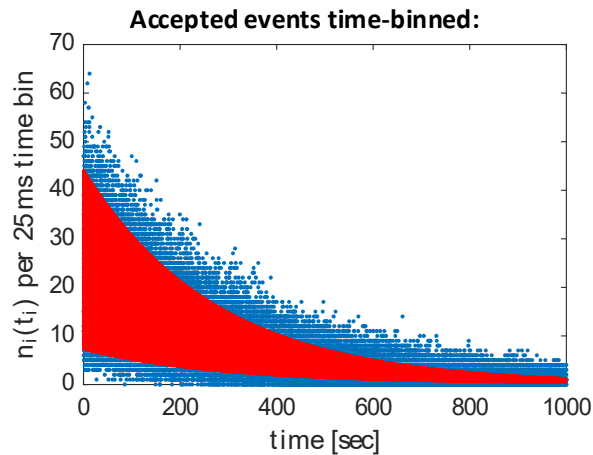
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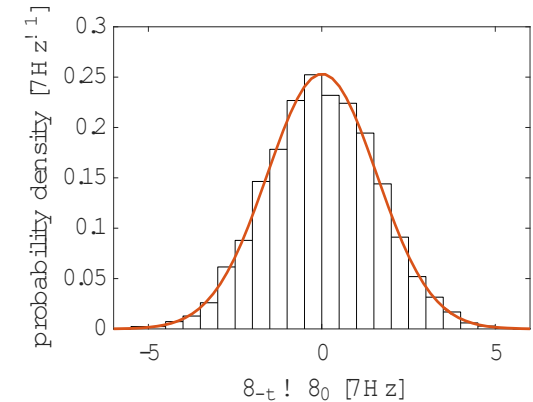
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UCN spectrum

Oscillating term due to previous n - ^3He absorption



Repeat
generation
& fit:



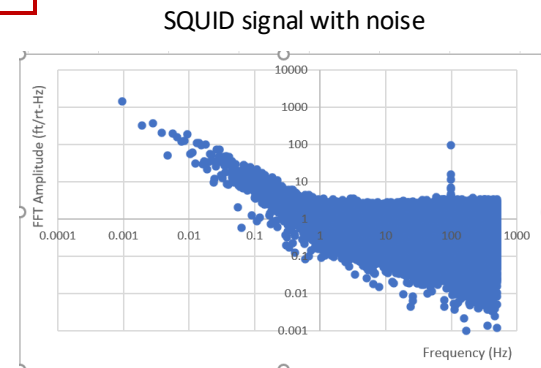
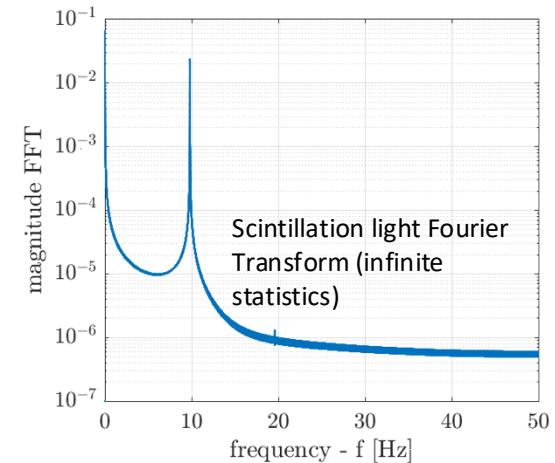
- Continuously measuring the UCN phase (relative to ^3He) → **continuous frequency measurement** (via derivative)!
- 300 live days of running (expected to take 3 years), get **1σ nEDM error = $3 \times 10^{-28} \text{ e.cm}$**

Data analysis simulations

- Neutron decay β -asymmetry
- **Spatial-variation scintillation light** detection efficiency
- **Oscillation in $N_n(t)$** due to history of n - ^3He absorption
- **Reduced parameter "contrast" fitting** to handle UCN energy-dependent wall loss
- Generation of scintillation light data with **magnetic field drifts**
- Generation of SQUID ^3He signal with noise and drifts (UKy student: Mojtaba Behzadipour)

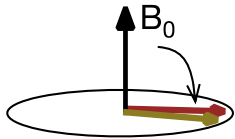
- **Simultaneously fitting SQUID signal and scintillation light signal** with global likelihood parameter (UKy student: Mojtaba Behzadipour)

- Fit **temporal field drifts** with **orthogonal polynomials**
- **Particle-by-particle neutron** scintillation data generation code
- **Magnetic field noise** in spin-dressing mode
- **Novel** spin dressing field modulation modes (Caltech grad: Raymond Tat)
- UCN spin-tracking on **Graphics Processing Units** (NCSU/Caltech/ORNL: Morano, Tat, Matthews)
- UCN center-of-mass **gravitational offset time-evolution**

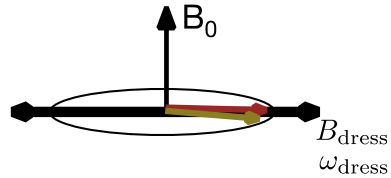


Dressed-spin mode

apply $\pi/2$ pulse



apply strong off-resonance dressing field perpendicular to B_0 to alter precession of **both** species



Cohen-Tannoudji:

$$\hat{H} = -\gamma B_0 \hat{S}_z + \hbar \omega_d \hat{a}^\dagger \hat{a} + \lambda \hat{S}_x (\hat{a} + \hat{a}^\dagger)$$

$$\lambda = \gamma B_d / 2\sqrt{n}.$$

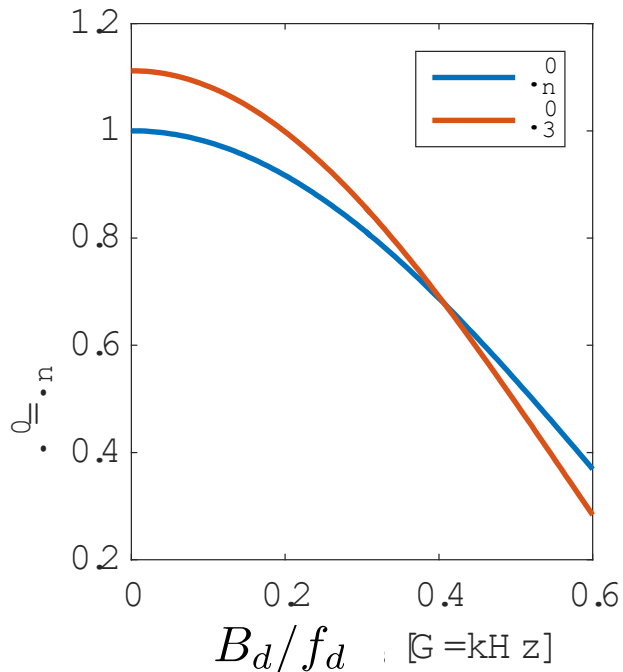
In the limit $B_{\text{dress}} \gg B_0$

effective
gyromagnetic
ratio

0th order Bessel
function

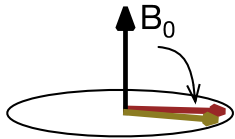
$$\gamma' = \gamma J_0 \left(\frac{\gamma B_{\text{dress}}}{\omega_{\text{dress}}} \right)$$

original

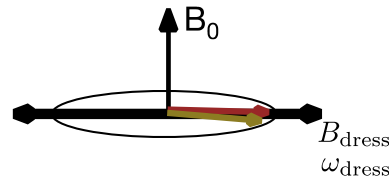


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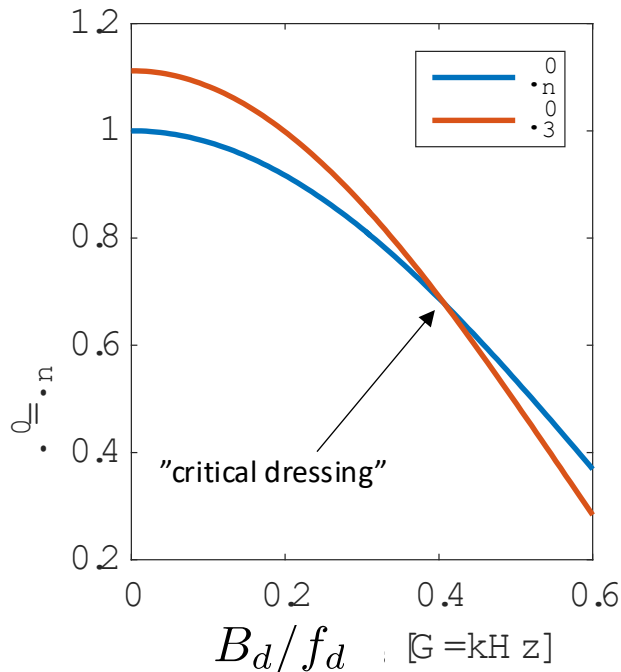
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original γ

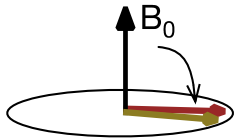
0th order Bessel function



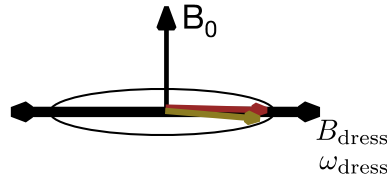
- Specific value of B_d/ν_d can make $\gamma'_3 = \gamma'_n$
- For instance, if $B_d = 1$ G is chosen, then $f_d \approx 2.5$ kHz is needed

Dressed-spin mode

apply $\pi/2$ pulse



apply strong off-resonance dressing field perpendicular to B_0 to alter precession of **both species**



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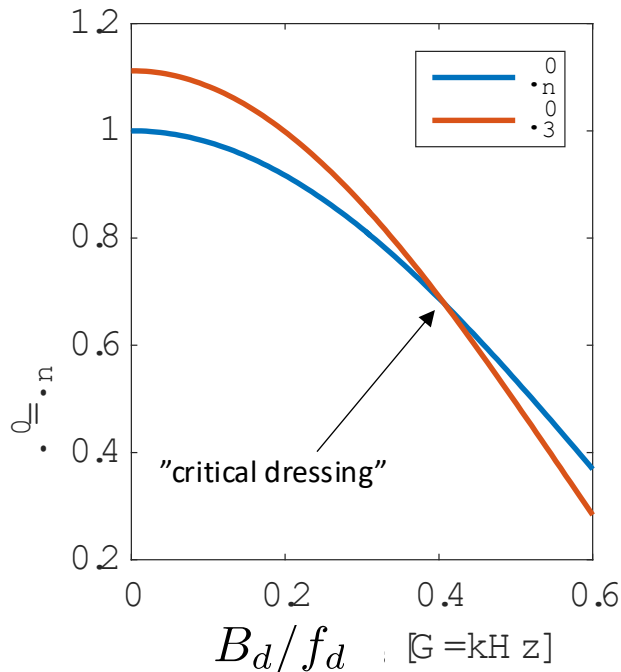
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- Specific value of B_d / ν_d can make $\gamma'_3 = \gamma'_n$
- For instance, if $B_d = 1$ G is chosen, then $f_d \approx 2.5$ kHz is needed
- If above or below critical dressing condition, then **can make neutrons effectively precess faster or slower than ^3He** as needed.

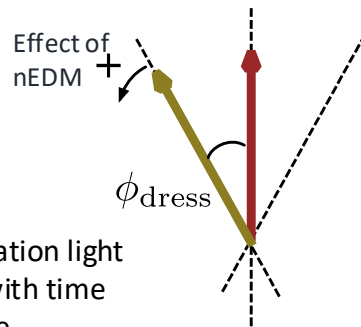
Critical modulated dressed-spin mode

- The effect of **neutron EDM** with spin dressing: $\gamma'_n B_0 \pm \frac{2ed_n E J_0 (\gamma_n B_{\text{dress}} / \omega_{\text{dress}})}{\hbar}$
- Example of modulation with “**square wave pulses**”. (Other modes possible.)

Critical dressing field to sit at a fixed ϕ_{dress}



(In rest frame of ^3He spin,
 B_0 coming out of screen)



If there's a nEDM scintillation light
 increases or decreases with time
 depending on E and cycle.

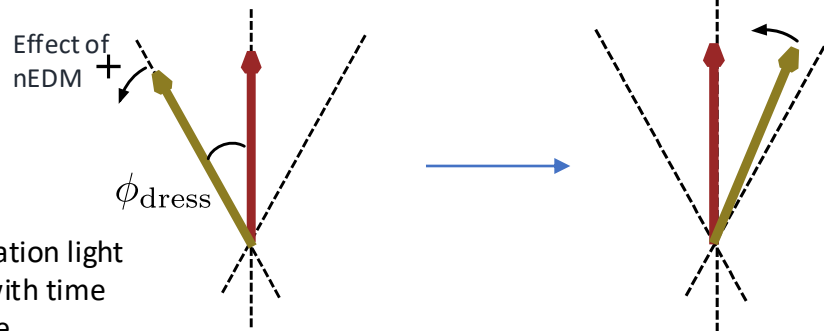
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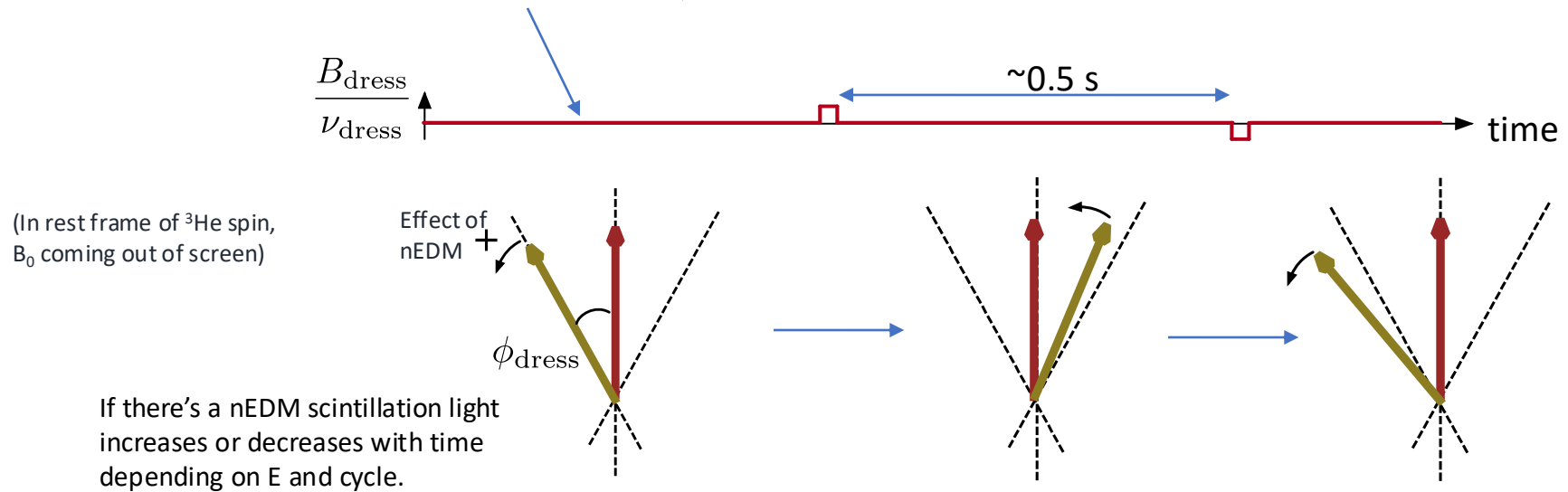


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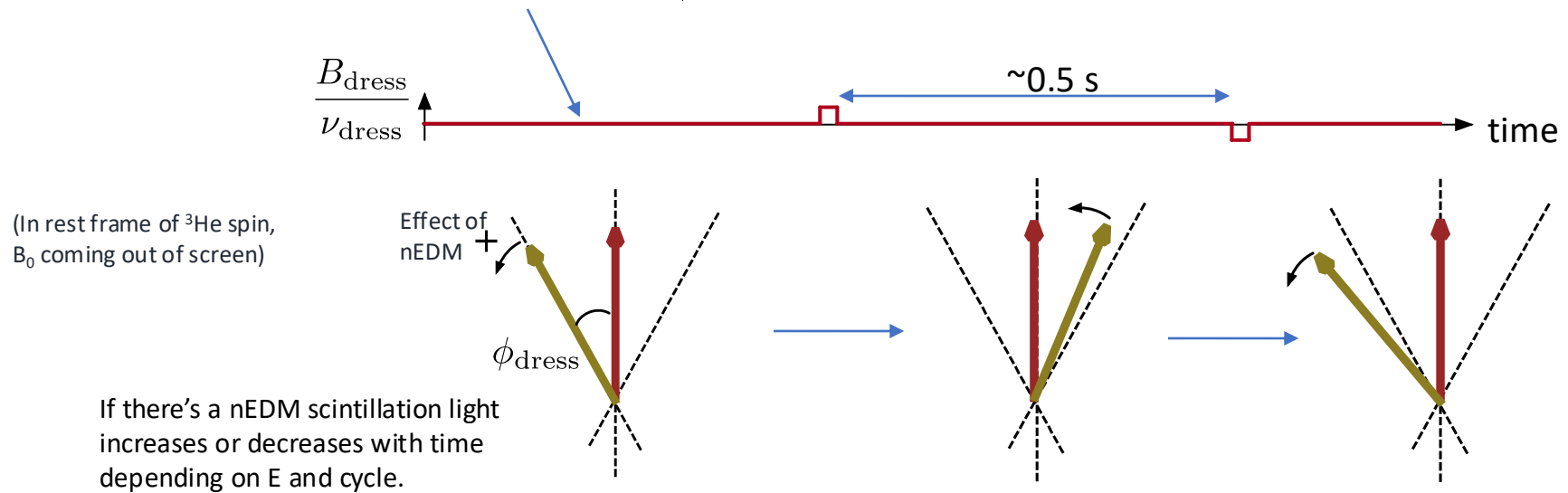
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- Example of modulation with “**square wave pulses**”. (Other modes possible.)

Critical dressing field to sit at a fixed ϕ_{dress}



- Can treat each pair as **asymmetry measurement**: $A_d(t_i) = \frac{N_- - N_+}{N_- + N_+}$ (Ahmed et al. JINST 2019)

Predicted sensitivity: 300 live days of data (e.g. 3 years running) $\sigma(d_n) = 1.7 \times 10^{-28} e \cdot \text{cm}$

- Less sensitive to static field in-homogeneities. Quality and control dressing field becomes main systematic.

nEDMSF apparatus/equipment status

Legacy Experiment Apparatus

Main subsystems

1. Superconducting magnet and shield: $B_0 = 1\text{--}3\text{ }\mu\text{T}$ + flip, trim & *dressing* coils ([Caltech](#), [UK](#), [ORNL](#), ...)
 - Large ($5 \times 5 \times 7\text{ m}^3$) magnetically shielded enclosure
2. Central detector system: acrylic measurement cells, HV, light collection, SQUIDS ([LANL](#), [ORNL](#), [Caltech](#), [Montclair](#), ...)
 - Dilution refrigerator I: $\sim 75\text{ mW}$ @ 250 mK
3. Polarized ^3He system: atomic beam source, injection, purification ([UIUC](#), [MIT](#), ...)
 - Dilution refrigerator II: $\sim 75\text{ mW}$ @ 250 mK
4. Cryostat vessels: aluminum, composite ([ORNL](#), ...)
5. Neutron transport ([UK](#), [ORNL](#), ...)
6. DAQ, slow controls ([ORNL](#), ...)
7. Simulation ([ORNL](#), [UK](#), [Caltech](#), [Montclair](#)...)
8. Systematics and operational systems apparatus ([NCSU](#), [Caltech](#), [Montclair](#) ...)

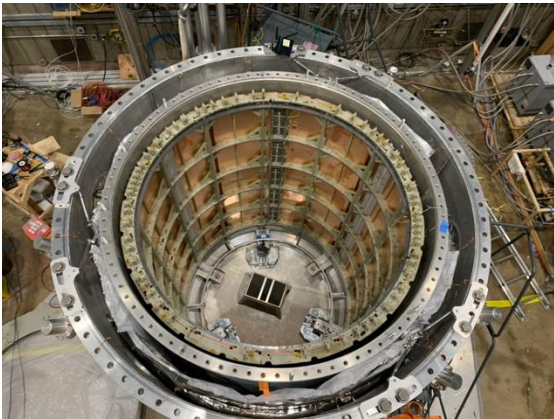
Hardware Status: Magnet System



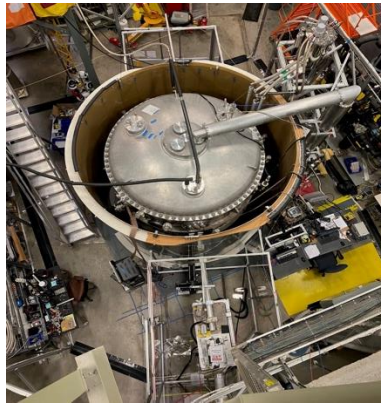
$B_0 \cos\theta$ coil



Magnet system closed for cold test at ORNL



Completed B_0 coil, sc shield, vacuum chamber at ORNL



Magnet system in cold beam test at ORNL (summer '23, '24)



Completed magnetic shielded room at Immedco: Hägendorf, CH

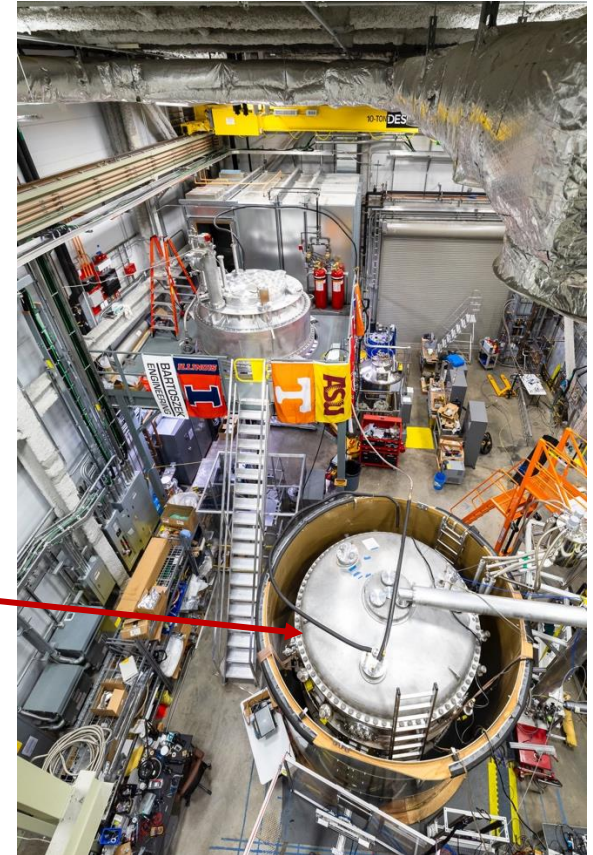
- To do: spin dressing coils

Spin-dress coil:

Brad's lab space yesterday



Oak Ridge National Lab



Hardware Status: Central Detector System



Completed cryostat top flange, tail for CDS testing



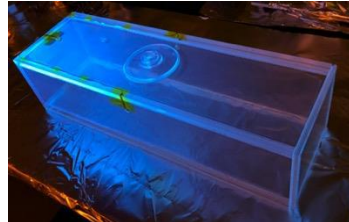
Half-scale HV test (T=0.4 K)



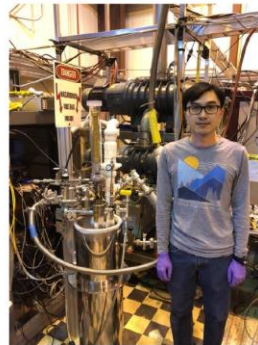
OVC composite vessel



Full-size stainless Cavallo electrodes

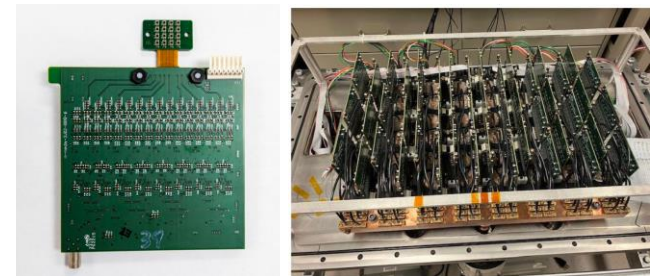
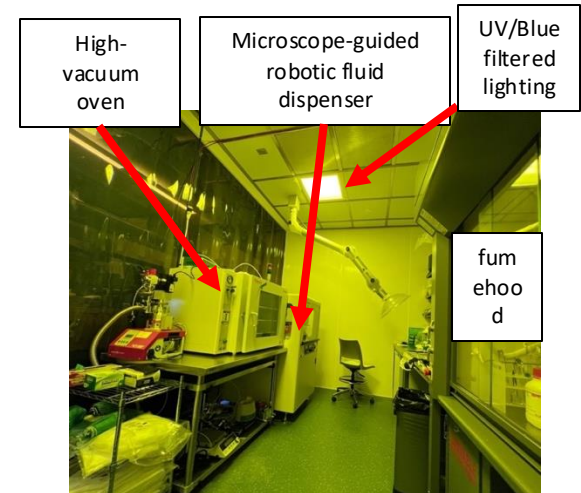


Measurement cell



Small scale HV

N. S. Phan, et al *J. Appl. Phys.* 129, 083301 (2021)



SiPM light collection

- To do: final electrode material choice & full-sized multilayered cells with thin windows

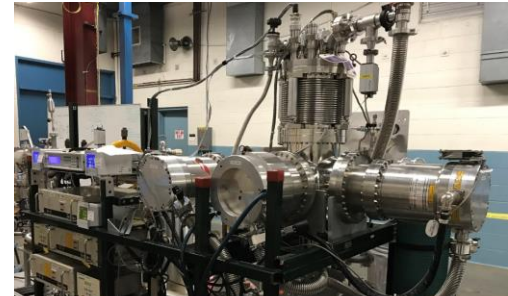
Hardware Status: Polarized ^3He System



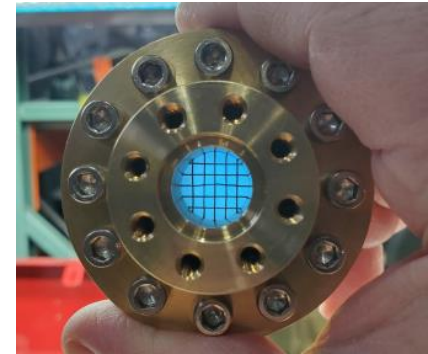
Dilution refrigerator/test cryostat



Dilution refrigerator + film burner
(75 mW at 250 mK)



Atomic beam source



Nucleopore filter: magnetic particles

- To do: injection system commissioning

Alternative physics:

Searching for axion-like particles (ALPs)
& improvement in neutron magnetic moment

time-oscillating nEDM-like signals from ALPs

- Axion-like particles couple with gluons to induce an oscillating nEDM signal [Abel et al., PRX 7 (2017)]:

$$\omega_n(t) = |\gamma_n|B_0 \pm \frac{2d_n|E|}{\hbar} + \frac{2|E|\alpha_{\text{ax}}}{\hbar} \cos(\omega_{\text{ax}}t + \phi_{\text{ax}})$$

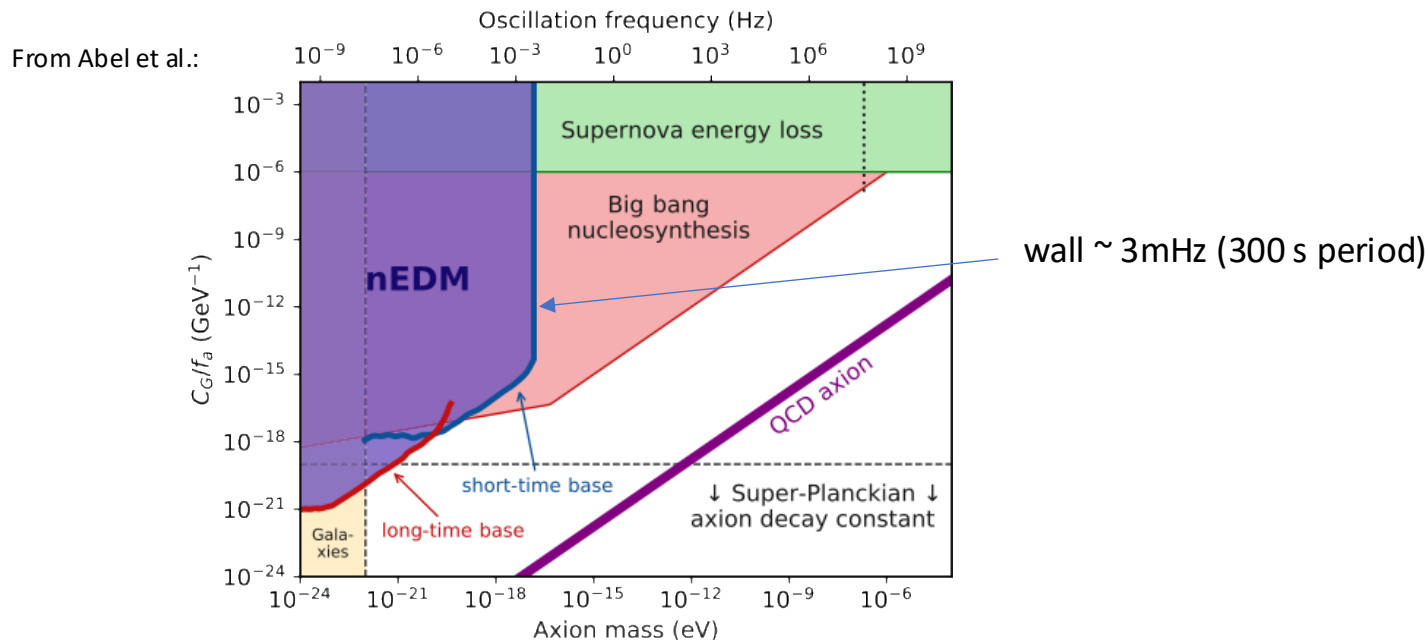
Amplitude of oscillation (units: e.cm) $\rightarrow \alpha_{\text{ax}}$

phase (free parameter in analysis) $\rightarrow \phi_{\text{ax}}$

Axion field coherently oscillates: $\omega_{\text{ax}} = m_a c^2 / \hbar$

axion mass $\rightarrow m_a$

- For “standard” Ramsey UCN search technique, axion sensitivity scales with nEDM sensitivity.
- Expected ~ 2 orders-of-magnitude improvement from “standard” Ramsey UCN search technique.



nEDMSF's ultra-short baseline searches

- $\phi_n(t)$ information allows searches for oscillating-nEDM signal at frequencies higher than the measurement cycling:

In free-precession mode
(can search up to ~ 10 Hz)

$$\phi_{3n}(t) = \left[(|\gamma_3| - |\gamma_n|) B_0 \pm \frac{2d_n|E|}{\hbar} \right] t + \frac{2|E|\alpha_{ax}}{\hbar\omega_{ax}} [\sin(\omega_{ax}t + \phi_{ax}) - \sin \phi_{ax}] + \phi_{3n0}$$

↑
determines n-3He scintillation light

In dressed-spin mode: each modulation cycle can be used to extract a nEDM value. Can search up to ~ 1 Hz.

nEDMSF's ultra-short baseline searches

- $\phi_n(t)$ information allows searches for oscillating-nEDM signal at frequencies higher than the measurement cycling:

In free-precession mode
(can search up to ~ 10 Hz)

$$\phi_{3n}(t) = \left[(|\gamma_3| - |\gamma_n|) B_0 \pm \frac{2d_n|E|}{\hbar} \right] t + \frac{2|E|\alpha_{ax}}{\hbar\omega_{ax}} [\sin(\omega_{ax}t + \phi_{ax}) - \sin \phi_{ax}] + \phi_{3n0}$$

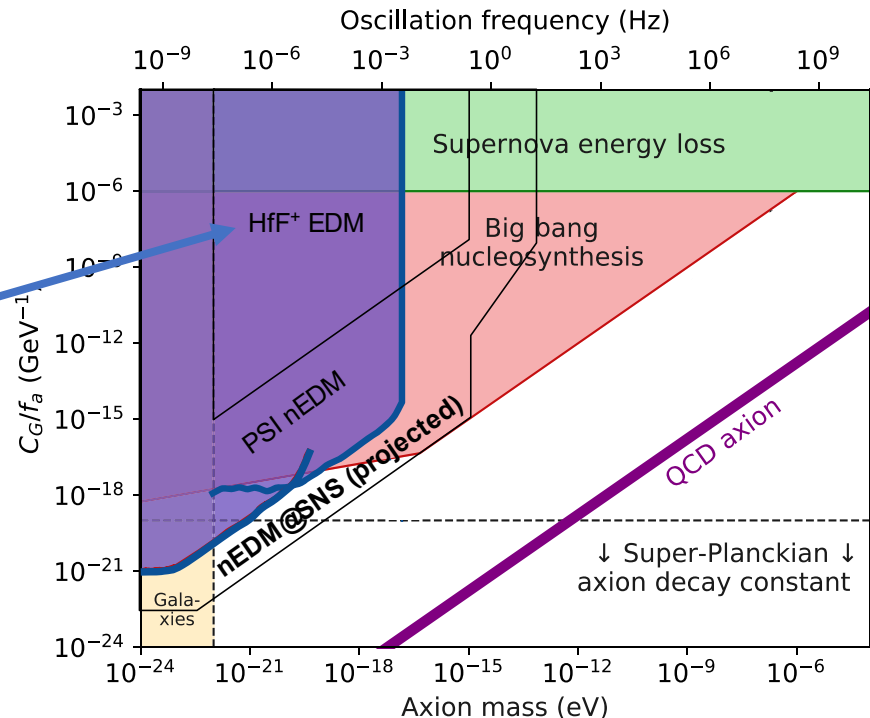
↑
determines n-3He scintillation light

In dressed-spin mode: each modulation cycle can be used to extract a nEDM value. Can search up to ~ 1 Hz.

Expected axion sensitivity exclusion from nEDMSF (from signal simulations & fit):

**"Experimental Constraint on Axionlike Particles
over Seven Orders of Magnitude in Mass"**

Roussy et al. PRL 2021



Measurement of $\gamma_n/\gamma_{3\text{He}}$ to improve γ_n

- Only minor improvements in knowledge of the neutron magnetic moment over the past 5 decades

PHYSICAL REVIEW D VOLUME 20, NUMBER 9 1 NOVEMBER 1979

Measurement of the neutron magnetic moment

G. L. Greene* and N. F. Ramsey + et al.
Harvard University, Cambridge, Massachusetts 02138

$$\frac{\sigma_{\gamma_n}}{\gamma_n} = 0.24 \text{ ppm via } \frac{\gamma_n}{\gamma_p} \text{ (Proton)}$$

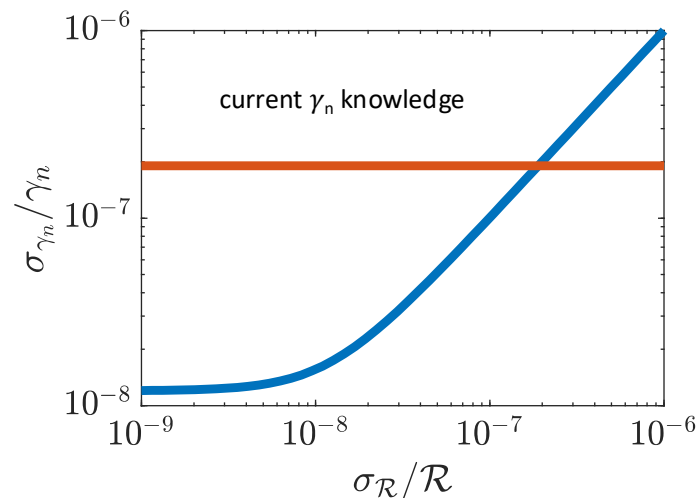
$$\left(\frac{\sigma_{\gamma_n}}{\gamma_n}\right)^2 = \left(\frac{\sigma_{\mathcal{R}}}{\mathcal{R}}\right)^2 + \left(\frac{\sigma_{\gamma_3}}{\gamma_3}\right)^2$$

$$\text{Afach et al., Phys. Lett. B 739 (2014) } \rightarrow \frac{\sigma_{\gamma_n}}{\gamma_n} = 0.19 \text{ ppm via } \frac{\gamma_n}{\gamma_{\text{Hg}}}$$

- In nEDMSF, we measure: $\mathcal{R} \equiv \frac{\gamma_n}{\gamma_{3\text{He}}}$ CODATA 2018: $\frac{\sigma_{\gamma_{3\text{He}}}}{\gamma_{3\text{He}}} = 1.2 \times 10^{-8}$

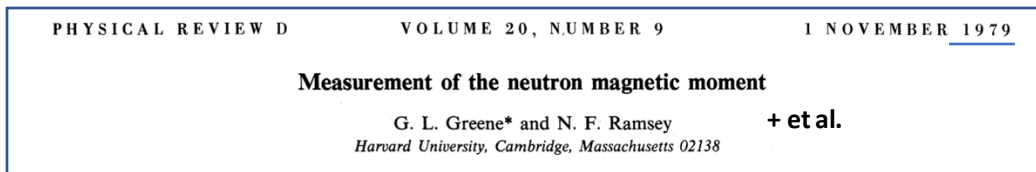
STATISTICAL ERROR:

- 11 days: $\frac{\sigma_{\mathcal{R}}}{\mathcal{R}} \approx \frac{\sigma_{\omega_{3n}}}{\omega_{3n}} \lesssim 1 \times 10^{-8} \text{ (0.01 ppm)} \text{ } (\sim 20\text{x improvement})$



Measurement of $\gamma_n/\gamma_{3\text{He}}$ to improve γ_n

- Only minor improvements in knowledge of the neutron magnetic moment over the past 5 decades



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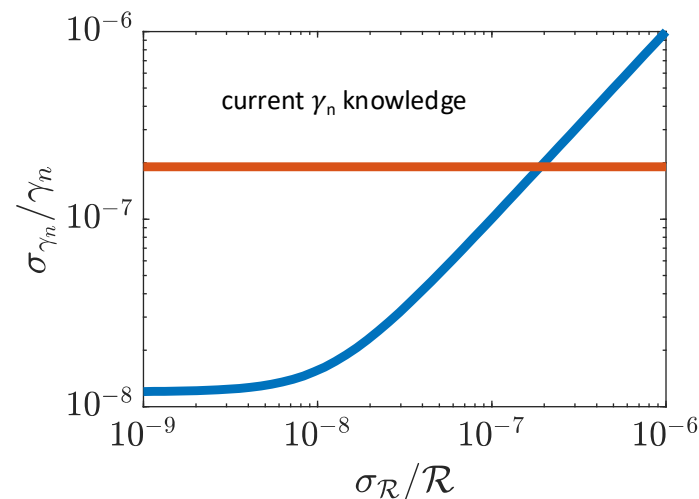
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BUT...SYSTEMATICS:

- Using comagnetometer for absolute field determination.
- UCN center-of-mass gravitational offset effect and knowledge of magnetic field gradients **very important** (need advanced magnetometry near cells)
- Does $\mu_{3\text{He}}$ renormalize when dissolved in superfluid ^4He ? I don't know. The ^3He has an effective (inertial) mass when in superfluid.

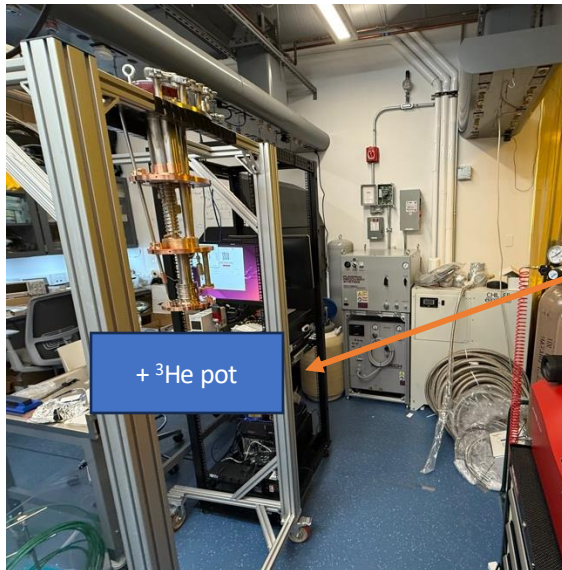
R&D for magnetometry near nEDMSF cell

- Currently ^3He **comagnetometer** and **fluxgate magnetometers** ~ 0.5 m away from cell (can't run during measurements due to noise and can't place too close due to distorting fields)
- Too cold for gas based magnetometry
- Technology gap for low-field **AND** low-temperature magnetometry

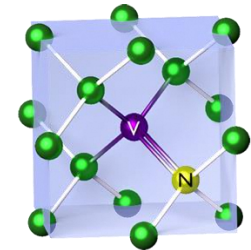
R&D for magnetometry near nEDMSF cell

- Currently ^3He comagnetometer and fluxgate magnetometers ~ 0.5 m away from cell (can't run during measurements due to noise and can't place too close due to distorting fields)
- Too cold for gas based magnetometry
- Technology gap for low-field AND low-temperature magnetometry
- Funding pending: development of near-cell magnetometers for nEDMSF

New 0.5 K platform @ Montclair



Read NV centers in diamond with SQUIDs

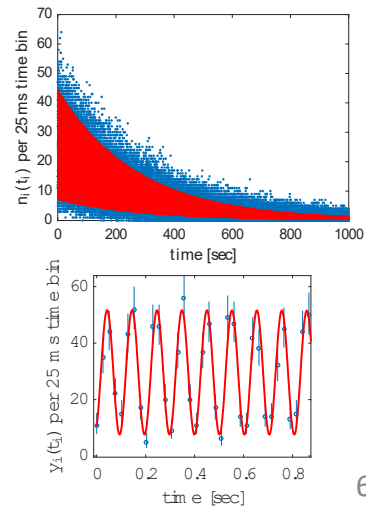
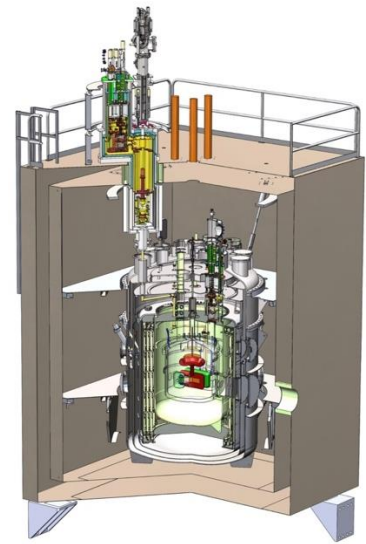


Hydrogen in plastic cell walls thermally polarized ($\mu B/kT$).
C.f. hospital MRI, $|B| \sim 10^6$ weaker, $T \sim 10^4$ lower \rightarrow highly detectable (especially with SQUIDs)



Summary

- There is a lot from the collaboration to cover, I apologize if I missed something for time's sake.
- The **cryogenic UCN + ^3He + superfluid** scheme offers many advantages to **reach 10^{-28} e.cm**
- In-situ produced UCN in **small cell with high-density and long storage**
- Supports high **electric field, SQUIDs** and **superconducting magnetic shielding**
- Can vary motion of our ^3He magnetometer with **small T changes to study key false EDM systematic effect**
- **Two measurement modes** with different systematic effects for **self-checking our own results**
- The live and in-situ UCN spin analysis is **a true frequency measurement**
- **New type of signals for the field**, on-going extensive work to understand its analysis
- Allows new **ultra-short baseline axion search to reach higher mass**
- **“Go east (relative to US), young man”**: new direction towards Europe.



Back up slides

Ultracold neutrons (UCNs)

- Slow neutrons undergo **scattering** from a nuclei via strong force: $\psi(\mathbf{r}) \propto \exp(i\mathbf{k} \cdot \mathbf{r}) + f(\theta) \frac{\exp(ikr)}{r} \rightarrow$ s-wave: $f(\theta) = -a$
- Coherent scattering off collection of nuclei: $\psi(\mathbf{r}) = \exp(i\mathbf{k}_0 \cdot \mathbf{r}) - \int \beta(\mathbf{r}') \psi(\mathbf{r}') \frac{\exp(ik_0 |\mathbf{r} - \mathbf{r}'|)}{|\mathbf{r} - \mathbf{r}'|} d^3r'$
- Apply Born approximation, where nuclei can be treated as δ -functions (nuclei size ~ 1 fm, where as slow neutron $\lambda > 1$ Å)
- Volume average gives **effective neutron “optical” potential**:

$$V_{\text{opt}} = \frac{2\pi\hbar^2}{m} \sum_i n_i b_i$$

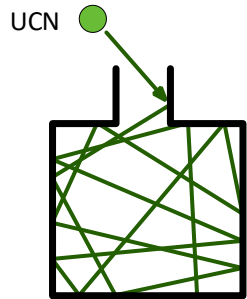
n_i ← nuclei-number density
 b_i ← bound coherent scattering length $b = \frac{m_n}{m} a$
 \sum_i ← sum over nuclear species in a material

If energy below V_{opt} of a material => total **external** reflection at all incident angles
 => **Ultracold neutrons** can be stored in a material “bottle”

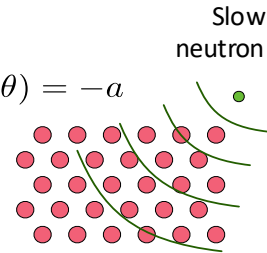
Material	V_{opt}
^{58}Ni	335 neV
Be	252 neV
Fluorocarbons	~100 neV
Al	54 neV
polyethylene	-9 neV
d8-polystyrene	160 neV

Idea: Fermi (1946); Published: Zeldovich (1959);
 Realization: Shapiro et al., Dubna (1968) & Steyerl et al., Munich, (1969)

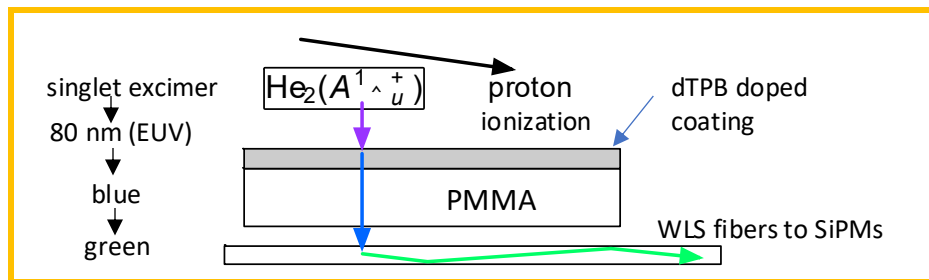
$$E \lesssim 300 \text{ neV}, \lambda \gtrsim 60 \text{ nm}, v \lesssim 7 \text{ m/s}, \text{“temperature”} \lesssim 2 \text{ mK}$$



- Typically 10^{-4} loss probability per reflection $\Rightarrow R_{\text{reflect}} \approx (v A_{\text{walls}}) / (4 \times \text{Volume})$
- Store UCNs with time constant (including β -decay loss) $\tau_{\text{tot}} \lesssim 100 - 600 \text{ s}$



Superfluid helium as a scintillator



- **Light signal gives live angle of neutron spin relative to ³He spin!**
 - **UCNs produced and spin analyzed in the cell** (i.e. no transport loss)
 - ~ 100% detection efficiency
 - more photo-electrons (PE) = better discrimination against backgrounds
 - Suppression of EUV in electric field: Ito et al. PRA (2012)

Light collection efficiency test apparatus @ ORNL with β -source.
Projected > 19 PE in final setup



The pseudomagnetic field

- The neutron & ^3He scattering length is spin-dependent
- Recall:

$$V_{\text{opt}} = \frac{2\pi\hbar^2}{m} \sum_i n_i b_i$$

number density n_i
 bound coherent scattering length b_i
 sum over nuclear species in a material

$b = \frac{m_n}{m} a$

- So the optical potential the UCNs are in becomes:

$$V_{\text{opt,tot}} = V_{\text{opt},^4\text{He}} + V_{\text{opt},^3\text{He}}$$

18.5 neV
 10^{-9} neV + const ($S_3 \cdot S_n$)
 (for $x_3 \sim 10^{-10}$)

- Recall: $V_{\text{mag}} \sim \mu \cdot B$
 so this effect is like a magnetic field
 ("pseudomagnetic field")

$$\vec{H}_{\text{pseudo}} = -\frac{4\pi\hbar}{m_n \gamma_n} \rho_3 b'_i \sqrt{\frac{I}{I+1}} \vec{P}_3$$

^3He density ρ_3
 Incoherent n- ^3He scattering length b'_i
 $\sim 0.1 \mu\text{G}$ ($\sim 10 \text{ pT}$)
 (for $x_3 \sim 10^{-10}$, $P_3 \sim 1$)

The pseudomagnetic field

- The neutron & ^3He scattering length is spin-dependent

- Recall:

$$V_{\text{opt}} = \frac{2\pi\hbar^2}{m} \sum_i n_i b_i$$

$b = \frac{m_n}{m} a$

Annotations for the equation above:

- n_i : number density
- b_i : bound coherent scattering length
- \sum_i : sum over nuclear species in a material

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$$V_{\text{opt,tot}} = V_{\text{opt},^4\text{He}} + V_{\text{opt},^3\text{He}}$$

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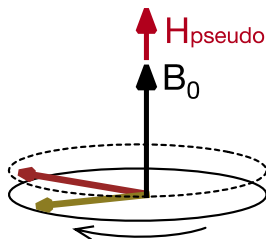
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- ρ_3 : ^3He density
- b'_i : Incoherent n- ^3He scattering length



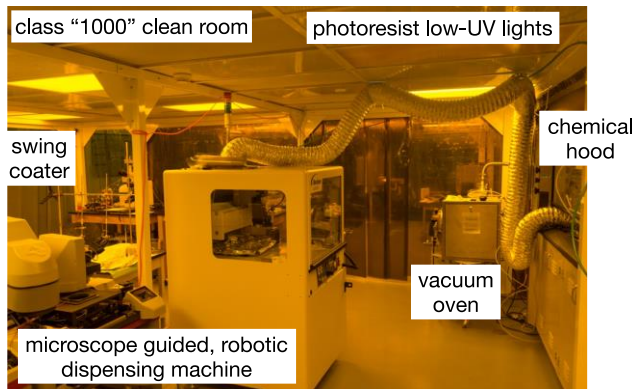
- ^3He does not see a pseudomagnetic field from UCNs
- After $\pi/2$ pulse, need ^3He spin accurately in perpendicular plane
- Since there are two cells, any frequency drifts cancel out
- Can add statistical noise between measurements though
- Once in perp. plane, frequency shifts are small because it's a Bloch-Siegert shift

Measurement cells

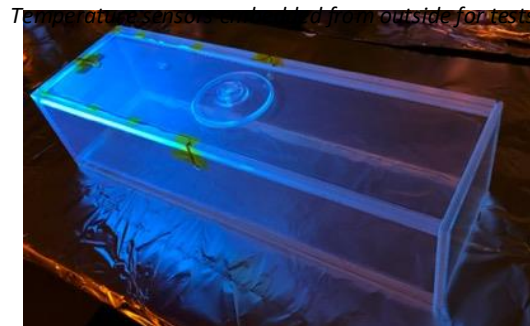
- Dimensions optimized for statistics and systematics. (e.g. E-field, diffusion-limit L^{-4} dependence of T_2 time & L^2 dependence of false EDM, magnetic field gradients, etc.).
- Cells separate from electrodes (cooling needed for heat flush)

Inner coating	~ 1 μm thick d-TPB embedded in d-PS matrix	Non-magnetic and not electrically conductive. Deuteration for high Fermi potential (165 neV from neutron reflectometry). TPB for detection of 80 nm ^4He scintillation light (see Ito, Cianciolo, Loomis talks)
Side walls	~ 1 cm thick p-PMMA plates	Optical photon transmission, mechanical strength, and purity. “Swing coating” on flat plates produce surface finish ~ 5 nm RMS roughness (AFM measured)
End windows	~ 5 mm thick d-PMMA	reduce activation caused by scattering of 9 Å beam
Sealed cell	deuterated solvent cemented	to separate 10^{-10} polarized ^3He inside cell from unpolarized natural-abundance ^3He
Cell hole	~ 1 cm opening	initial ^4He filling then loading/unloading ^3He via heat flush. Low UCN loss.

Cell production facilities in clean room



Full-sized test measurement cell (illuminated by 360 nm UV lamp).



*p- & d- refers to “protonated” or fully-deuterated versions of a material.

PMMA = poly(methyl methacrylate), more commonly-known by tradenames acrylic or plexiglass.

PS = polystyrene.

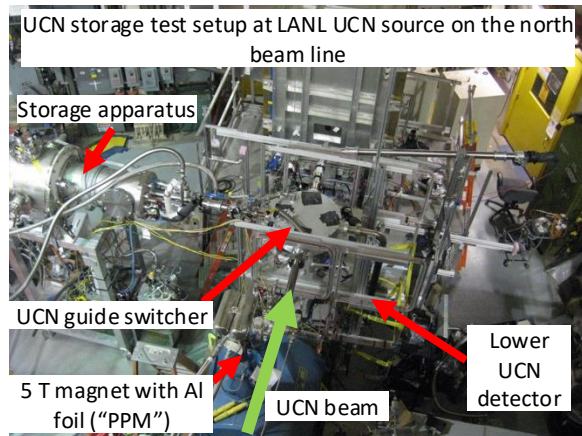
TPB = tetraphenyl butadiene, an electroluminescent dye

Cryogenic UCN storage tests with external UCN source

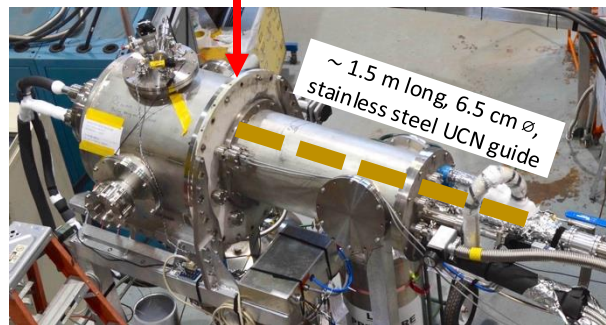


Thanks to collaborators: Cianciolo, Clayton, Cooper, Curie, Golub, Griffith, Huffman, Ito, Korobkina, Makela, McDonald, O'Shaughnessy, Pentilla, Ramsey, Smith, Stanislaus, Tang, Weaver, Wei

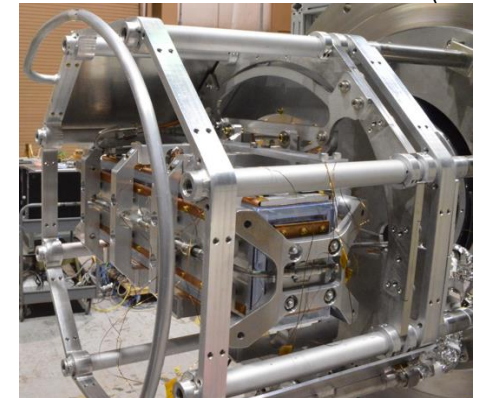
- A **cryogenic UCN storage apparatus** used with the **Los Alamos National Lab. solid-D₂ UCN source**.
- **Standard fill-and-empty setup** with cell valve & UCN guide switcher to connect (1) UCN source to cell or (2) cell to lower UCN detector
- **Cells cooled by spring-loaded cooling plates coupled to flowing L⁴He line**. Cells reach nominal 30 K in ~ 3 days (but can reach 15 K with high L⁴He rate).
- Dry & clean vacuum system: in vacuum outside cell, $P \sim 5 \times 10^{-7}$ mbar before starting to cool cell. Outgassing dominated by plastic.
- Pumping inside cell restricted due to 1-cm cell hole & UCN guides. Put UCN guide switcher in “intermediate” position for increased cell pumping speed.
- **Cells are heated** to 50 degC (coating limit) and **pumped for 10-14 days** before cooling.
- The input L4He flow is split, **a region of stainless steel UCN guide ~20 cm away from the cell is maintained to be coldest point in the system** (nominally > 20 K colder than cell). This SS surface **acts as a cryopump** for pumping out cell and **traps condensable-contamination** from reaching inside cell.



Slide-openable vacuum vessel with flowing LN₂ radiation shield (80 K & ~ 3π sr coverage)



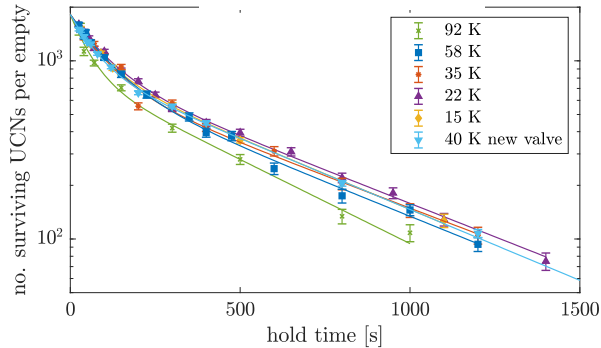
Spring loaded cooling plates used to maintain contact with cell
Return L4He cools another radiation shield (50 K)



UCN storage results

(Data normalized so number of UCNs at $t=0$ extrapolated with double exponential fit are matched. The value chosen at $t=0$ is average of the different storage curves with same cell, measured over consecutive days)

2nd generation cell #2



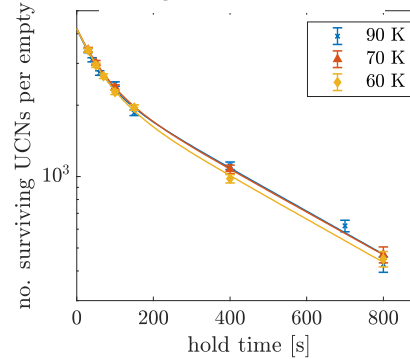
- Beam position: behind UCNr
- Measured with old and new valve systems

Fitting 40 K data:

$$N_{\text{short}} = (50 \pm 2) \%, \tau_{\text{short}} = (73 \pm 7) \text{ s}$$

$$N_{\text{long}} = (50 \pm 2) \%, \tau_{\text{long}} = (546 \pm 16) \text{ s}$$

2nd generation cell #3



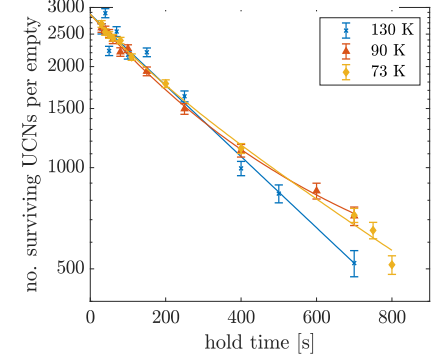
- Beam position: north beam line
- dPMMA valve system

Fitting 60 K data:

$$N_{\text{short}} = (45 \pm 4) \%, \tau_{\text{short}} = (66 \pm 16) \text{ s}$$

$$N_{\text{long}} = (55 \pm 4) \%, \tau_{\text{long}} = (477 \pm 47) \text{ s}$$

2nd generation cell #4



- Beam position: north beam line
- dPMMA valve system
- Reduced dust contamination during cell production

Fitting 73 K data:

$$N_{\text{short}} = (24 \pm 28) \%, \tau_{\text{short}} = (170 \pm 190) \text{ s}$$

$$N_{\text{long}} = (76 \pm 28) \%, \tau_{\text{long}} = (590 \pm 190) \text{ s}$$

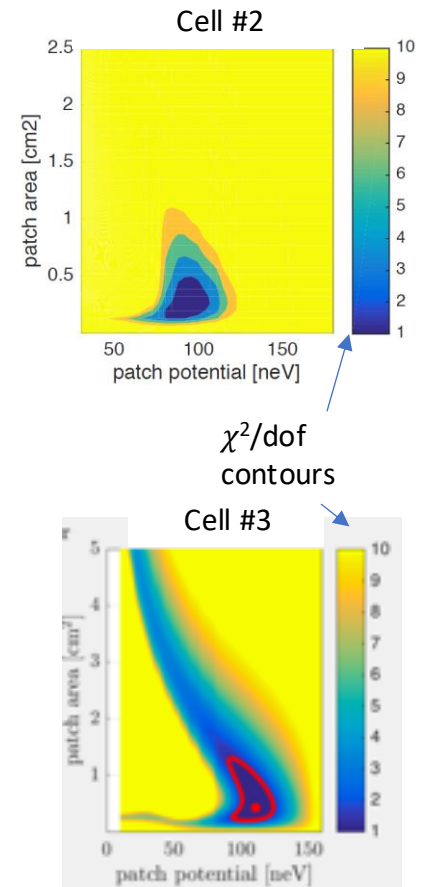
OR

$$\tau_{\text{single-exp}} = (490 \pm 11) \text{ s}$$

- Either sum of two exponential decays ("long" and "short") or single-exponential fit.
- When sum of two exp decays, strong correlations between fitted parameters
- When cell not cooled, $\tau_{\text{single-exp}} = (130 \pm 5) \text{ s}$
- Storage times here are β -decay + cell wall losses

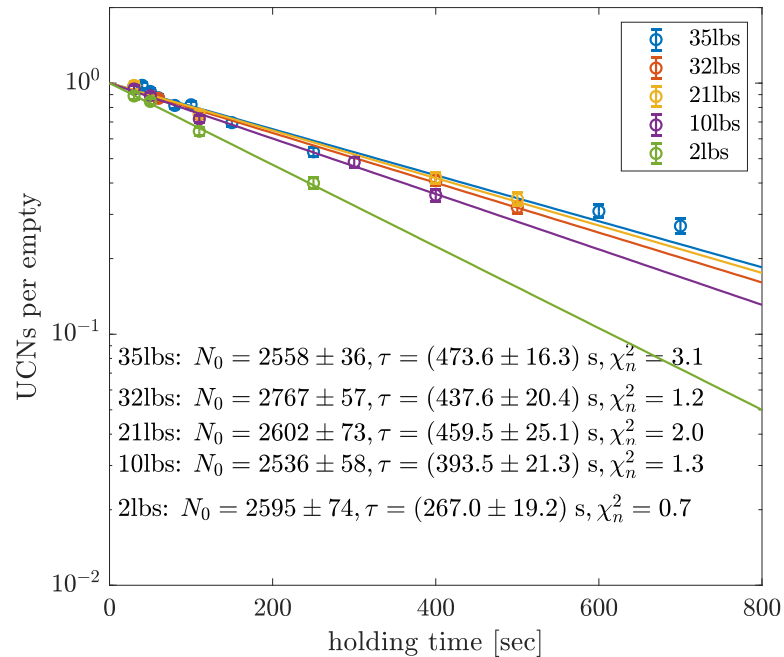
Interpretation of results

- High quality set of cryogenic UCN storage data with a new type of deuterated polymer coated surface ($V_F = 165$ neV, fluorescent, ^3He polarization friendly) exhibiting excellent UCN storage properties. Cell #4 with $\tau_{\text{single-exp}} = (490 \pm 11)$ s \Rightarrow **loss per reflection $\sim 2 \times 10^{-5}$**
- **Standard UCN-energy-dependent wall loss theory** (i.e. $f = W/V_F$ description) **insufficient** for describing data exhibiting **strong double-exponential decay**.
- From crude information on UCN spectrum, we deduce it is the higher-energy UCNs in the τ_{short} component. Commonly explained away by above-trapping threshold UCNs but the time scales here are very long for our rectangular cells.
- Most promising model found so far has been fitting with a small area “patch” (~ 0.3 cm²) with **low V_F** (~ 100 neV) exposed to UCNs
- Patch parameters very similar between cells #2 & #3
- Possible theories for cell #4:
 - (1) τ_{short} component very short, so not visible since shortest hold time is 25 sec (patch area $> 3\times$ larger)
 - (2) τ_{short} component very long or non-existent (difficult to fit for. Would need patch $5\times$ smaller)
 - (3) higher UCN energies not loaded into the cell in the first place.
- Currently developing UCN tracking simulations to increase understanding. Possibly use UCN spectrometer in future experiments.



dPMMA UCN cell valve results

- These results show we have **learned to implement a UCN valve with negligible UCN loss**
- The leak tightness of the valve is more stringent for containing ^3He inside the cell in final experiment ($v \sim 30$ m/s)



9 Å neutron beam induced cell activation in final experiment

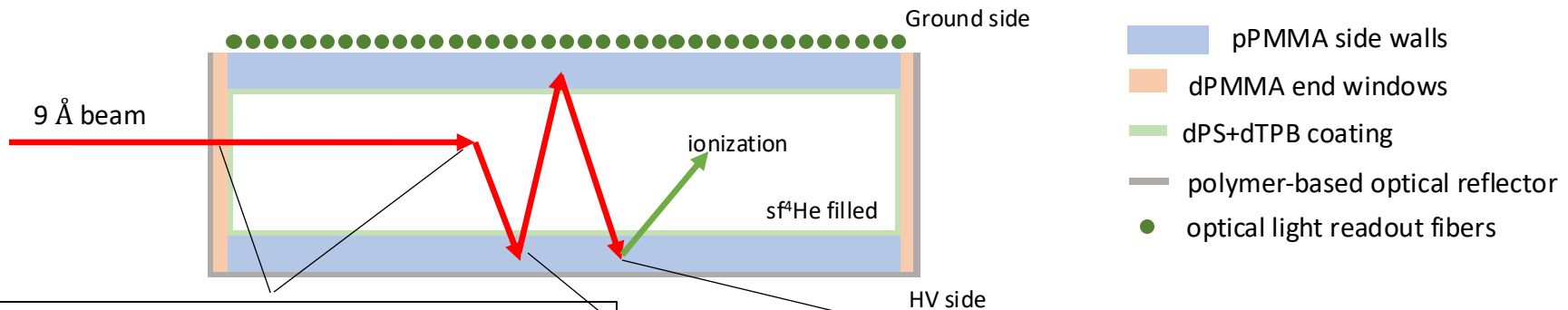
- In-situ UCN production via super-thermal production in sf^4He by 9 Å beam (via choppers) offers many advantages:

High UCN density $\sim 150 \text{ UCN/cm}^3$
after 1000 s of “filling time”

No UCN or depolarization loss from transport

Spatial spread of beam and filling time reduces phase-space evolution at start of precession measurements

- The beam intensity is $5 \times 10^5 \text{ s}^{-1}$ in each cell (UCN production $P = 0.31 \text{ UCN/cm}^3/\text{s}$), so $T_{\text{fill}} = 1000 \text{ s} \rightarrow 5 \times 10^5$ neutrons incident



Scattering in sf^4He (coherent inelastic scattering).

Mean-free-path = 20 m, 40 cm cell, 2%.

OR

Scattering in dPMMA windows (mix coherent elastic scattering & incoherent elastic scattering). $\sim 1\%$ depending on window thickness

OR

Direct activation of impurities in dPMMA with decay chains

Note: we developed L^4He MCNP neutron scattering kernels to 4 K. (Thanks to C. Lavelle from Johns Hopkins!)

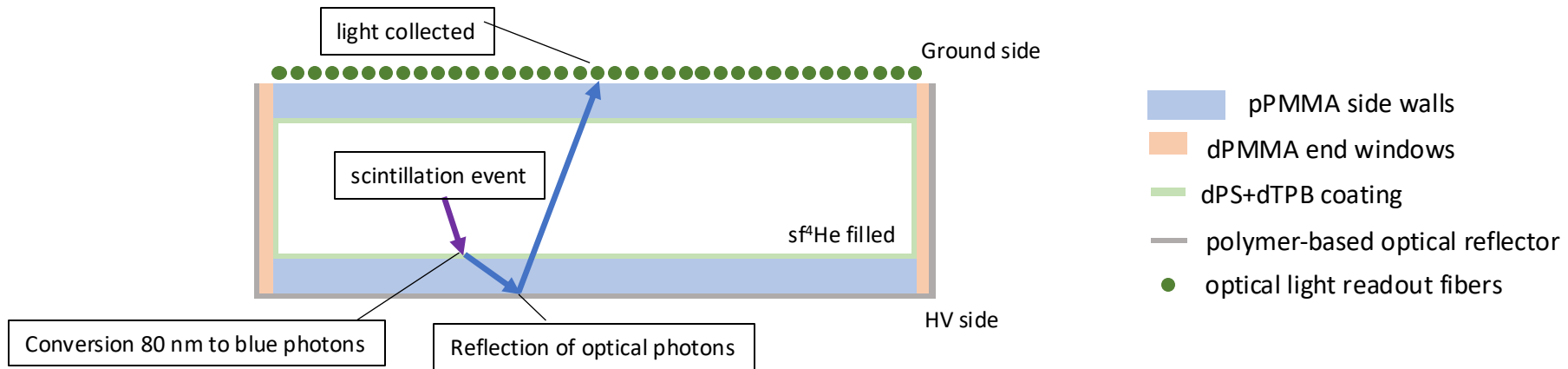
Cold neutrons
bounce around via
incoherent elastic off
 ^1H

Eventually $n+p \rightarrow D + \gamma$ (2.2 MeV) + other prompt-reactions
 \Rightarrow can produce free charge in cell
OR
 $n + \text{contaminant} \rightarrow$ short or long lived decay chain \Rightarrow delayed background scintillation events

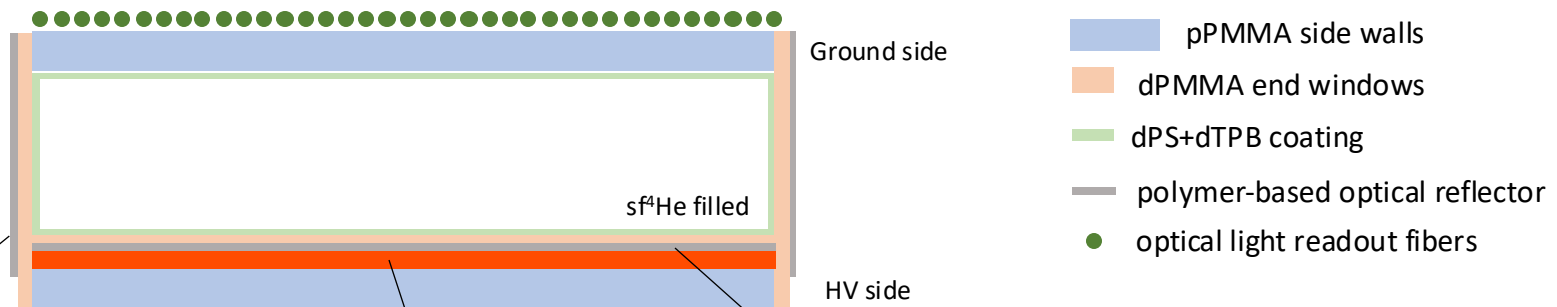
(see Korsch & Loomis talks)

How scintillation light is collected in measurement cell

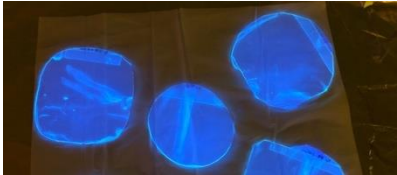
(see Cianciolo talk)



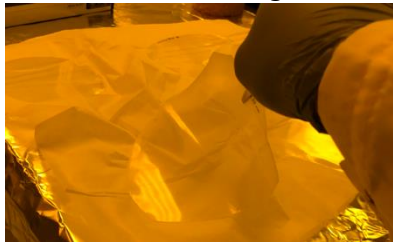
Solution to problem



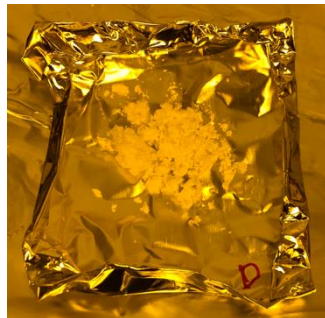
Coated thin end PMMA windows



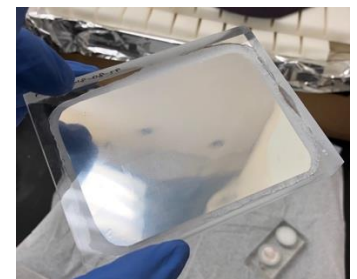
LN2 dunk and strength tests



In-house synthesis of high-purity dPMMA



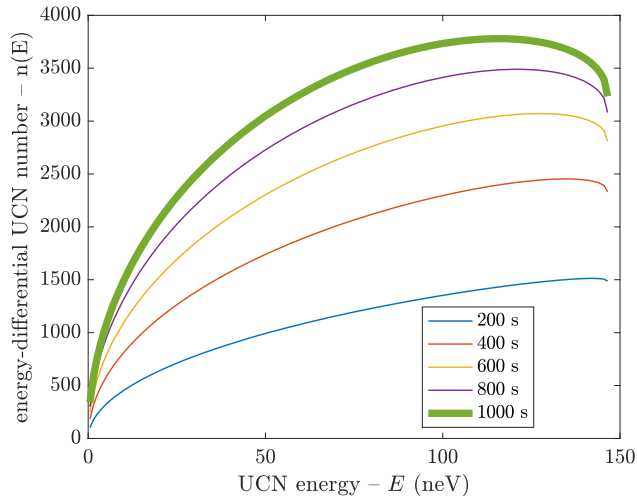
1/3rd scale prototype side wall with embedded LiOH absorber and optical reflector



Evolution of UCN spectrum during filling and precession

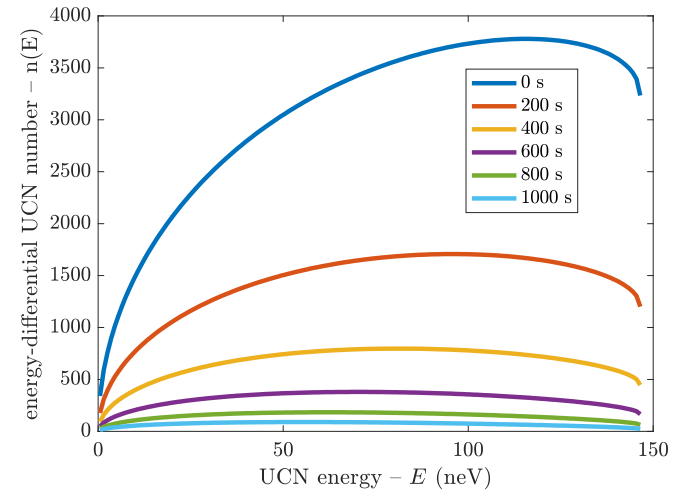
Example: used $f = W/V = 0.8 \times 10^{-5}$

UCN spectrum during “filling” (low ^3He absorption)



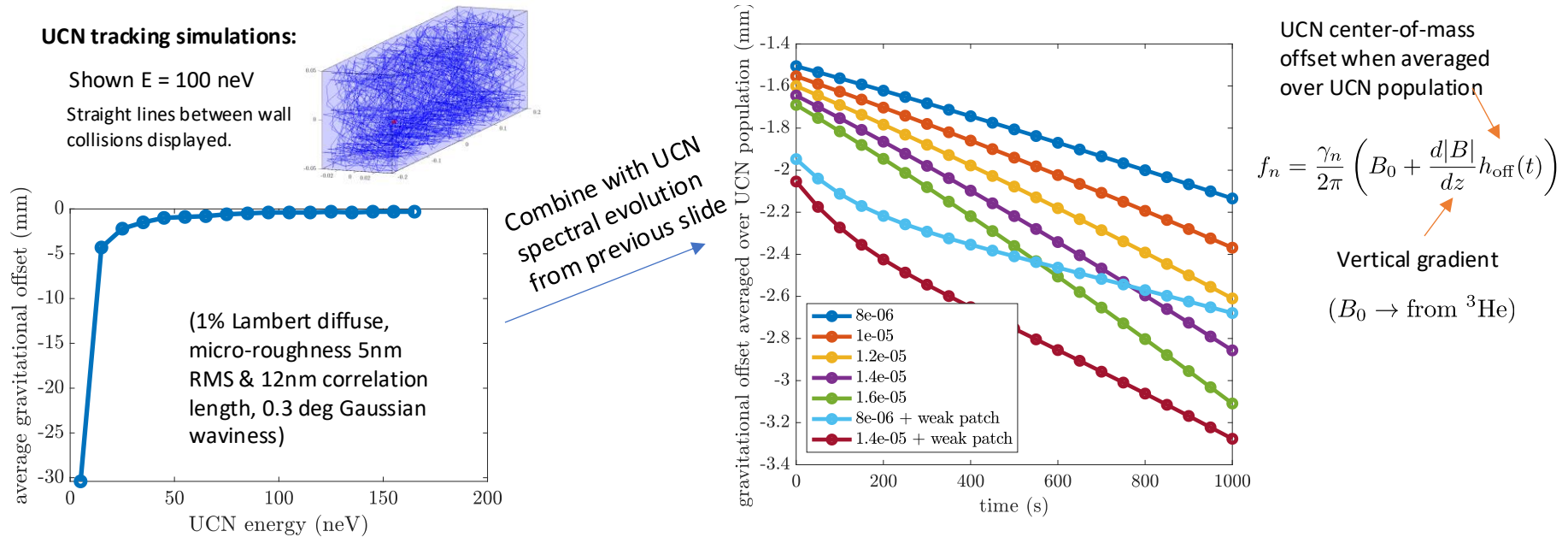
Statistically optimized $T_{\text{fill}} = 1000$ s:

UCN spectrum evolution during free precession measurement time (includes $\tau_3 = 500$ s)



- Produced **UCN spectrum in sf^4He is well-described**. Transport in guides not so much.
- Above assumes mechanical equilibrium. **Phase-space evolution will be small in nEDM@SNS** (3 L cell, with UCNs produced with approximately isotropic momentum and close to uniformly throughout cell, filled over 1000 s). Next step is to confirm with simulations
- These are the UCN spectra inside the cell. **Since spin analysis is in-situ, no need to correct for UCN-energy dependent transport loss** (and depolarization loss) during transport to interpret any UCN spectral measurements
- Change filling time (with reduction on statistics) to change initial UCN spectrum slightly for systematics

Use time-evolution of UCN gravitational offset to study UCN spectrum



- Apply known gradient, measure $f_n(t)$ to get information on the time-evolution of average UCN energy during precession.
- **Combine with total number UCN storage time measurement:** just watch β -decay without loading ^3He .
- More direct than spin-echo since since measuring precession frequency
- Phase-space evolution (which will impact this systematic effect as well as others!) less in the nEDM@SNS experiment
- A good chance for controlling this important systematic effect! Currently in early days of development still.

Photoelectron spectrum

