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After $(g - 2)_\mu$, what next for $U(1)_{L_\mu - L_\tau}$?

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Based on work with M.W. Li, X.G. He, X. Chu, and TRIDENT collaboration

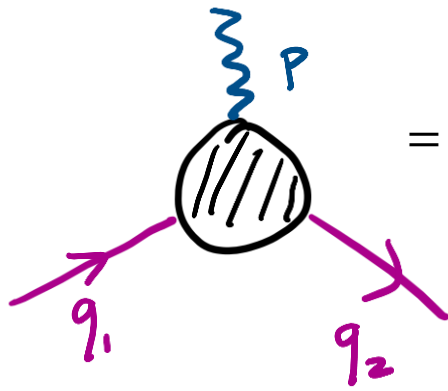
$g - 2$: 80 years of precision

From the Dirac equation one can derive magnetic moment of a Fermion

$$\vec{\mu} = g \frac{q}{2m} \vec{S}$$

Where $g = 2$, q is the electric charge $\vec{S} = \frac{1}{2} \vec{\sigma}$. Experimental hints that $g \neq 2$ in 1947!

Comes from $F_{\mu\nu} \sigma^{\mu\nu}$ term \Rightarrow quantum corrections will come from loops with this form.



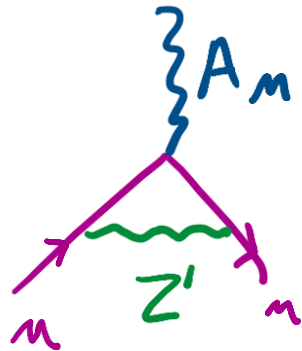
$$= (-ie) \bar{u}(q_2) \left[F_1 \left(\frac{p^2}{m^2} \right) \gamma^\mu + \frac{i\sigma^{\mu\nu}}{2m} p_\nu F_2 \left(\frac{p^2}{m^2} \right) \right] u(q_1) \quad \Rightarrow g \rightarrow 2 + 2 F_2 \left(\frac{p^2}{m^2} \right)$$

1-loop order $F_2(0) = \frac{\alpha}{2\pi}$ (Schwinger 1948, then Feynman and Tomonaga)

Vital for development of QFT!

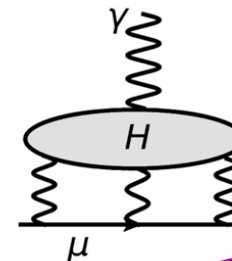
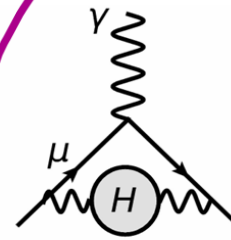
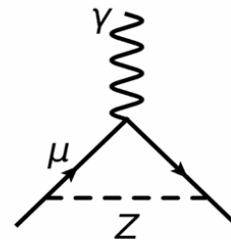
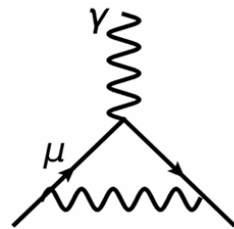
The muon ($g - 2$)

Measuring $(g - 2)_\mu$ is a clean test of new physics, and contribution is enhanced compared to electron $\Delta a_\mu \propto \left(\frac{m_\mu}{m_{NP}}\right)^2$



$$\rightarrow \Delta a_\mu^{Z'} = \frac{\tilde{g}^2}{8\pi^2} \times \left(\frac{m_\mu}{m_{Z'}}\right)^2 \int_0^1 \frac{2x^2(1-x)dx}{1-x + \left(\frac{m_\mu^2}{m_{Z'}^2}\right)x^2} \xrightarrow{m_{Z'} \gg m_\mu} \frac{\tilde{g}^2}{12\pi^2} \times \left(\frac{m_\mu}{m_{Z'}}\right)^2$$

SM contribution calculated to many loop level

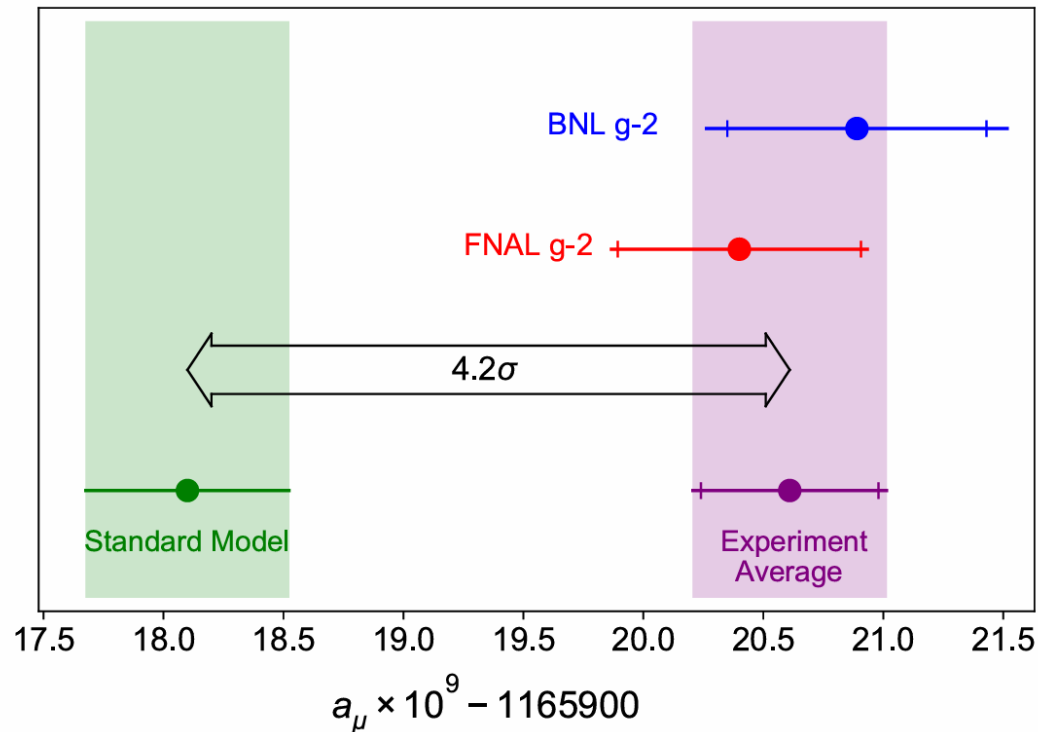


Largest error

Hints of new physics from $(g - 2)_\mu$

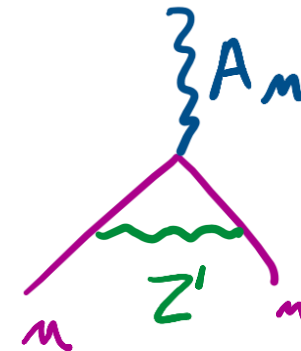
Since 2006, E821 at BNL saw hints of deviations from SM.

In 2021, Fermilab Muon $g - 2$ experiment started to confirm this



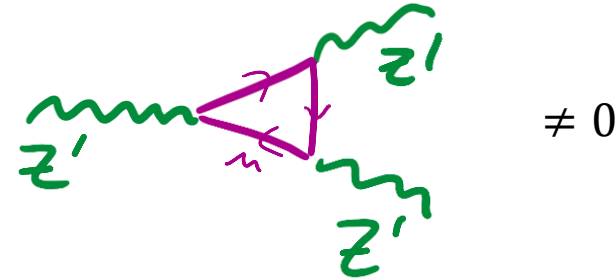
New physics to explain this is varied

The simplest is by introducing only 1 new particle



The $U(1)_{L_\mu - L_\tau}$ extension

Coupling only to the muon is anomalous, i.e.



We can either add more particle content

- Solutions lead to flavor changing constraints (see Carena et.al. [arXiv:1908.04818](https://arxiv.org/abs/1908.04818))

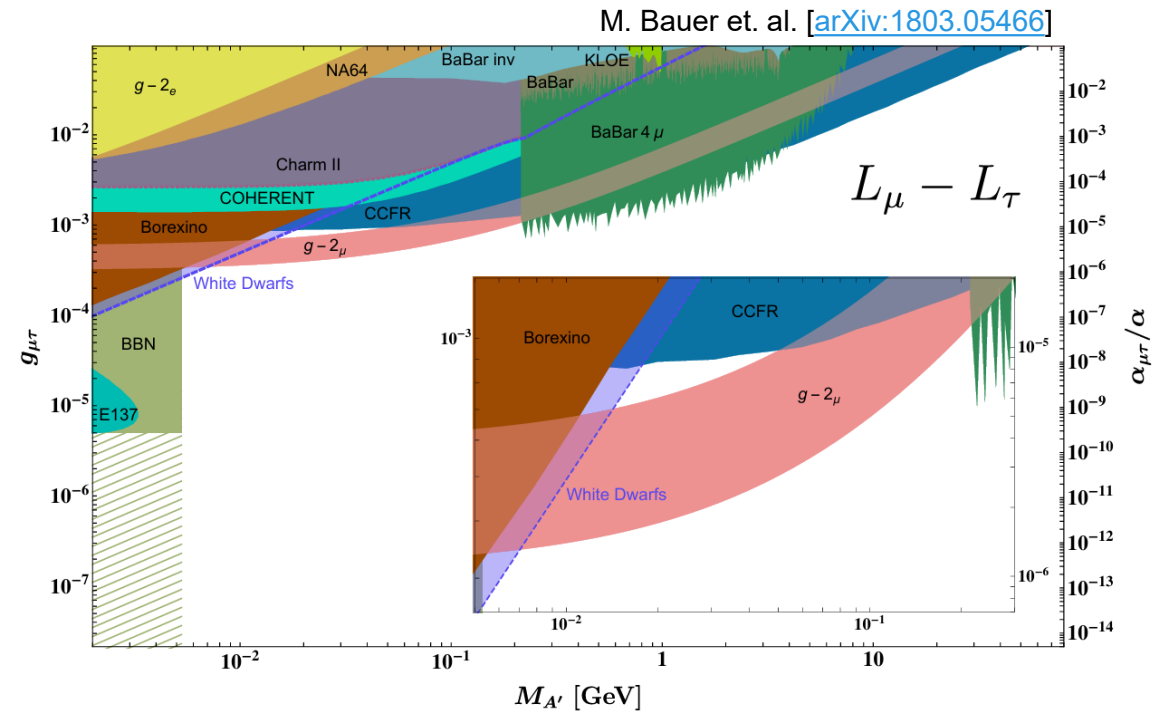
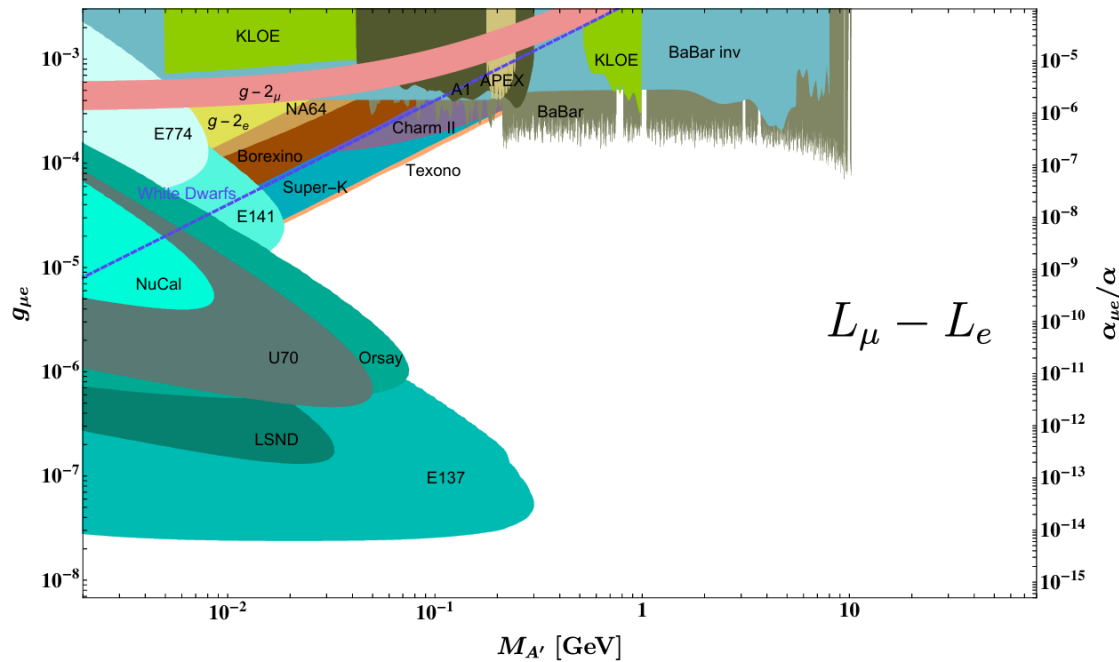
More simply, look for anomaly-free models

- Only two that involve muon coupling (see XG He et. al. [PRD 1991](https://arxiv.org/abs/hep-th/9905221)).

$U(1)_{L_\mu - L_\tau}$ or $U(1)_{L_\mu - L_e}$ where the gauge current is

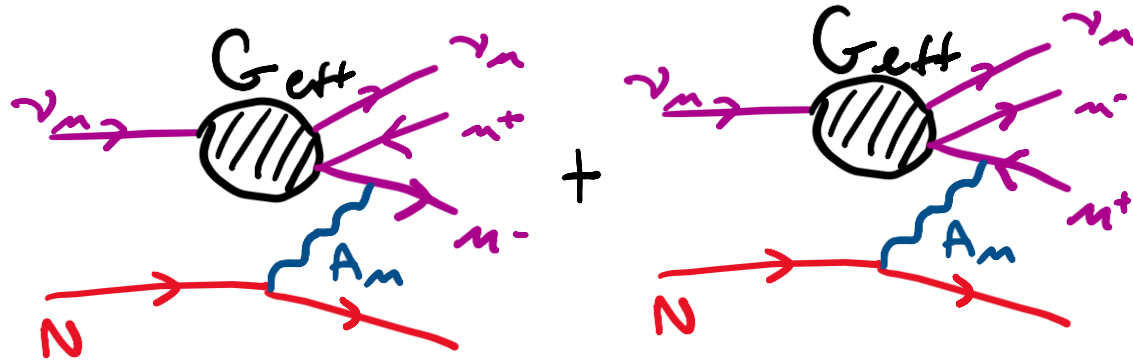
$$j_\alpha^{L_i - L_j} = \bar{L}_i \gamma_\alpha L_i + \bar{\ell}_{i,R} \gamma_\alpha \ell_{i,R} - \bar{L}_j \gamma_\alpha L_j - \bar{\ell}_{j,R} \gamma_\alpha \ell_{j,R} \quad \text{with } L_i = \begin{pmatrix} \nu_{i,L} \\ \ell_{i,L} \end{pmatrix}$$

The $U(1)_{L_\mu - L_\tau}$ extension



Focus on low Z' mass

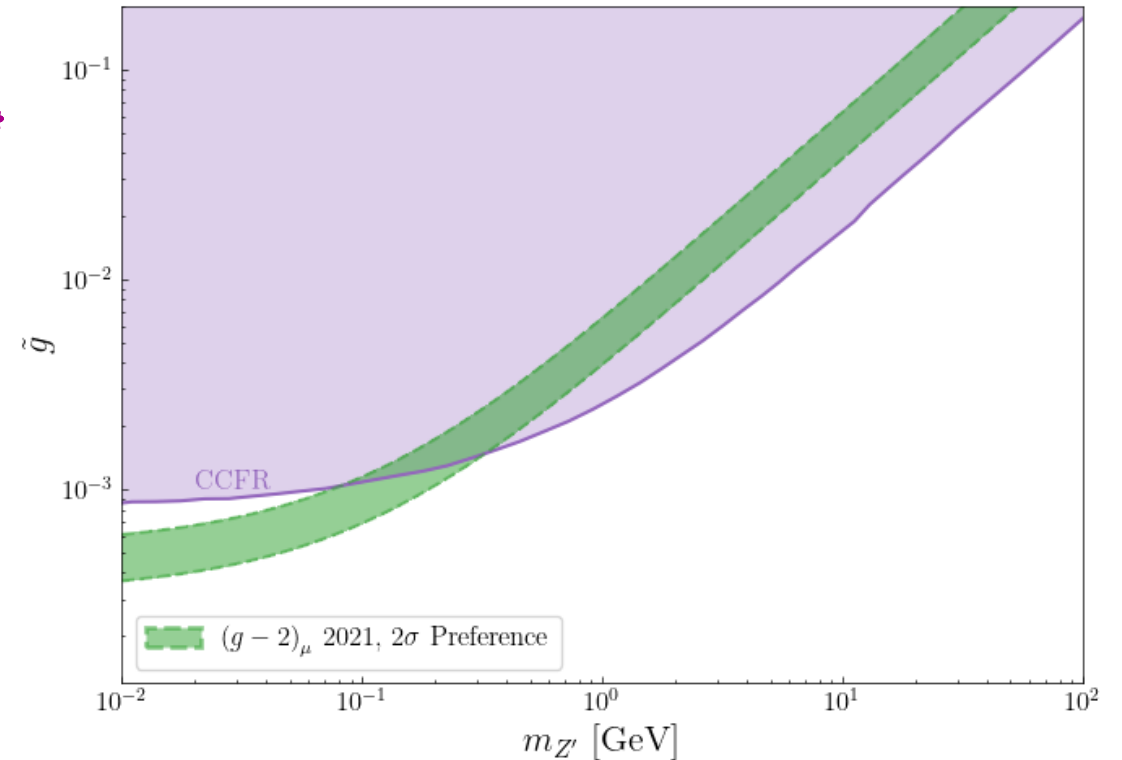
Muon Neutrino Trident (MNT) measurements limit how heavy Z' can be



$$\mathcal{L}_{\text{total}}^{\text{eff}} = -g_V \left(\bar{\nu}_\mu \gamma^\mu P_L \nu_\mu \right) \left(\bar{\mu} \gamma_\mu \mu \right) - g_A \left(\bar{\nu}_\mu \gamma^\mu P_L \nu_\mu \right) \left(\bar{\mu} \gamma_\mu \gamma^5 \mu \right)$$

where $g_V = \frac{G_F}{\sqrt{2}} (1 + 4 \sin^2 \theta_W) + G_{Z'}$ and $g_A = -\frac{G_F}{\sqrt{2}}$

$$\left. \frac{\sigma_{Z'}}{\sigma_{\text{SM}}} \right|_{\text{trident}} = \frac{(1 + 4s_W^2 + 8\tilde{g}^2 m_W^2 / g^2 m_{Z'}^2)^2 + 1}{1 + (1 + 4s_W^2)^2}$$

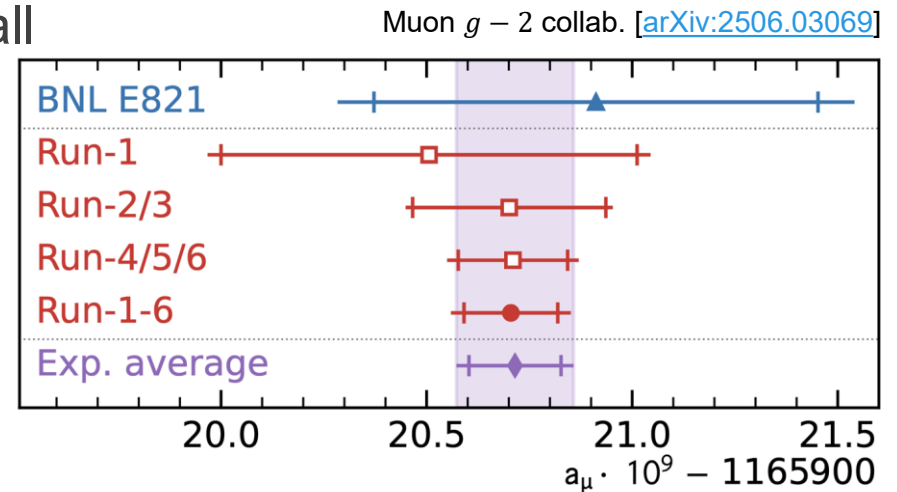
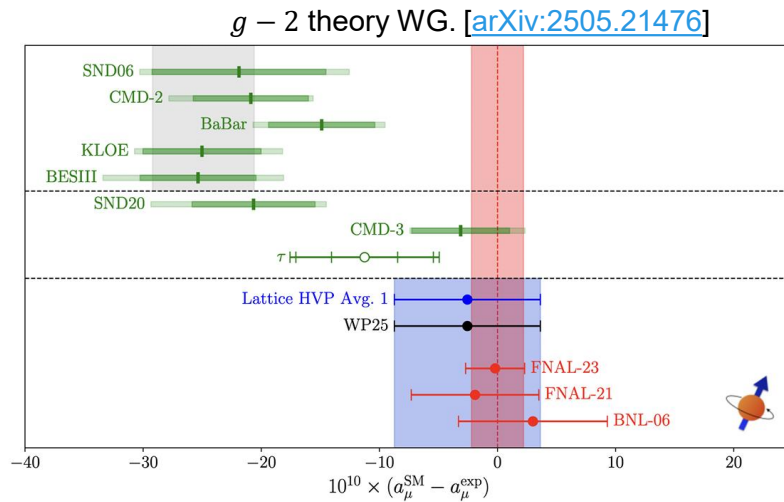


2025: The end of an anomaly

The Muon $g - 2$ collaboration has completed all their runs.

Much more precise, and consistent with 2021!

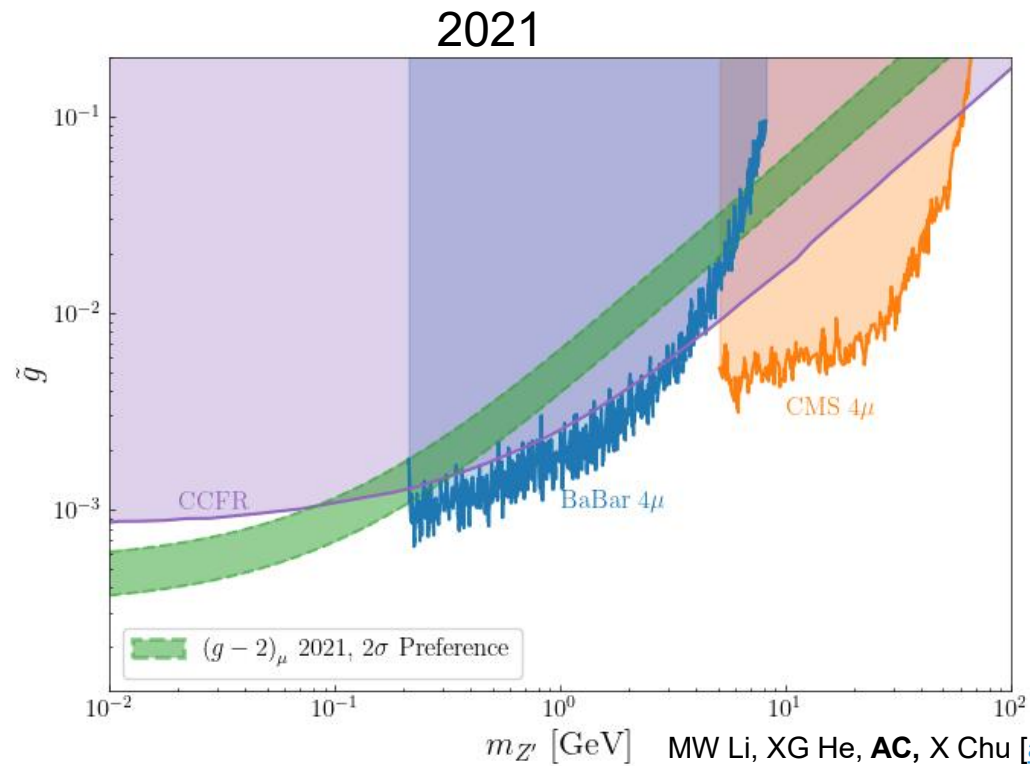
But the SM theory prediction has changed.



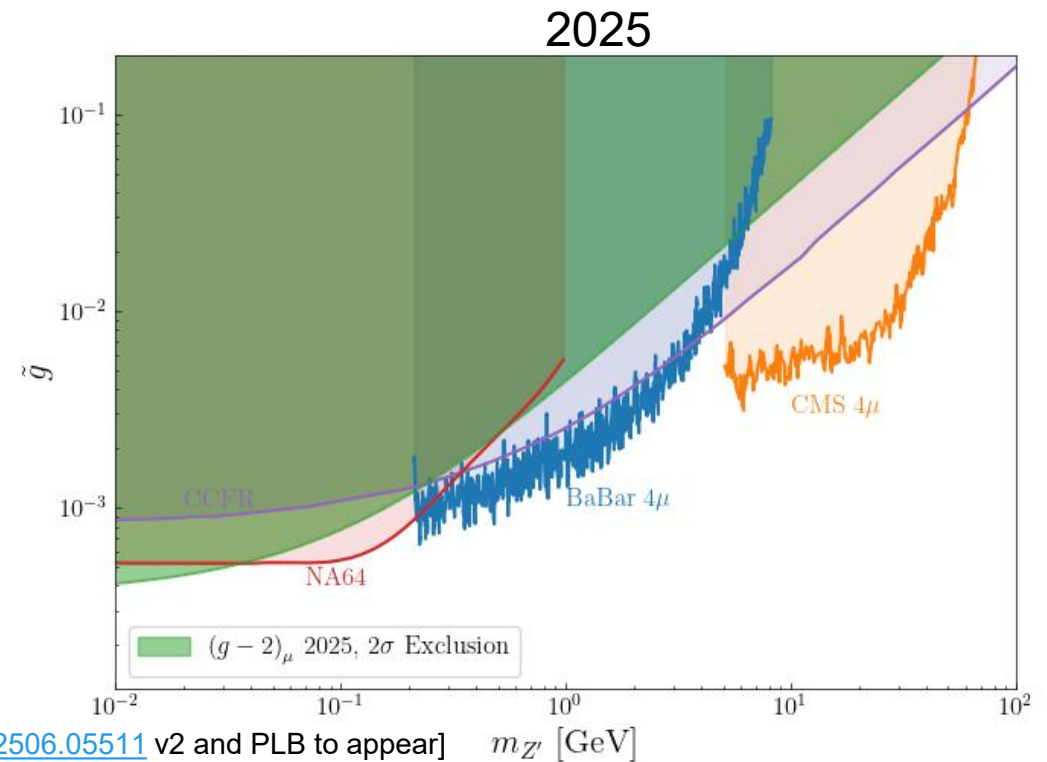
Green: Is dispersive methods for calculating the hadronic vacuum polarization.

CMD-3 measures $e^+e^- \rightarrow \pi^+\pi^-$ and shows large discrepancy.

From tension to exclusion



Green preferred!

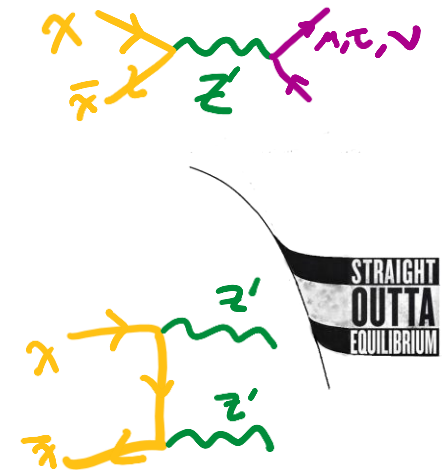
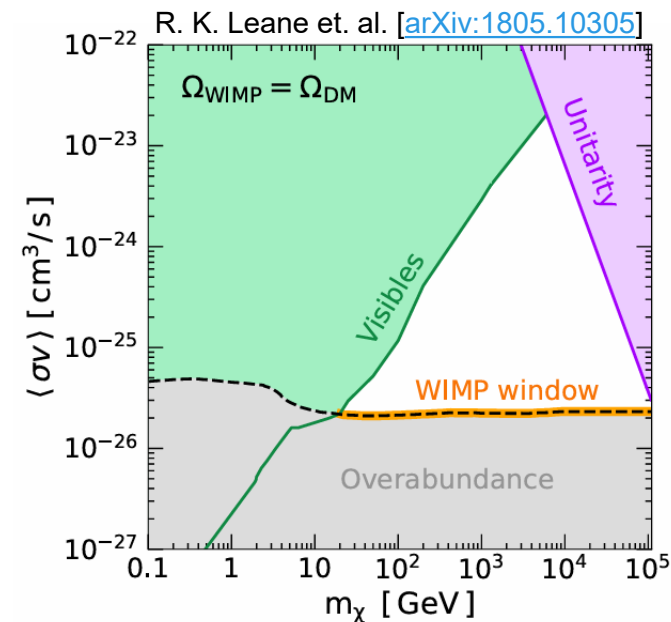
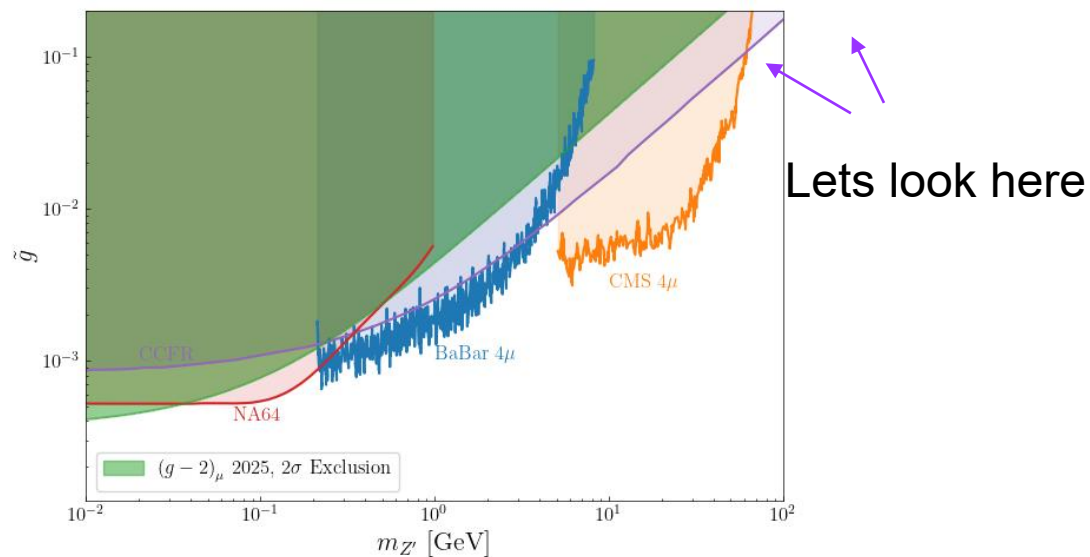


Green disfavored!

Time to give up on $U(1)_{L_\mu-L_\tau}$?

No!

1. As one of the simplest consistent extension to the SM, it should be thoroughly tested.
2. Thermal dark matter is still not ruled out!



$U(1)_{L_\mu - L_\tau}$ at a future muon collider

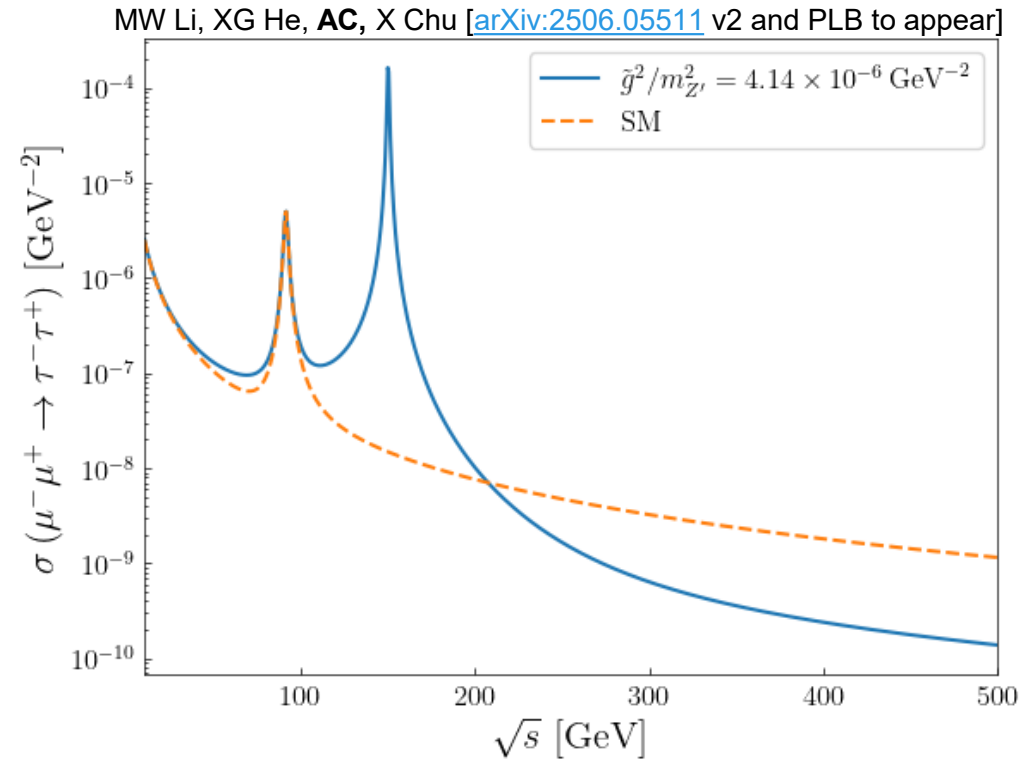
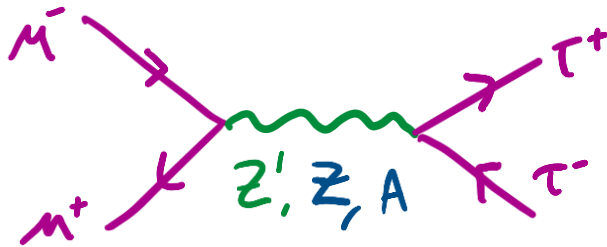
Taking the MNT limit as a benchmark,

CHARM-II: $\sigma_{\text{exp}}/\sigma_{\text{SM}} = 1.58 \pm 0.64$

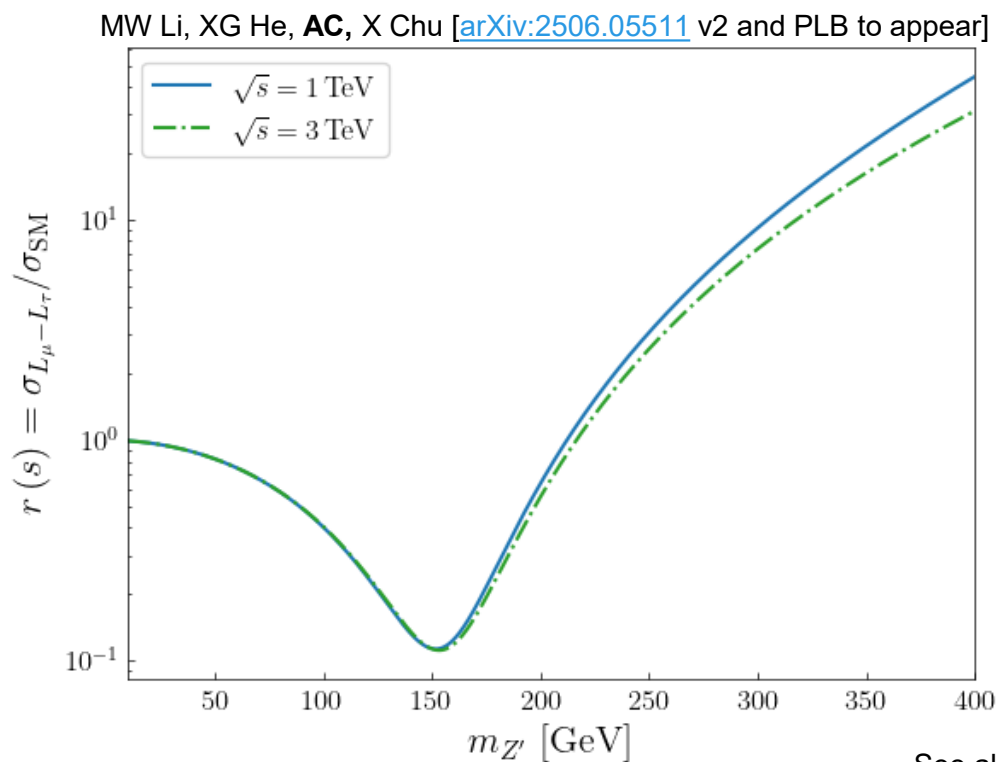
CCFR: $\sigma_{\text{exp}}/\sigma_{\text{SM}} = 0.82 \pm 0.28$

NuTeV: $\sigma_{\text{exp}}/\sigma_{\text{SM}} = 0.72^{+1.73}_{-0.72}$

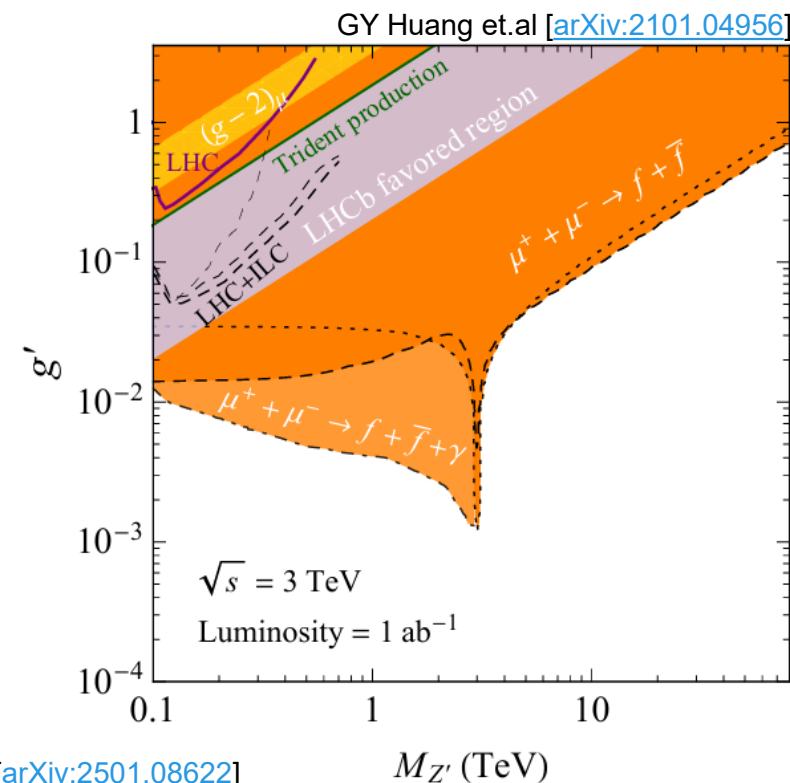
We would see a large deviation from SM in future muon collider



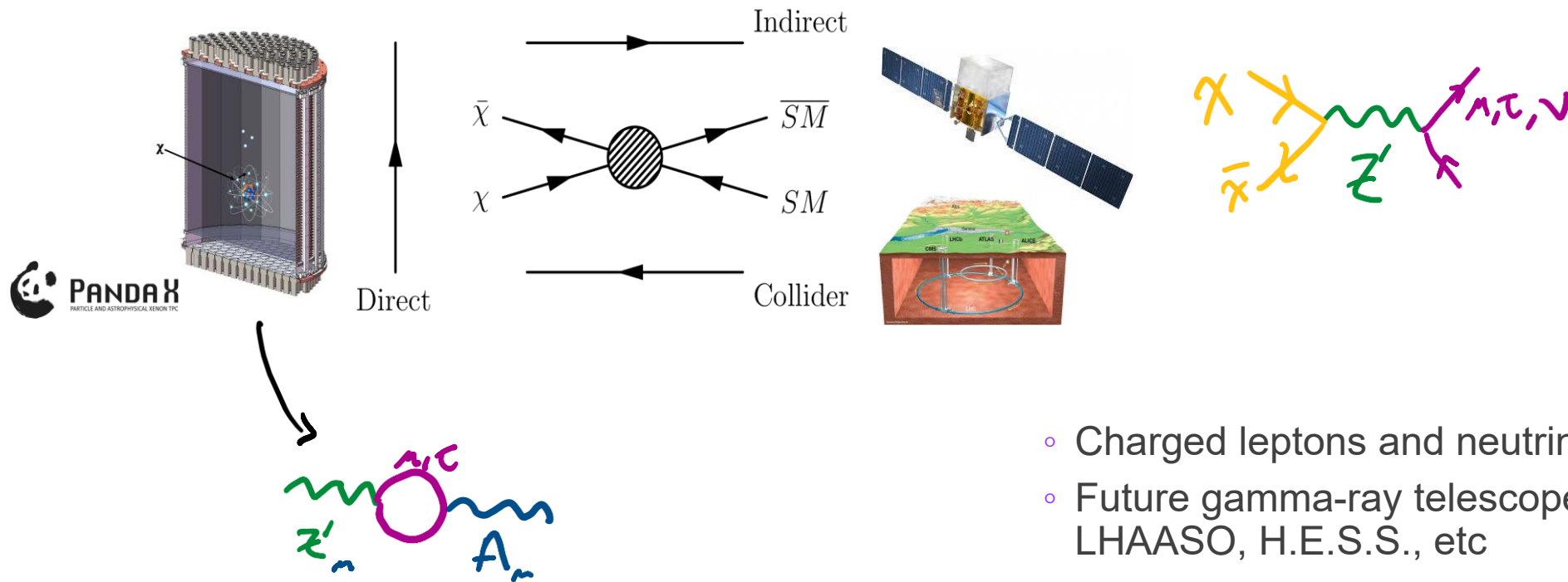
$U(1)_{L_\mu-L_\tau}$ at a future muon collider



See also ZW Wang et.al [[arXiv:2501.08622](https://arxiv.org/abs/2501.08622)]

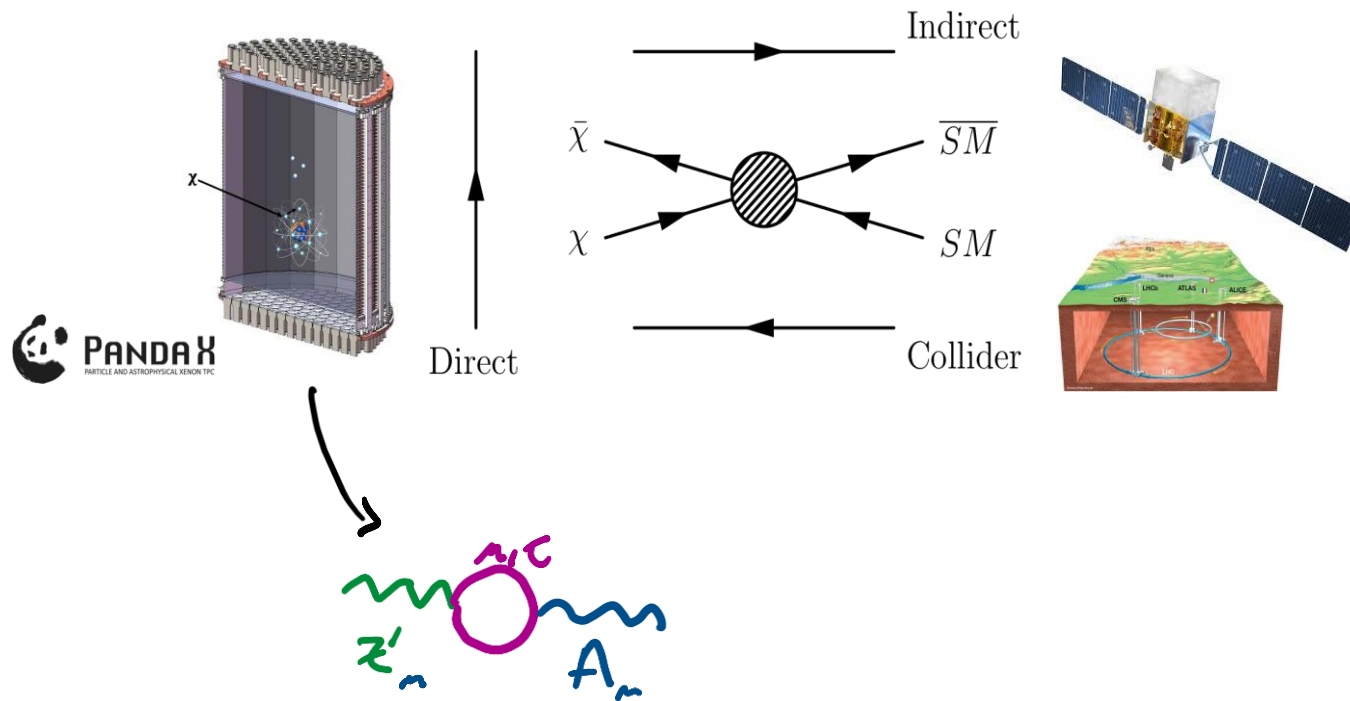


Dark matter and $U(1)_{L_\mu-L_\tau}$

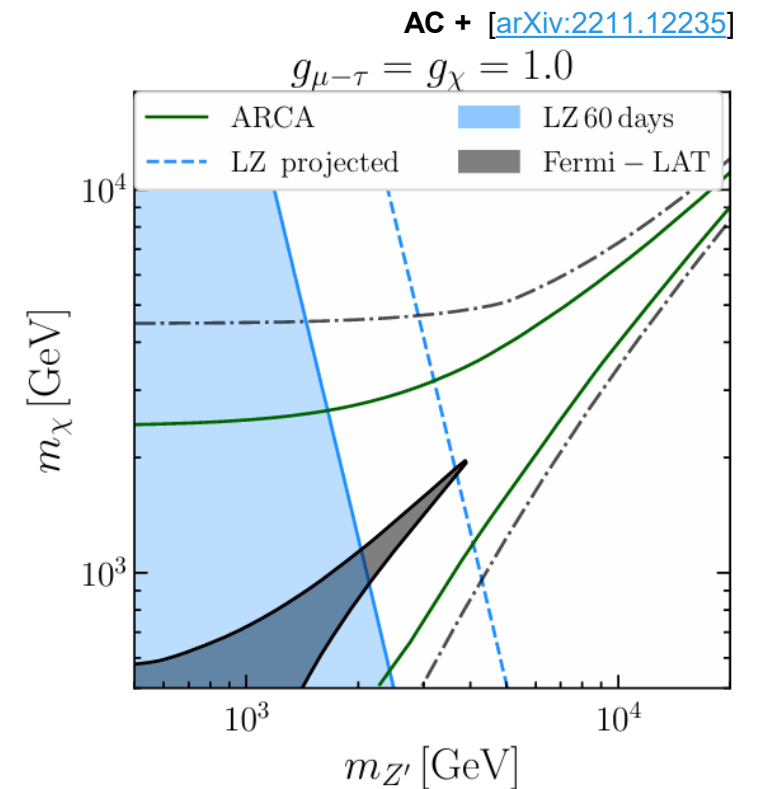


Interaction with nucleons come from kinetic mixing

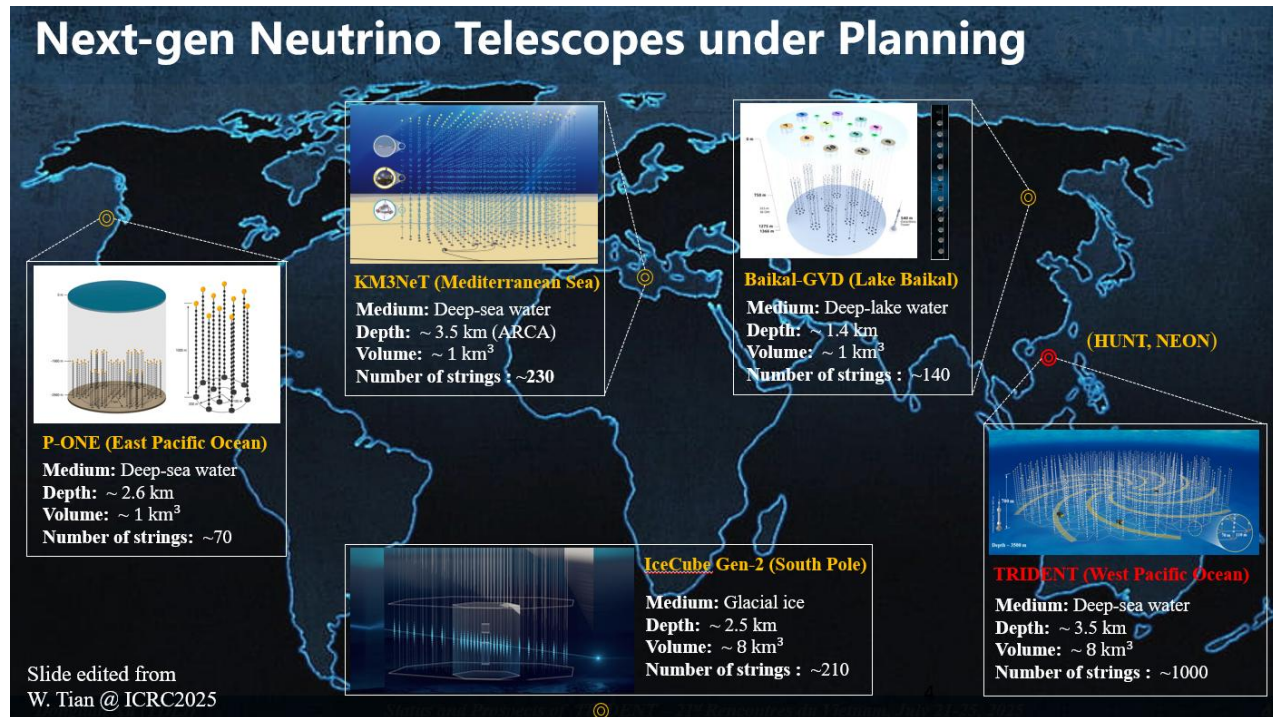
Dark matter and $U(1)_{L_\mu - L_\tau}$



Interaction with nucleons come from kinetic mixing



Neutrino telescopes



At TDLI! Currently at phase-I (build and deploy 10 strings by 2026).

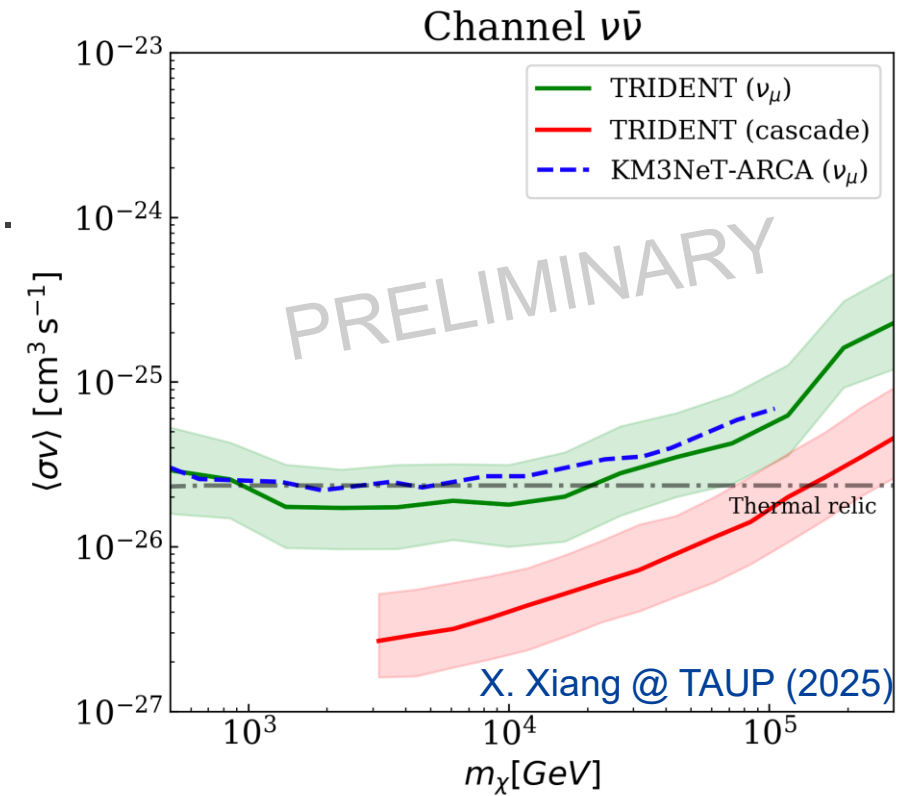
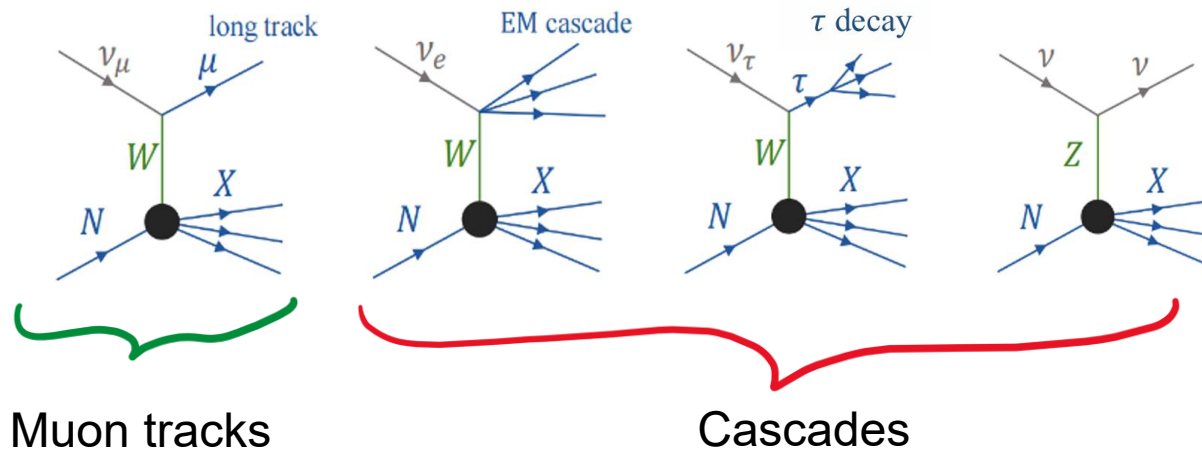
Full conceptual design = 1200 strings.

Trident prospects

TRIDENT will be roughly 10 times as big as KM3NeT.

Visibility of galactic center a little less.

Sea-water detectors are better at observing cascades.



Conclusion

Tension of $(g - 2)_\mu$ appears to be gone thanks to updates in SM theory calculations.

Previous emphasis on light Z' no longer required, encourages exploration of high-scale extensions.

We see this as a strong motivation for a muon collider.

One of the last allowed simple thermal DM models is the $U(1)_{L_\mu - L_\tau}$ and will be probed by future neutrino telescopes such as TRIDENT.