

# STRUCTURE OF A LOOP COMPUTATION

courtesy of Simone Devoto

## Process definition

*Feynman Amplitudes*

*Computation of the interference terms*

*Reduction to a set of Master Integrals*

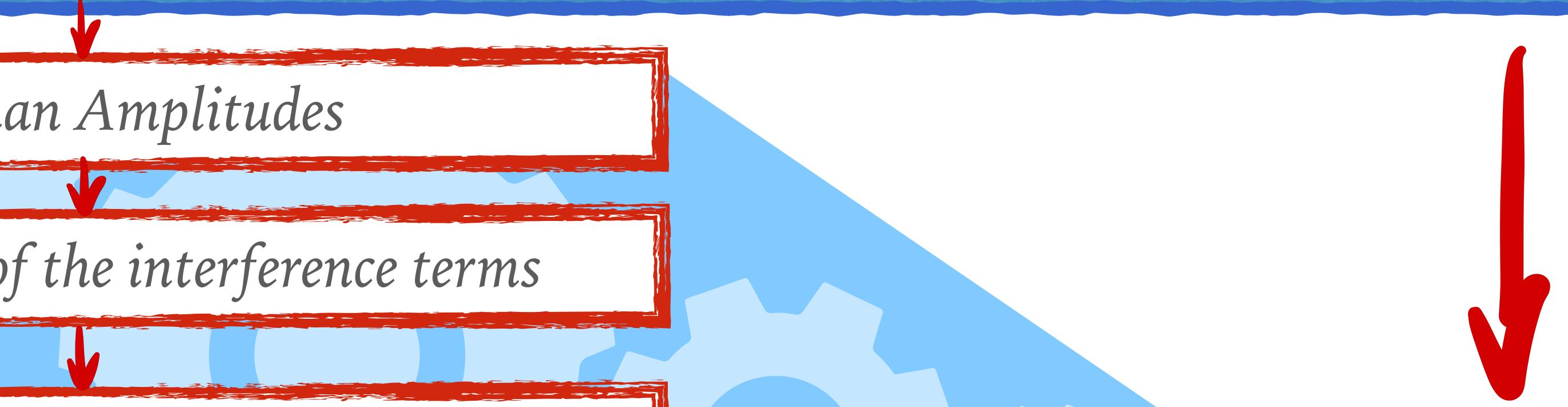
*Evaluation of the Master Integrals*

*Subtraction of the UV poles (renormalisation)*

*Subtraction of the IR poles*

*Numerical evaluation in phase-space points*

*Numerical grid*



# Solving Feynman integrals

$$I[\alpha_0, \alpha_1, \dots, \alpha_l] = \int \frac{d^n k_1}{(2\pi)^n} \int \frac{d^n k_2}{(2\pi)^n} \frac{1}{[k_1^2 - m_0^2]^{\alpha_0} [(k_1 + p_1)^2 - m_1^2]^{\alpha_1} \dots [(k_1 + k_2 + p_j)^2 - m_j^2]^{\alpha_j} \dots [(k_2 + p_l)^2 - m_l^2]^{\alpha_l}}$$

where  $I$  is the name of the “integral family”  $\rightarrow$  a specific value of momentum and mass of the  $k$ -th denominator is understood

The solution of a Feynman integral can be achieved by different techniques:

- direct numerical integration (naive, sector decomposition,...)
- integration of auxiliary parameters (Feynman, Schwinger,...)
- differential equations

A Feynman integral can be considered as an unknown function  $I = I(s, t, u, \dots, m_0, m_1, \dots, \varepsilon)$

which satisfies a (system of) linear first-order differential equations, with respect to kinematical invariants and masses, e.g.  $\frac{\partial \vec{I}}{\partial s} = \mathbf{A} \cdot \vec{I}$

How can we write the matrix  $\mathbf{A}$  ? How do we solve the system ?

# Differential equations and IBPs

M.Argeri and P.Mastrolia, “[Feynman Diagrams and Differential Equations](#)”  
*Int.J.Mod.Phys.A* 22 (2007) 4375-4436 • e-Print: [0707.4037](#)

- Not all the Feynman integrals in one amplitude are independent  
 → exploit Integration-by-parts (IBP) and Lorentz identities to reduce to a basis of independent Master Integrals

$$\int \frac{d^n k_1}{(2\pi)^n} \int \frac{d^n k_2}{(2\pi)^n} \frac{\partial}{\partial k_1^\mu} \frac{(k_1^\mu, k_2^\mu, p_r^\mu)}{[k_1^2 - m_0^2]^{\alpha_0} [(k_1 + p_1)^2 - m_1^2]^{\alpha_1} \dots [(k_1 + k_2 + p_j)^2 - m_j^2]^{\alpha_j} \dots [(k_2 + p_l)^2 - m_l^2]^{\alpha_l}} = 0$$

$$\int \frac{d^n k_1}{(2\pi)^n} \int \frac{d^n k_2}{(2\pi)^n} \frac{\partial}{\partial k_2^\mu} \frac{(k_1^\mu, k_2^\mu, p_r^\mu)}{[k_1^2 - m_0^2]^{\alpha_0} [(k_1 + p_1)^2 - m_1^2]^{\alpha_1} \dots [(k_1 + k_2 + p_j)^2 - m_j^2]^{\alpha_j} \dots [(k_2 + p_l)^2 - m_l^2]^{\alpha_l}} = 0$$

- The independent Master Integrals (MIs) satisfy a system of first-order linear differential equations with respect to each of the kinematical invariants / internal masses

When considering the complete set of MIs, the system can be cast in homogeneous form:  $d\vec{I}(\vec{s}; \varepsilon) = \mathbf{A}(\vec{s}; \varepsilon) \cdot \vec{I}(\vec{s}; \varepsilon)$

$$\frac{d}{dk^2} \text{ (loop diagram)} + \frac{1}{2} \left[ \frac{1}{k^2} - \frac{(D-3)}{(k^2 + 4m^2)} \right] \text{ (loop diagram)} = -\frac{(D-2)}{4m^2} \left[ \frac{1}{k^2} - \frac{1}{(k^2 + 4m^2)} \right] \text{ (loop diagram)}$$

- Henn’s conjecture (2013): if a change of basis exists which leads to  $d\vec{J}(\vec{s}; \varepsilon) = \varepsilon \tilde{\mathbf{A}}(\vec{s}) \cdot \vec{J}(\vec{s}; \varepsilon)$  then the solution is expressed in terms of iterated integrals (Chen integral representation) depending only on the results at previous orders in the  $\varepsilon$  expansion

<https://github.com/TommasoArmadillo/SeaSyde>

T.Armadillo, R.Bonciani, S.Devoto, N.Rana, A.Vicini,  
"Evaluation of Feynman integrals with arbitrary complex masses via series expansions",  
*Comput.Phys.Commun.* 282 (2023) 108545 • e-Print: [2205.03345](https://arxiv.org/abs/2205.03345)

Tommaso Armadillo,  
"Evaluating Feynman Integrals through differential equations and series expansions"  
[2502.14742](https://arxiv.org/abs/2502.14742)

# Evaluation of the Master Integrals by series expansions

T.Armadillo, R.Bonciani, S.Devoto, N.Rana, AV, 2205.03345

## A Simple Example

$$\begin{cases} f'(x) + \frac{1}{x^2 - 4x + 5} f(x) = \frac{1}{x+2} \\ f(0) = 1 \end{cases}$$

$$f_{hom}(x) = x^r \sum_{k=0}^{\infty} c_k x^k$$

$$f'_{hom}(x) = \sum_{k=0}^{\infty} (k+r) c_k x^{(k+r-1)}$$

$$\begin{cases} r c_0 = 0 \\ \frac{1}{5} c_0 + c_1(r+1) = 0 \\ \frac{4}{25} c_0 + \frac{1}{5} c_1 + c_2(2+r) = 0 \\ \dots \end{cases}$$

$$f_{hom}(x) = 5 - x - \frac{3}{10}x^2 + \frac{11}{150}x^3 + \dots$$

$$\begin{aligned} f_{part}(x) &= f_{hom}(x) \int_0^x dx' \frac{1}{(x'+2)} f_{hom}^{-1}(x') \\ &= \frac{1}{2}x - \frac{7}{40}x^2 + \frac{2}{75}x^3 + \dots \end{aligned}$$

$$f(x) = f_{part}(x) + C f_{hom}(x)$$

$$f(0) = 1 \rightarrow C = \frac{1}{5}$$

*Expanded around  $x' = 0$*

# Evaluation of the Master Integrals by series expansions

T.Armadillo, R.Bonciani, S.Devoto, N.Rana, AV, 2205.03345

- **Taylor expansion:** avoids the singularities;
- **Logarithmic expansion:** uses the singularities as **expansion points**.
- Logarithmic expansion has larger convergence radius but requires longer evaluation time. **We use Taylor expansion as default.**

