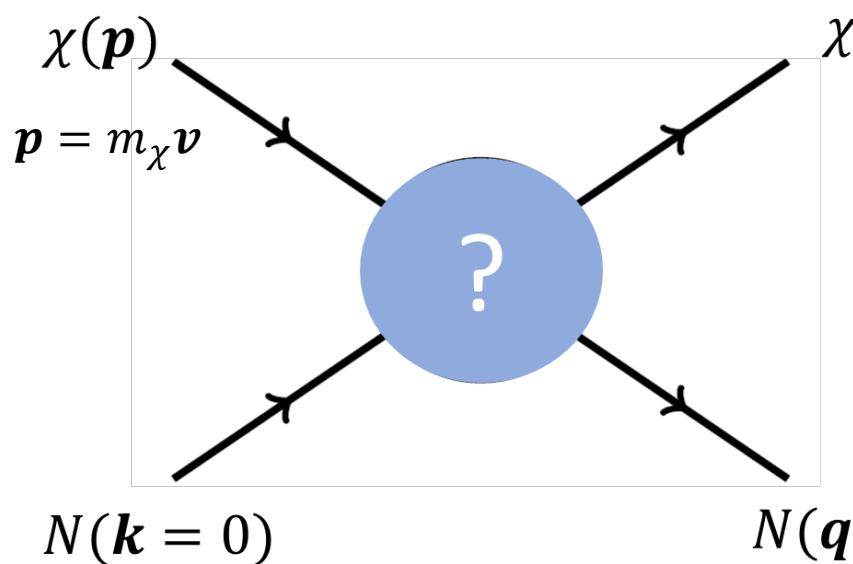


# kinematics

Detector @  $T \approx 300K \rightarrow K \approx K_B T \sim 0.025 \text{ eV} \rightarrow$  In the **lab frame** we can safely neglect the nucleus thermal motion



$$E_i = E_f \rightarrow \frac{p^2}{2m_\chi} = \frac{q^2}{2m_N} + \frac{(\mathbf{p} - \mathbf{q})^2}{2m_\chi}$$

$$\rightarrow \mathbf{p}\mathbf{q} = \frac{q^2 m_\chi}{2\mu_{\chi N}} \quad \text{Where } \mu_{\chi N} = \frac{m_\chi m_N}{m_\chi + m_N}$$

$$\text{Now, } \mathbf{p}\mathbf{q} = pq \cos\theta_{lab} \rightarrow q = \frac{2\mu_{\chi N}}{m_\chi} p \cos\theta_{lab}$$

$$E_R = \frac{q^2}{2m_N} = \frac{2\mu_{\chi N}^2}{m_N} v^2 \cos^2\theta_{lab}$$

$$\text{Max momentum transfer } (\theta_{lab} = 0) \rightarrow q_{max} = \frac{2p\mu_{\chi N}}{m_\chi} \quad \Rightarrow \quad E_R^{max} = \frac{q_{max}^2}{2m_N} = \frac{2\mu_{\chi N}^2}{m_N} v^2$$

For a given Energy threshold  $E_R$  there is a minimum velocity of the WIMP to be visible in the detector

$$v_{min} = \sqrt{\frac{E_R m_N}{2\mu_{\chi N}^2}}$$

# Expected rate in the detector

---

The expected WIMP rate in the detector is given by  $R \approx N_T \times n_\chi \times \langle v \rangle \times \sigma$

where:

$N_T = M_{det} / m_N$  is the number of target nuclei (NOTE: here  $M_{det}$  is in the same units as  $m_N$ )

$n_\chi = \rho_\chi / m_\chi$  is the number density of WIMPs and  $m_\chi$  is the WIMP mass

$\rho_\chi$  is the DM local density

$\langle v \rangle$  is the mean relative velocity WIMP / detector

$\sigma$  is the cross section, that in general depends on the transferred moment

If  $f(\mathbf{v})$  is the WIMP velocity distribution in the **detector's frame**, so  $\langle v \rangle = \int_0^\infty v f(\mathbf{v}) d^3 \mathbf{v}$  so:

**differential rate per unit of recoil energy:**

$$\frac{dR}{dE_R} = \frac{M_{det} \rho_\chi}{m_N m_\chi} \int_{v_{min}}^\infty \frac{d\sigma}{dE_R} v f(\mathbf{v}) d^3 \mathbf{v}$$

where  $v_{min} = \sqrt{\frac{E_R m_N}{2 \mu_{\chi N}^2}}$

# The Master formula

$$\frac{dR}{dE_R} = \frac{M_{det}\rho_\chi}{m_N m_\chi} \int_{v_{min}}^{v_{esc}} \frac{d\sigma}{dE_R}(q, v) v f(v) d^3v$$

$$\frac{d\sigma}{dE_R}(v, q) = \frac{d\sigma}{dE_R}(v, 0) F^2(q)$$

nuclear form factor

$$\frac{d\sigma}{dE_R}(v, 0) = \frac{\sigma^0}{E_R^{max}} = \frac{m_N}{2\mu_{\chi N}^2 v^2} \sigma^0$$

DM-nucleon  
SI cross-section

$$(\text{consider only SI}) \quad \sigma^0 = \sigma_{SI}^0 = \frac{A^2 \mu_{\chi N}^2}{\mu_{\chi n}^2} \sigma_{SI}$$

DM-nucleon  
SI cross-section



$$\frac{dR}{dE_R} = \frac{M_{det}\rho_\chi}{2m_\chi \mu_{\chi n}^2} A^2 \sigma_{SI} F_{SI}^2(q) \int_{v_{min}}^{v_{max}} \frac{f(v, t)}{v} d^3v$$

WIMP mass

$\eta(v_{min}, t) = \int_{v_{min}}^{v_{esc}} \frac{f(v, t)}{v} d^3v$  is the **mean inverse speed function** (note that it depends on time due to

the Earth rotation around the Sun) and  $v_{min} = \sqrt{\frac{E_R m_N}{2\mu_{\chi N}^2}}$ ,  $\mu_{\chi n} = \frac{m_\chi m_n}{m_\chi + m_n}$

# Velocity in the lab reference system: time dependency

$$\mathbf{v}_{gal} = \mathbf{v} + \mathbf{v}_\odot + \mathbf{v}_\oplus(t)$$

Where

$\mathbf{v}_{gal}$ : **WIMP velocity in the GALACTIC reference frame**

$\mathbf{v}$ : **WIMP velocity in the laboratory reference frame**

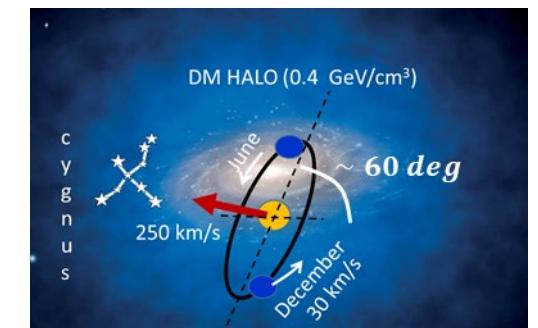
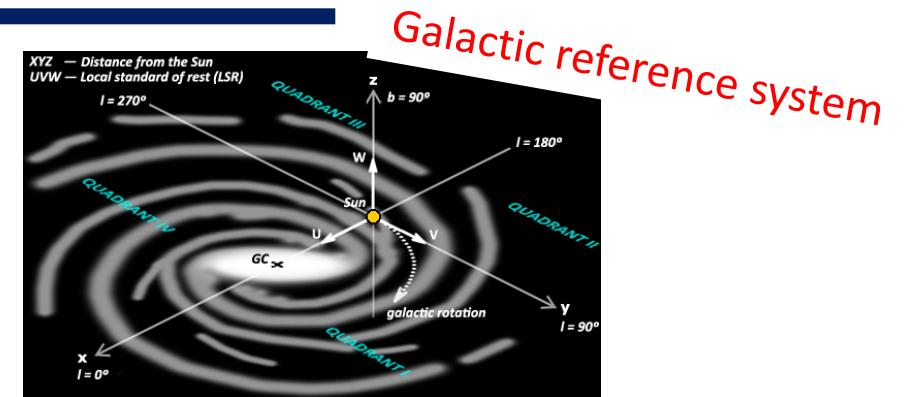
$\mathbf{v}_E$ : **Earth velocity with respect to the Galaxy frame**,  $\mathbf{v}_E = \mathbf{v}_\odot + \mathbf{v}_\oplus(t)$

$\mathbf{v}_\odot = \mathbf{v}_0 + \mathbf{v}_{pec}$ : velocity of the Sun with respect to the galaxy

$\mathbf{v}_0$ : **local standard of rest velocity ( $\mathbf{v}_0$ )**

$\mathbf{v}_{pec}$ : solar peculiar velocity

$\mathbf{v}_\oplus(t)$ : **velocity of the Earth with respect to the Sun**



**Table 1** Suggested Standard Halo Model parameters. Vectors are given as  $(v_r, v_\phi, v_\theta)$  with  $r$  pointing radially inward and  $\phi$  in the direction of the Milky Way's rotation. Analyses insensitive to annular modulation can approximate  $\mathbf{v}_\oplus(t)$  with Eq. 12

Parameter	Description	Value	References
$\rho_\chi$	Local dark matter density	$0.3 \text{ GeV}/c^2/\text{cm}^3$	[9]
$v_{esc}$	Galactic escape speed	544 km/s	[45]
$\langle  \mathbf{v}_\oplus  \rangle$	Average galactocentric Earth speed	29.8 km/s	[41]
$\mathbf{v}_\odot$	Solar peculiar velocity	$(11.1, 12.2, 7.3) \text{ km/s}$	[46]
$\mathbf{v}_0$	Local standard of rest velocity	$(0, 238, 0) \text{ km/s}$	[47,48]

$$\mathbf{v}_\oplus(t) = \langle |\mathbf{v}_\oplus| \rangle \times \begin{pmatrix} 0.9941 \cos(\omega \Delta t) - 0.0504 \sin(\omega \Delta t) \\ 0.1088 \cos(\omega \Delta t) + 0.4946 \sin(\omega \Delta t) \\ 0.0042 \cos(\omega \Delta t) - 0.8677 \sin(\omega \Delta t) \end{pmatrix}$$

where  $\omega = 0.0172 \text{ d}^{-1}$  and  $t$  is the number of days since March 22, 2018

D. Baxter et al. "Recommended conventions for reporting results from direct dark matter searches" Eur. Phys. J. C (2021) 81:907 [2105.00599]

# Standard halo model (SHM)

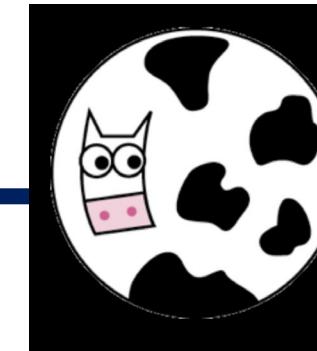
The standard halo model is an isotropic, isothermal sphere **in the galactic frame**, with density profile  $\rho(r) \propto r^{-2}$ . In this case the solution to the collisionless Boltzmann equation is a Maxwellian velocity distribution ( $\nu_{gal}$ :DM  $\nu$  in the galactic reference system)

$$f_{gal}(\nu_{gal})d^3\nu_{gal} = \frac{1}{\nu_0^3 \pi^{3/2}} e^{-\nu_{gal}^2/\nu_0^2} d^3\nu_{gal}$$

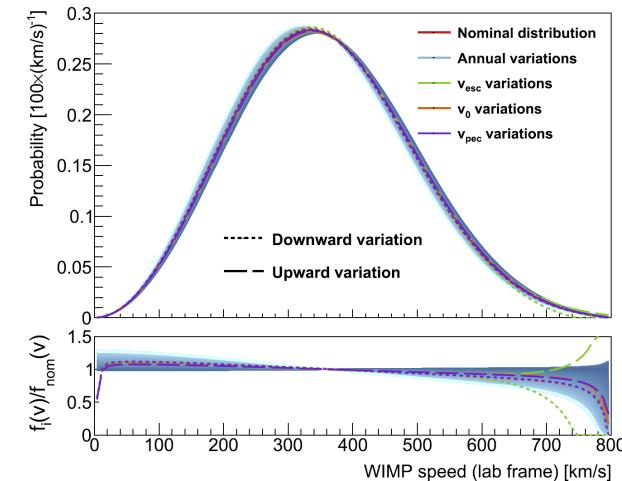
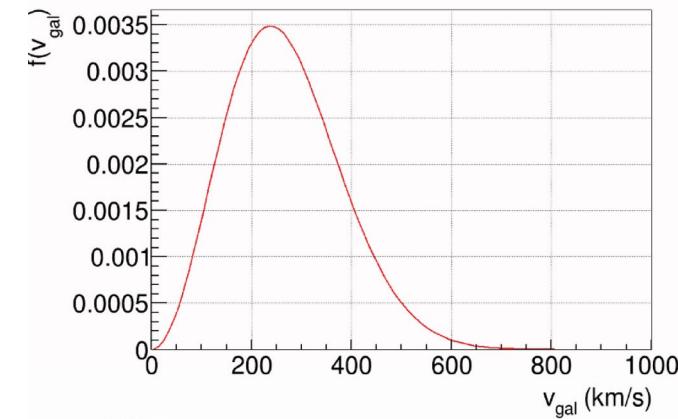


Move to lab. frame

$$f(\nu, t) = f_{gal}(\nu + \nu_{\odot} + \nu_{\oplus}(t))$$



*The spherical cow of direct WIMP searches (Gelmini)*



# The mean inverse speed function

---

$$\eta(v_{min}, t) = \int_{v_{min}}^{\infty} \frac{f(v, t)}{v} d^3v$$

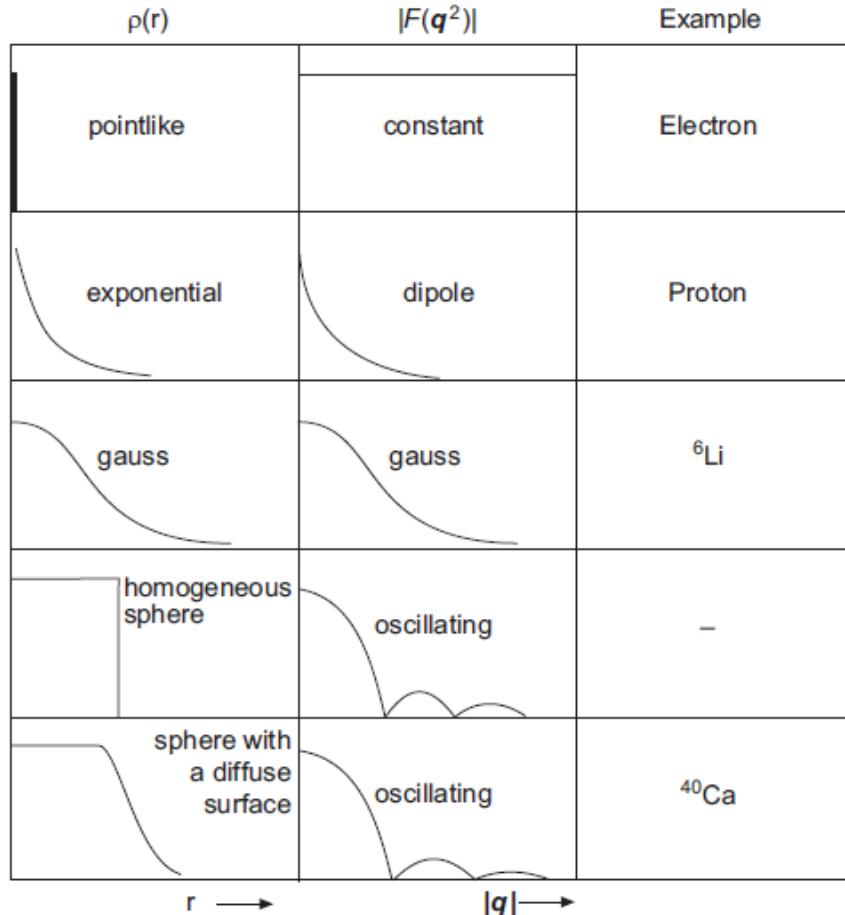
This is the integral to be solved for the DM halo model considered. For the SHM, this case the integral can be solved analytically. Defining:

$$x = \frac{v_{min}}{v_o}, y = \frac{v_{\odot} + v_{\oplus}(t)}{v_o}, z = \frac{v_{esc}}{v_o}$$

$$\eta(v_{min}, t) = \frac{1}{2yv_o} \frac{1}{N} \begin{cases} \text{erf}(x + y) - \text{erf}(x - y) - \frac{4}{\sqrt{\pi}} ye^{-z^2} & 0 \leq x \leq z - y \\ \text{erf}(z) - \text{erf}(x - y) - \frac{2}{\sqrt{\pi}} (z + y - x) e^{-z^2} & z - y < x \leq z + y \\ 0 & x > z + y \end{cases}$$

Here,  $\text{erf}()$  is the error function and  $N = \text{erf}(z) - \frac{2z}{\sqrt{\pi}} e^{-z^2}$  is a normalization factor

# Spin independent form factor



An analytic expression for the FF (and the most commonly used in DD calculations) is that from Helm:

$$F_{\text{SI}}^2(q) = \left( \frac{3j_1(qR_1)}{qR_1} \right)^2 e^{-q^2 s^2}$$

$$\text{where } j_1(x) = \frac{\sin x}{x^2} - \frac{\cos x}{x}$$

is a spherical Bessel function of the first kind,  **$R_1$  is an effective nuclear radius** and  **$s$  is the nuclear skin thickness**, These parameters that need to be fit separately for each nucleus, but good results are obtained with

$$R_1 = \sqrt{R^2 - 5s^2}, \quad R \approx 1.2A^{1/3} \text{ fm}, \quad s = 1 \text{ fm}$$

Other parametrizations can give more precise results for high  $q$ , see for example:

G. Duda, A. Kemper and P. Gondolo, "Model Independent Form Factors for Spin Independent Neutralino-Nucleon Scattering from Elastic Electron Scattering Data", JCAP 0704:012,2007 [arXiv:hep-ph/0608035]

# Functions in the python code

---

`vearth(t)`: Earth velocity

`eta(E, t, A, mW)`: mean inverse speed function

`FF(E, A)`: SI form factor

`rate(E, t, A, mW, sigmasI)`: differential rate  $\frac{dR}{dE} = \frac{M_{det} \rho_\chi}{2 m_\chi \mu_{\chi n}^2} A^2 \sigma_{SI} F_{SI}^2 \eta$

`totalRate(Ei, Ef, t, A, mW, sigmaSI)`:  $R = \int_{E_{min}}^{E_{max}} \frac{dR}{dE} dE$

`t` in days from March 22

`E` in keV

`mW` in GeV/c<sup>2</sup>

`sigmasI` in cm<sup>2</sup>

`A` is the mass number of the target

+ examples to plot Earth velocity, form factor, meas inverse speed function and differential rate

# Exercises

---

- Plot the differential rate vs recoil energy for several targets: Ar(40), Ge(72), Xe(132)(\*)
- Plot the differential rate vs recoil energy for several Wimp masses: 10, 100, 1000 GeV/c<sup>2</sup>
- Plot the differential rate vs recoil energy for 2<sup>nd</sup> June (timeMax) and 1<sup>st</sup> December (timeMin)
- For Xenon target and a WIMP of  $m_\chi = 70 \text{ GeV}$ ,  $\sigma_{SI} = 10^{-41} \text{ cm}^2$  and two different energy intervals ([0-60] keV and [10-60] keV):

- 1) Implement a function to plot the TOTAL rate in  $(E_i, E_f)$ : `ratevsTime (Ei, Ef, A, mW, sigmasI)`
- 2) Plot rate total vs time for 1 year
- 3) Compute the maximum rate ( $R_{max}$ ), the minimum rate ( $R_{min}$ ), the average rate ( $R_0$ ) and the day corresponding to the maximum ( $t_{max}$ )
- 4) In the same figure as before, plot the following approximation for the rate:

$$R_{approx} = R_0 + R_{mod} \cos(\omega(t - t_{max})) \quad \text{where } R_0 \text{ is the average rate and}$$
$$R_{mod} = \frac{1}{2}(R_{max} - R_{min}) \text{ and } \omega = 2\pi/365 \text{ d}^{-1}$$

- 4) Calculate the error (in %) of the approximation in  $t_{max}$  as:

$$err(t_{max}) = \frac{R(t_{max}) - R_{approx}(t_{max})}{R(t_{max})} \times 100$$

(\*) in the python code you will find an example, just modify it