

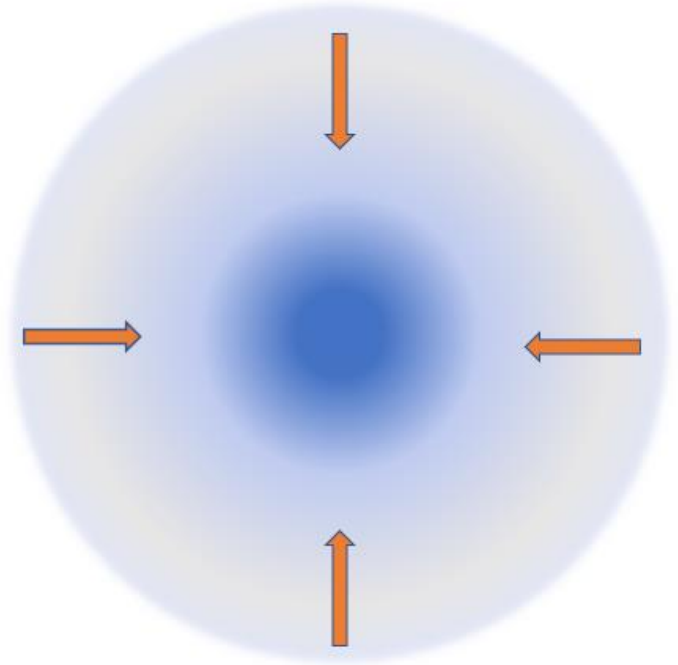
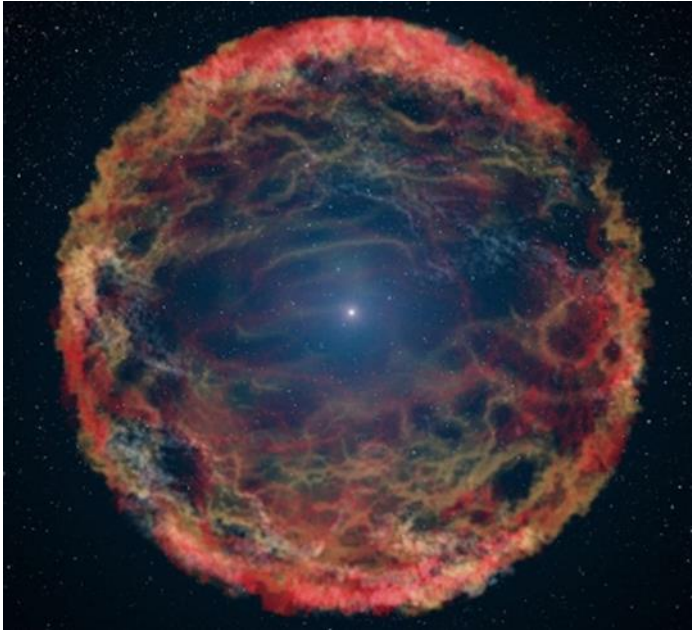


# Numerical simulations of fallback core-collapse supernovae: the lower mass gap and multi messenger signals

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# Core-collapse supernovae and fallback



Massive stars ( $> 8 M_{\odot}$ ) would end their lives as a core-collapse supernova (CCSN).

Many CCSN explosion calculations have confirmed the existence of fallback and studied its dynamics and effects.

# Fallback in core-collapse supernovae

Fallback might be related to a number of the observed phenomena:

- the peculiar supernovae
- the late-time neutrino emission
- the neutron star /black hole kick
- the long-duration gamma-ray bursts

The intensity of the fallback:

supernova explosion energy VS the binding energy of the star

Powerful fallback



Weak explosion energy

Faint electromagnetic signals

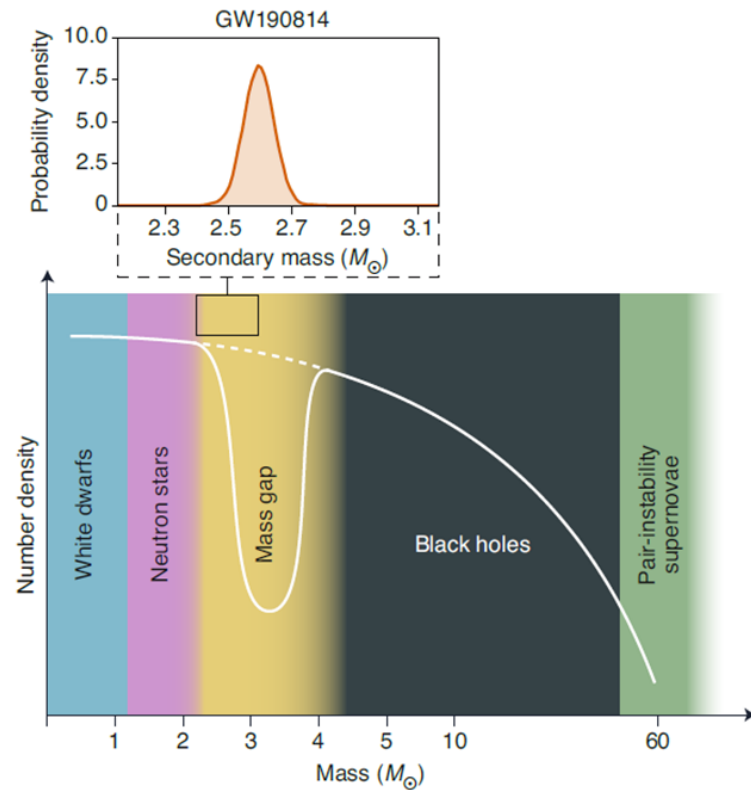
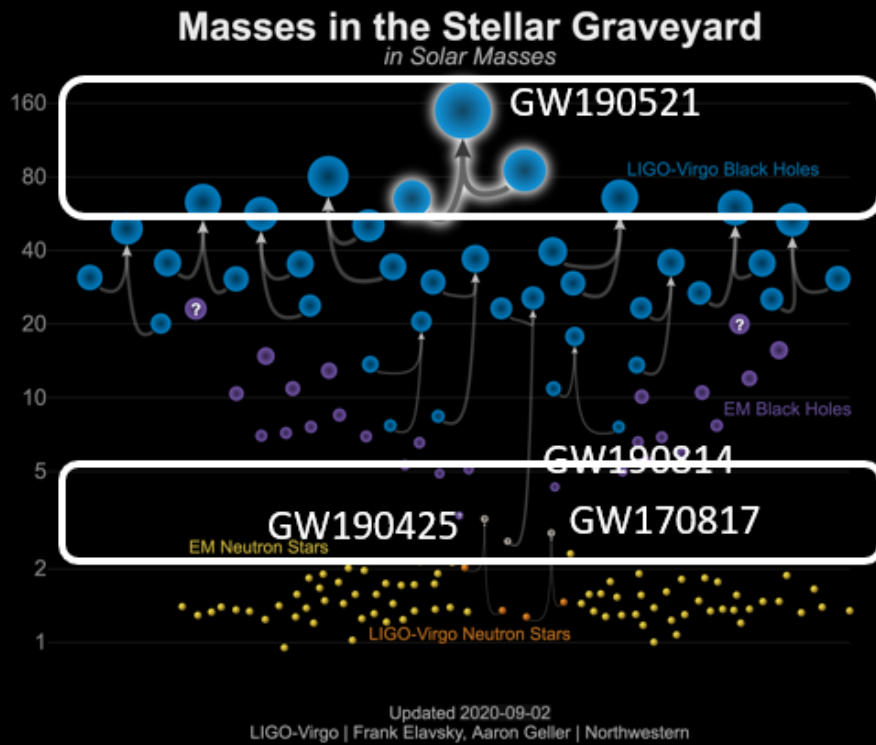


neutrinos and gravitational waves

# Two works in fallback supernova

- 1. The lower mass gap.
- 2. Neutrinos and gravitational waves signals from black hole neutrino-dominated accretion flows in fallback core-collapse supernovae.

# The lower mass gap

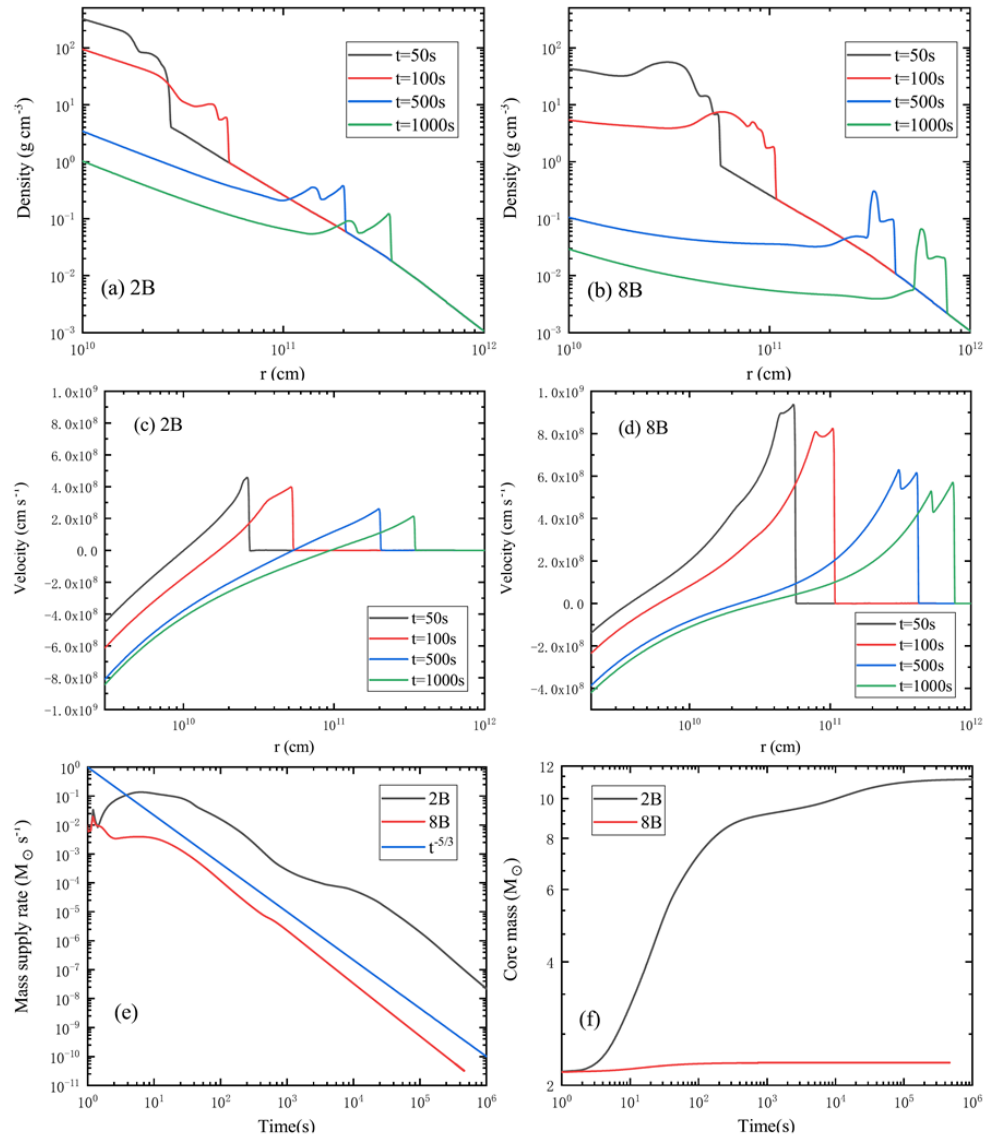


Safarzadeh (2020)

# CCSN simulation (piston approach)

- pre-SN models (KEPLER code)
- initial progenitor mass in range of 20-40  $M_{\odot}$
- metallicities  $Z/Z_{\odot} = 0, 0.01, 1$
- Athena++ code
- piston approach (one-dimensional CCSN simulations)
- initial explosion energies  $E=2, 4, 8B$  (1B=  $10^{51}$  erg )

# Results



**Figure 1.** Profiles of density and velocity at 50 s, 100 s, 500 s, and 1,000 s and time evolutions of mass supply rate and core mass for  $E = 2B$  and  $8B$  with progenitor mass  $30 M_{\odot}$  and  $Z/Z_{\odot} = 0.01$ .

# Results

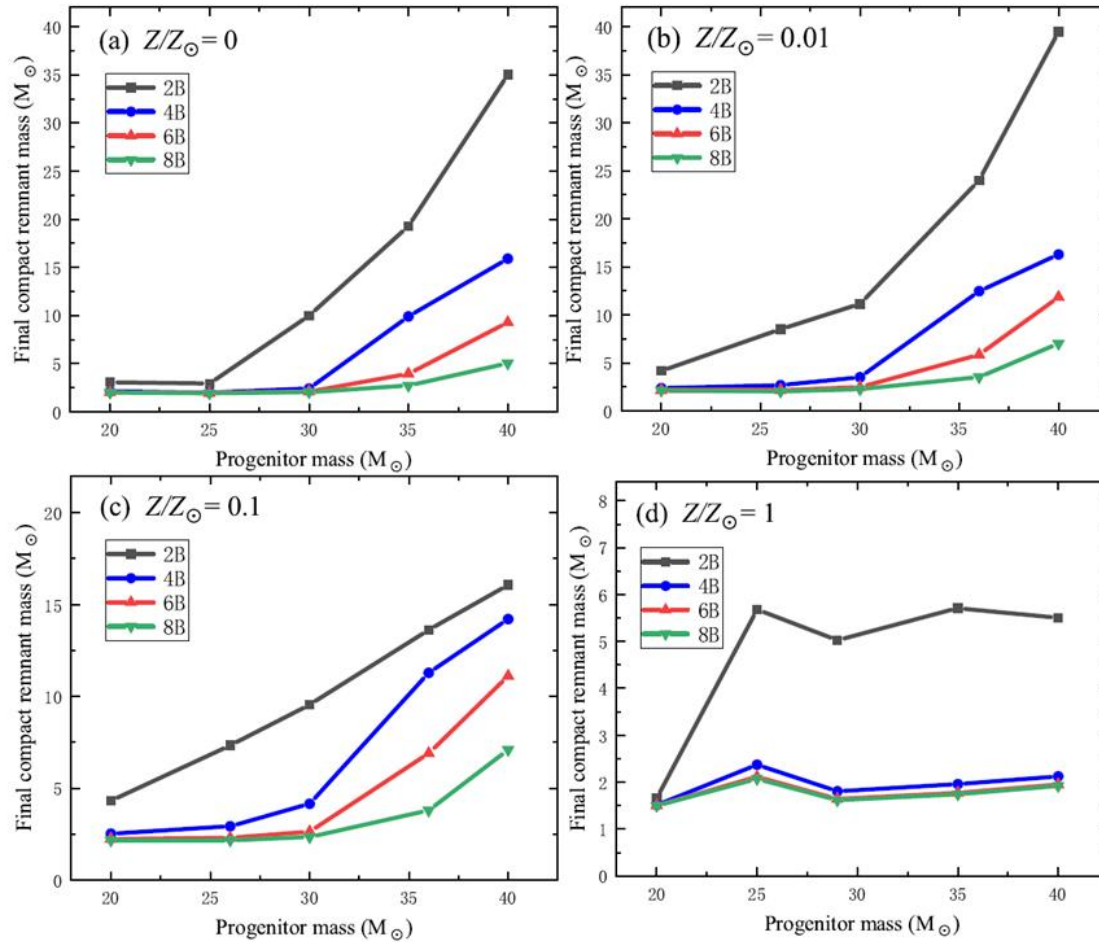
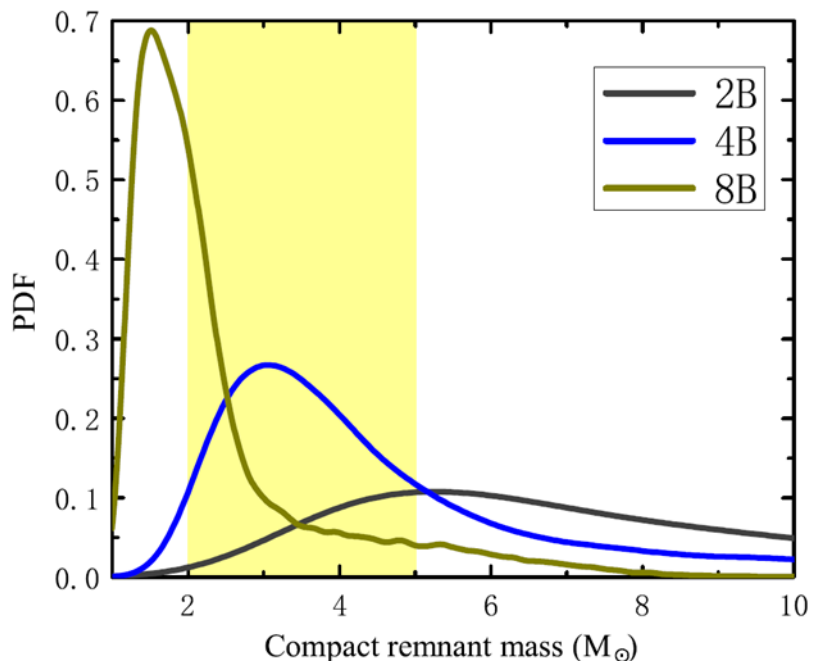


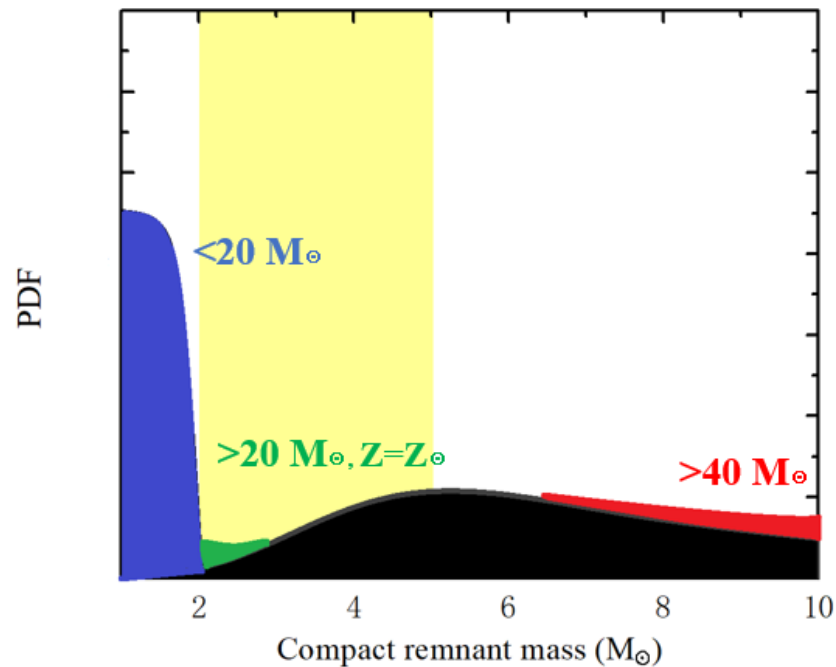
Figure 1. Final compact remnant masses in the center of the CCSNe with different explosion energies (2B, 4B, 6B, and 8B, where 1B =  $10^{51}$  erg) for the progenitor masses in the range of 20 – 40  $M_{\odot}$  and with a metallicity of  $Z/Z_{\odot} = 0, 0.01, 0.1,$  and 1.



# Applications on the lower mass gap



The lower explosion energy naturally causes more efficient fallback accretion for low-metallicity collapsars, and then the newborn BHs in the center of the CCSNe can escape from the gap, but neutron stars cannot easily collapse into BHs in the gap; nevertheless, the final remnants of the solar-metallicity progenitors stick to the gap.

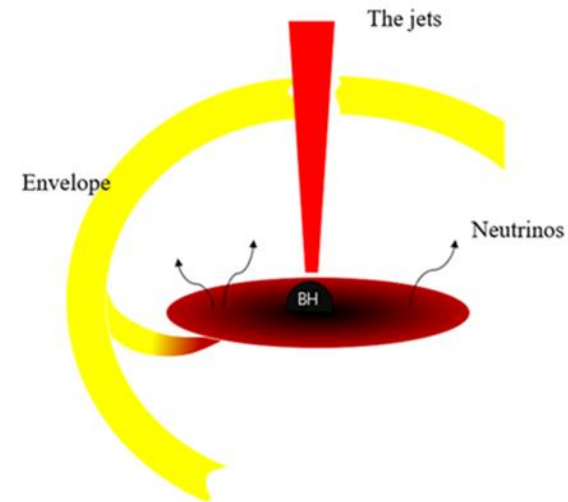
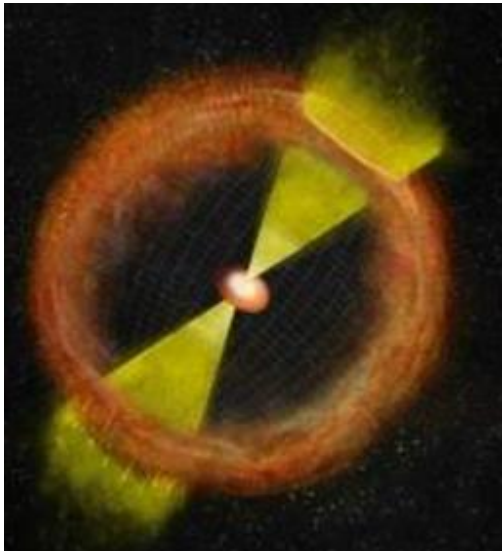


**The process of black hole accretion ?**

# Black hole (BH) hyperaccretion in CCSN

Infalling material with **enough angular momentum** would be slowed by the rotation and be piled into a disk.

If **the accretion rate is extremely high and neutrino cooling is dominant**, the hyperaccretion should be in the phase of the neutrino-dominated accretion flows (NDAFs)

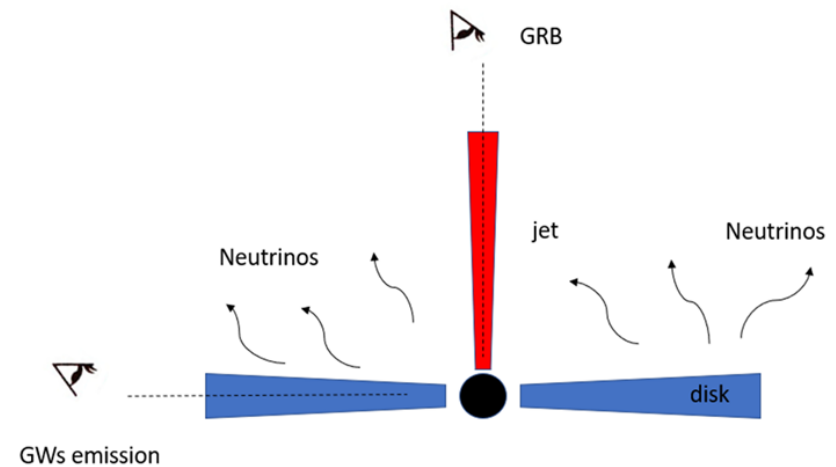
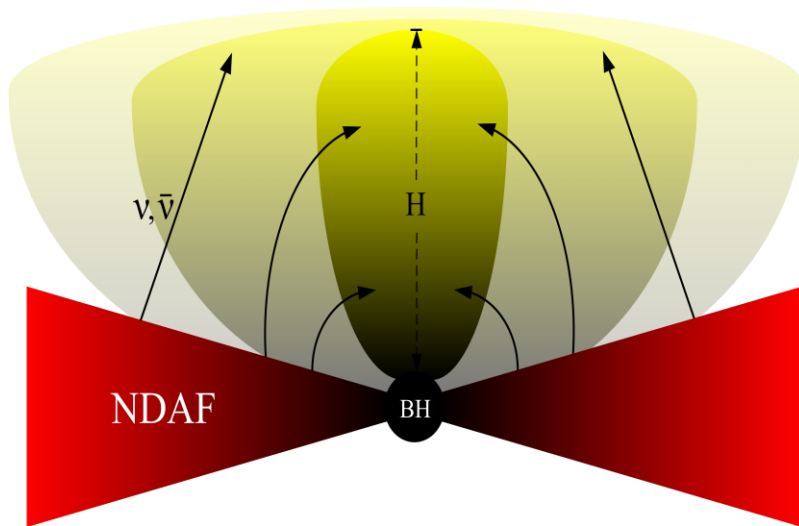


# Neutrino dominated accretion flows (NDAFs)

NDAFs around stellar-mass black holes are plausible candidates for the central engine of gamma-ray bursts.

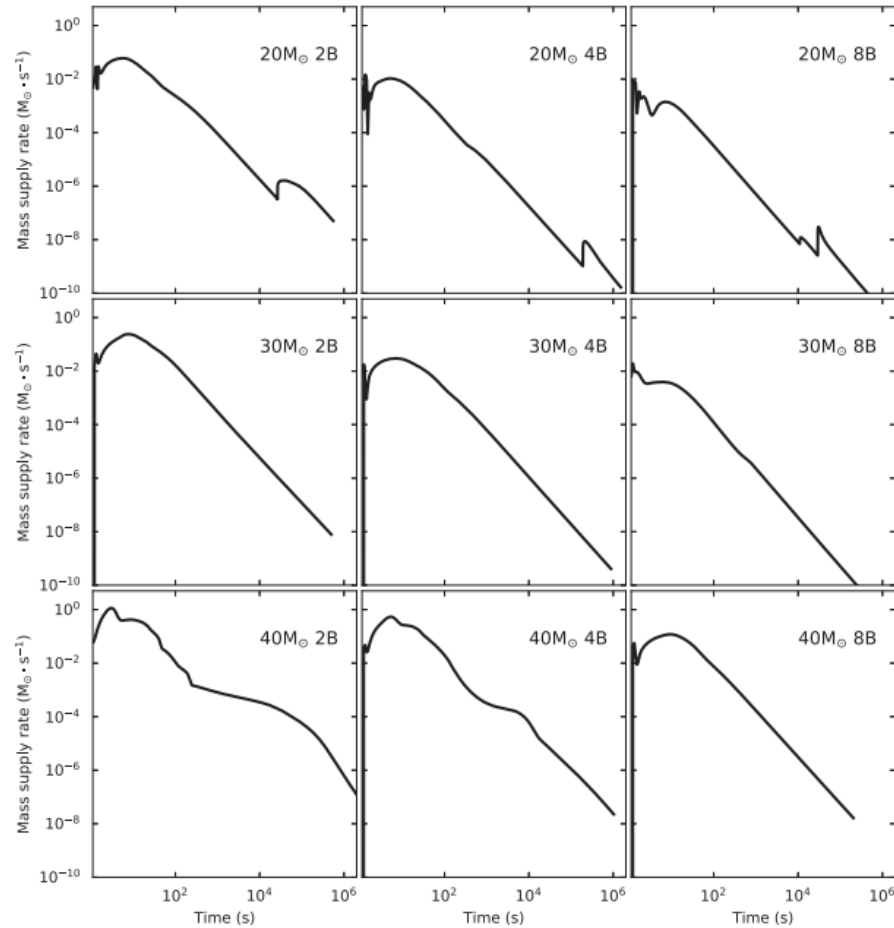
In the inner region of hyperaccretion disk  $T \sim 10^{10} - 10^{11} \text{K}$ ,  $\rho \sim 10^6 - 10^{13} \text{g cm}^{-3}$

- Neutrino physics (Urca processes; electron-positron pair annihilation;...)
- Gravitational wave emission (anisotropic neutrino radiation)



Liu et al. (2015)

# Mass supply rate



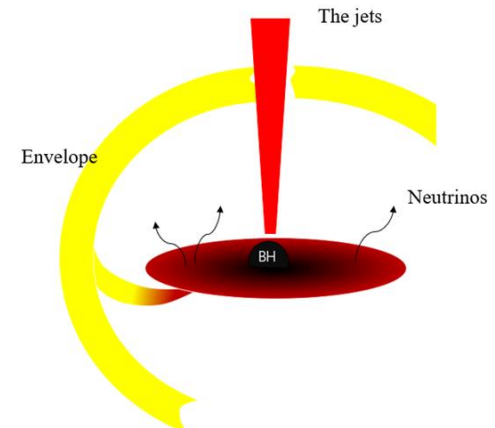
According to the high accretion rates, the NDAFs would be triggered.

# BH evolution

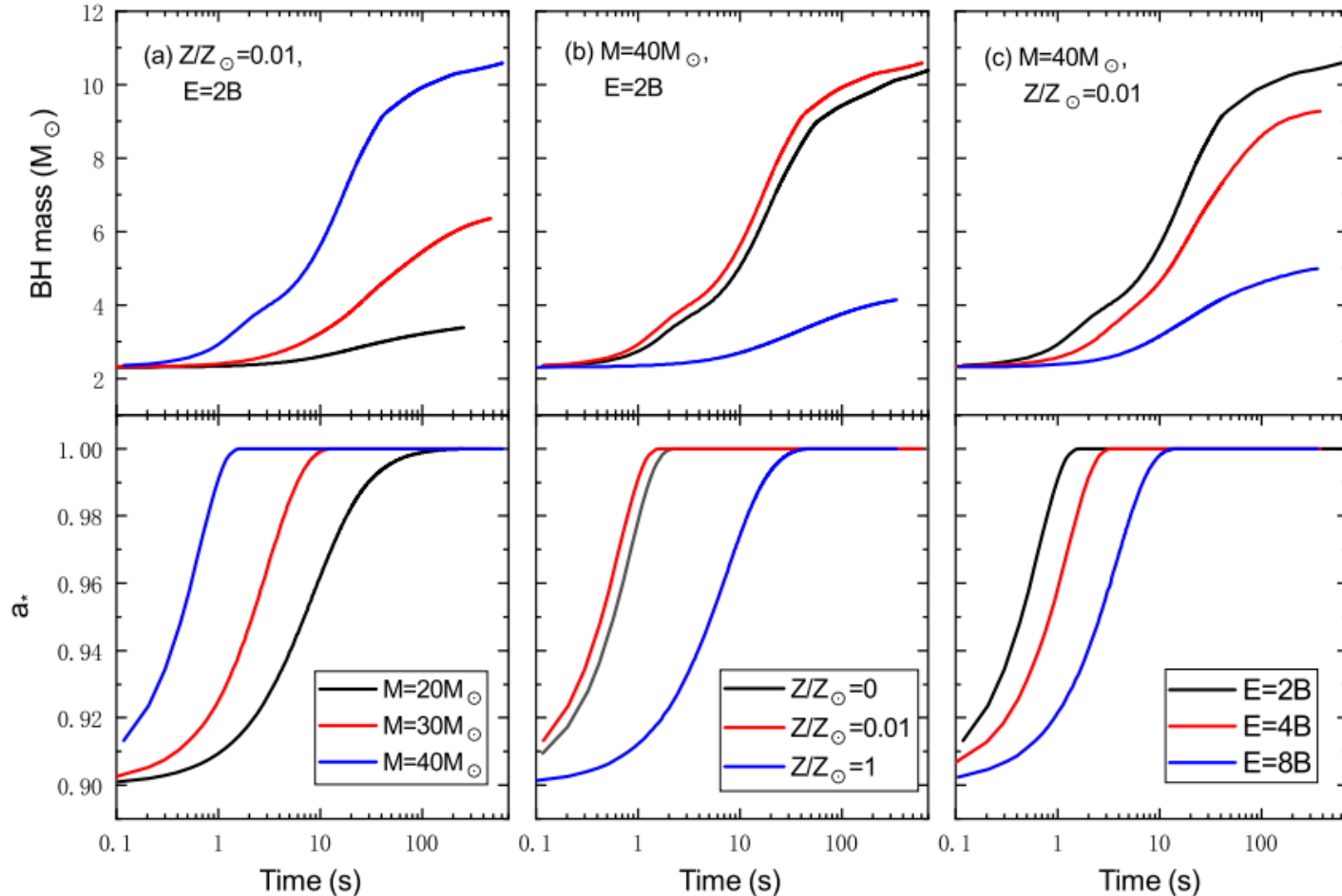
The mass and spin of a BH surrounded by a hyperaccretion disk will violently evolve within a short period.

$$\frac{dM_{\text{BH}}}{dt} = \dot{M} e_{\text{ms}}, \quad e_{\text{ms}} = \frac{1}{\sqrt{3x_{\text{ms}}}} \left( 4 - \frac{3a_*}{\sqrt{x_{\text{ms}}}} \right),$$
$$\frac{dJ_{\text{BH}}}{dt} = \dot{M} l_{\text{ms}}, \quad l_{\text{ms}} = \frac{2\sqrt{3}GM_{\text{BH}}}{c} \left( 1 - \frac{2a_*}{3\sqrt{x_{\text{ms}}}} \right),$$

$$\frac{da_*}{dt} = \frac{2\sqrt{3}\dot{M}}{M_{\text{BH}}} \left( 1 - \frac{a_*}{\sqrt{x_{\text{ms}}}} \right)^2$$



# Evolutions of mass and spin of BHs



The initial BH mass and spin are set as  $M_{\text{BH}} = 2.3 M_{\odot}$  and  $a = 0.9$ , respectively.

# Neutrinos from NDAFs

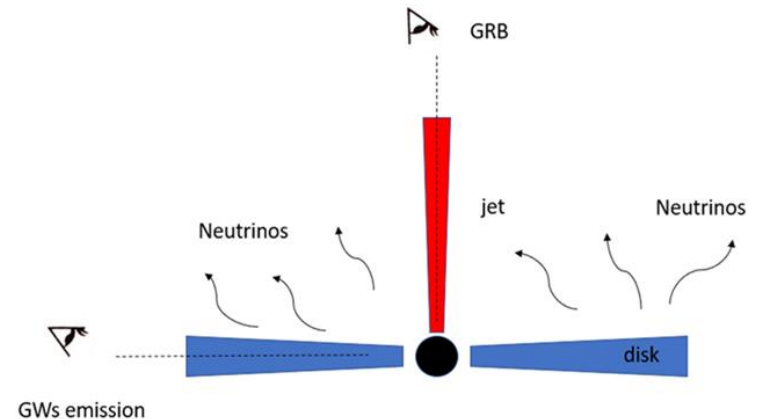
$$\log Q_{\bar{\nu}_e} = 41.40 - 0.23m_{\text{BH}} + 0.58a_* + 1.85 \log \dot{m} \\ - 3.96 \log r,$$

$$\log T = 11.23 - 0.4m_{\text{BH}} + 0.10a_* + 0.23 \log \dot{m} \\ - 0.86 \log r,$$

$$\log L_\nu = 52.80 - 0.03m_{\text{BH}} + 1.01a_* + 1.08 \log \dot{m},$$

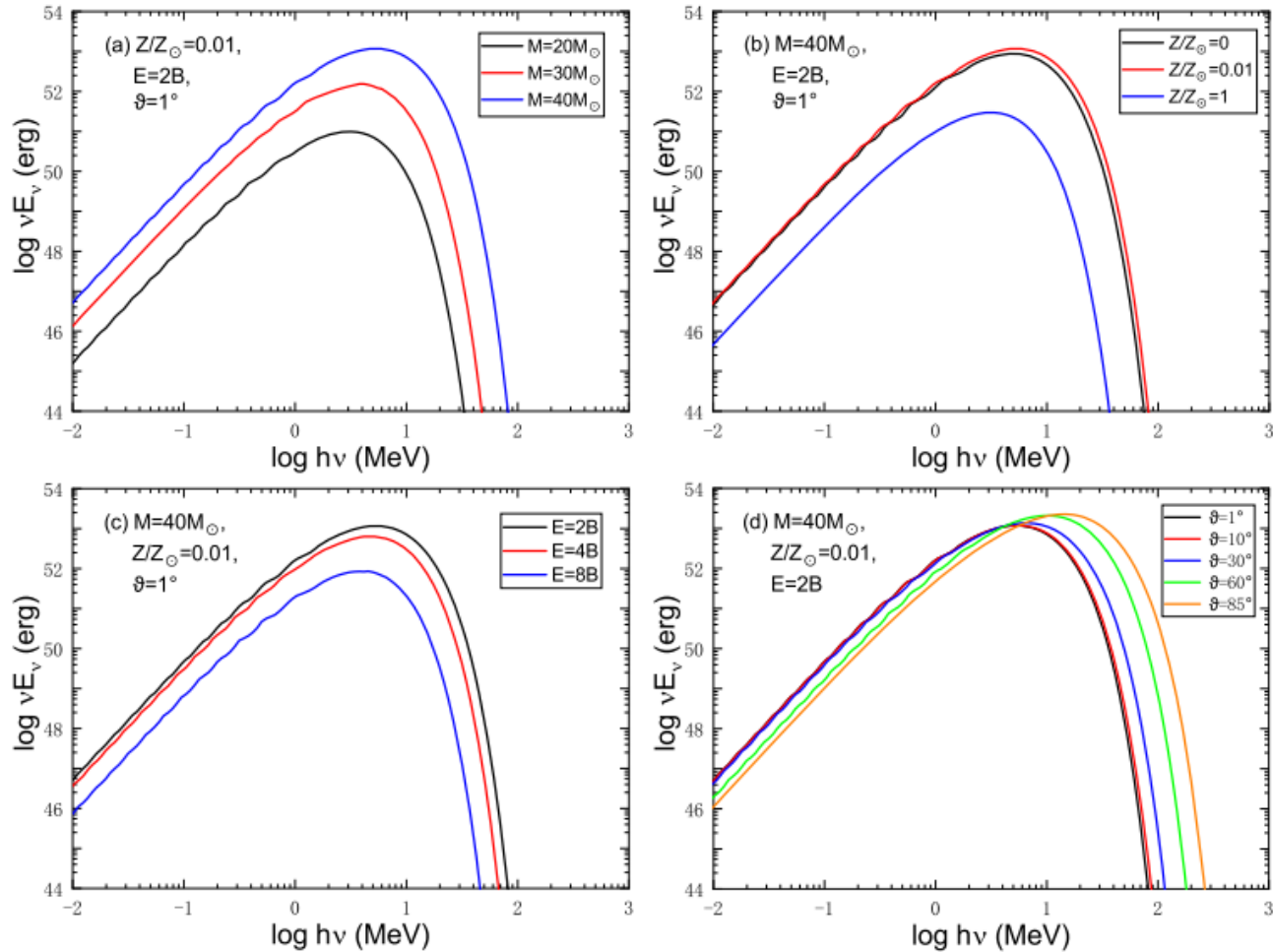
$$F_{E_{\text{obs}}} = \int_{\text{image}} g^3 I_{E_{\text{em}}} d\Omega_{\text{obs}},$$

$$I_{E_{\text{em}}} = Q_{\bar{\nu}_e} \frac{F_{E_{\text{em}}}}{\int F_{E_{\text{em}}} dE_{\text{em}}},$$

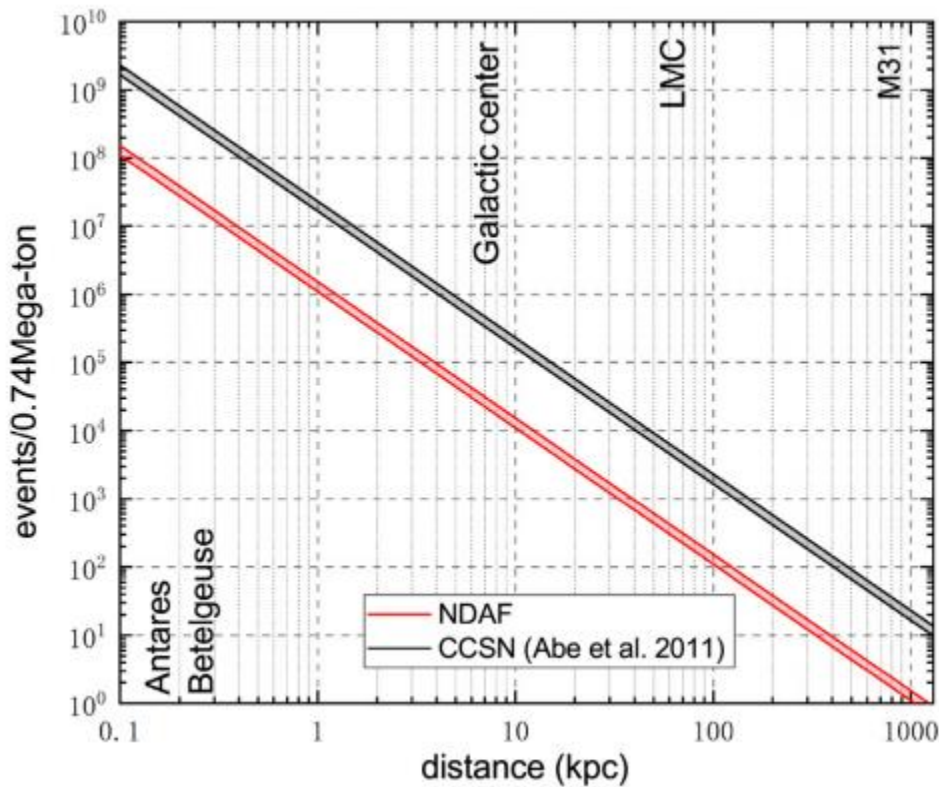




# Time-integrated electron antineutrino spectra of NDAFs



# MeV neutrino detection



Electron antineutrinos can be detected via inverse beta decay (IBD) reaction by the upcoming Hyper-Kamiokande (Hyper-K) detector



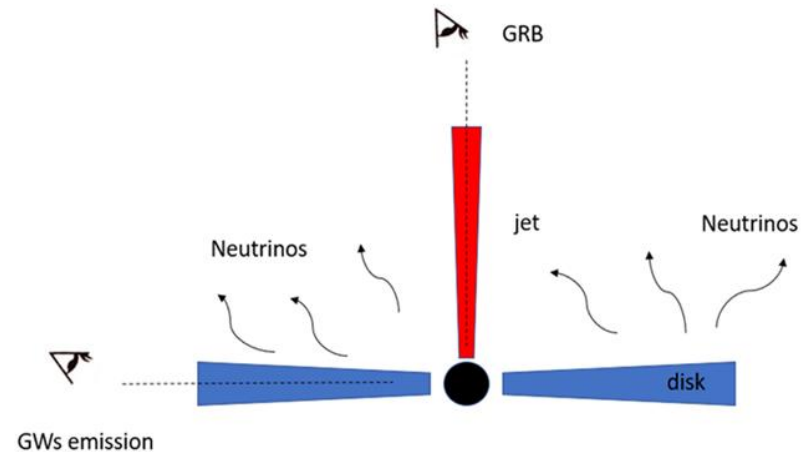
Approximately 165 000–230 000 inverse beta events are expected to be detected by Hyper-K for a typical CCSN at a distance of 10 kpc.

# GWs from NADFs

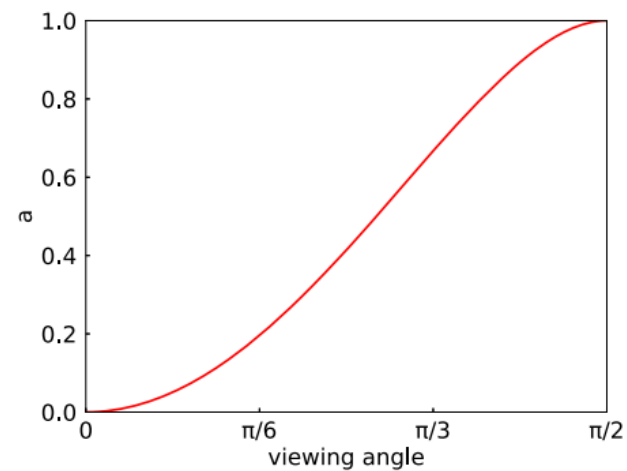
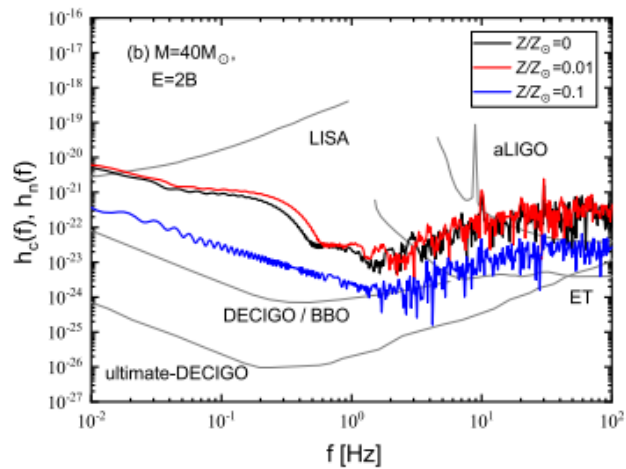
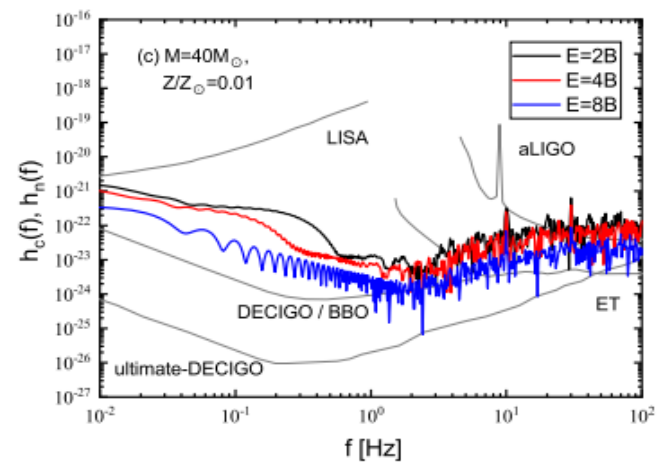
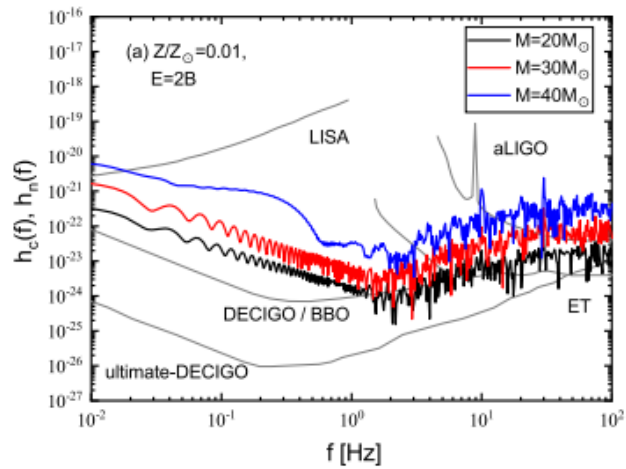
$$\log L_\nu = 52.80 - 0.03m_{\text{BH}} + 1.01a_* + 1.08 \log \dot{m}$$

$$h_{+, \text{eq}}^{TT}(t) = \frac{2G}{c^4 D} \int_{-\infty}^{t-D/c} \alpha(t') L_\nu(t') dt',$$

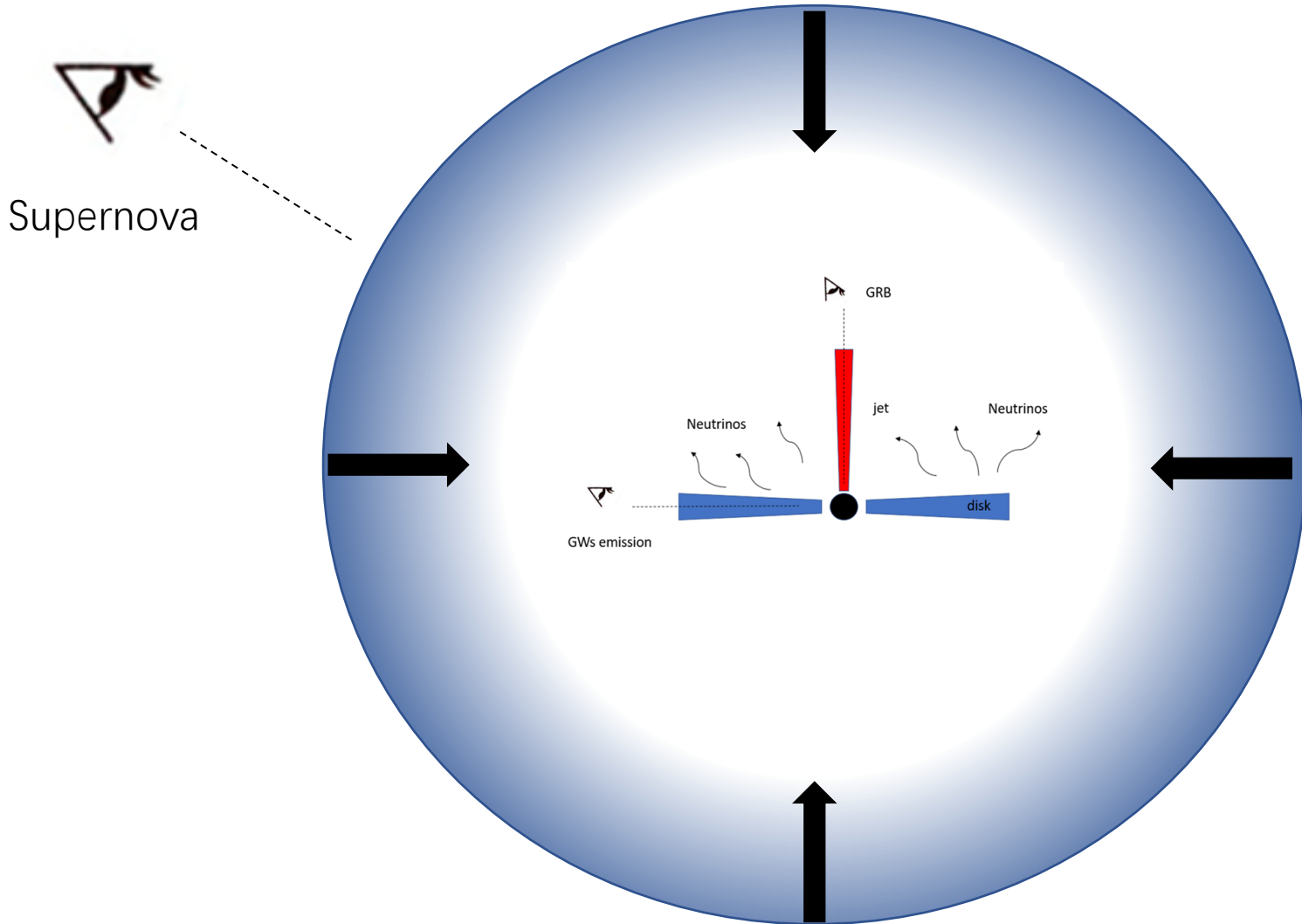
$$\alpha(t) = \frac{1}{L_\nu(t)} \int_{4\pi} \Psi(\vartheta', \varphi') \frac{dL_\nu(\vec{\Omega}', t)}{d\Omega'} d\Omega'$$



# Characteristic amplitudes of GWs from NDAFs



# Viewing angle



# Summary

- If the typical initial explosion energy of the CCSNe is much **lower than that of the observable ones**, the lower mass gap will naturally exist.
- The effects of the **initial explosion energies, viewing angles, and masses and metallicities of progenitor stars** on time-integrated spectra of neutrinos and GW spectra are investigated.
- According to the current CCSN observations, **faint or failed CCSNe with low initial explosion energy might be universal**, which is beneficial to the detection of the neutrino emission and GW radiation of NDAFs.
- The **joint detections of neutrinos and GWs**, including the observations of LGRBs associated with CCSN, are meaningful and would help to verify the **existence of the central BH hyperaccretion disc**.

Thanks