

ACCELERATING EARTH-BOUND DARK MATTER

DM+NU FORUM AT TDLI & SJTU

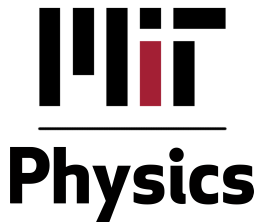
MARIANNE MOORE

TRIUMF & UBC & MIT

2202.08840

WITH DAVE McKEEN, DAVID MORRISSEY, MAXIM POSPELOV & HARI RAMANI

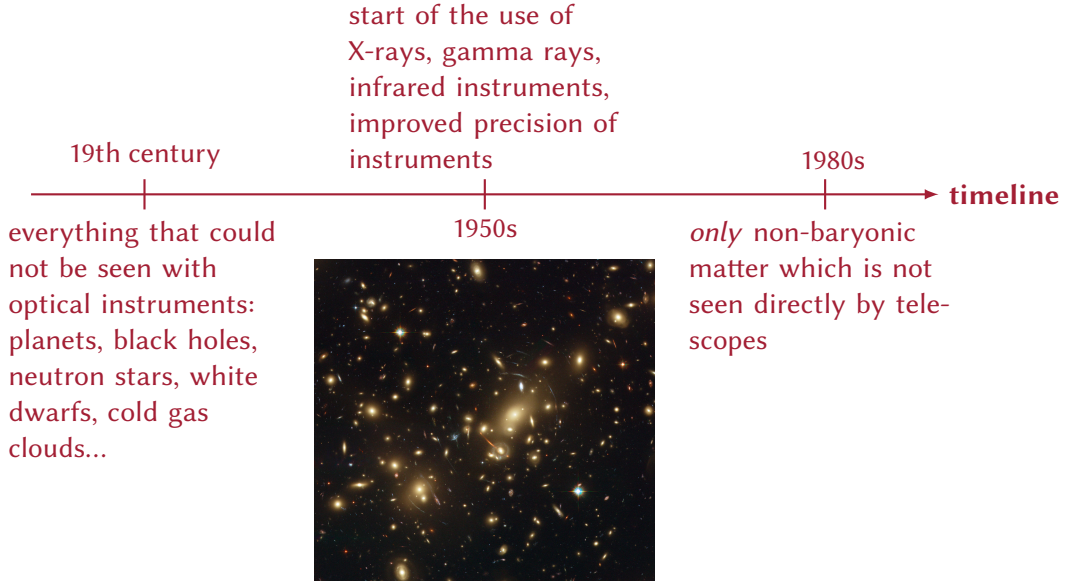
MARCH 23, 2022



- Dark matter & strongly interacting dark matter
- Dark matter accumulation
- Dark matter acceleration and detection
 - ▶ Underground nuclear accelerators
 - ▶ Thermal sources
- Summary

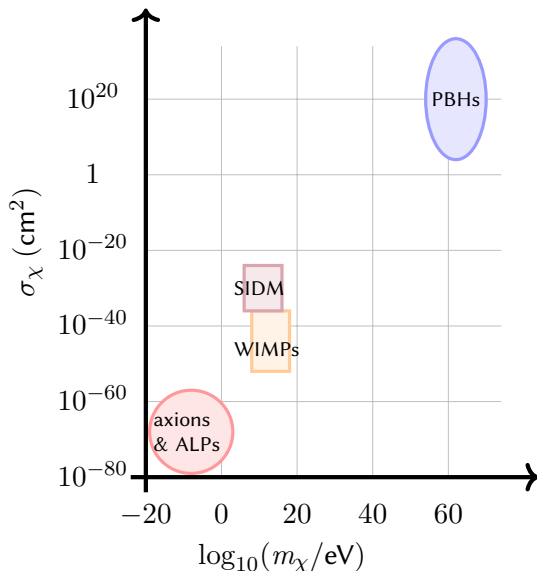
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WHAT IS DARK MATTER?



NASA, ESA, and Johan Richard

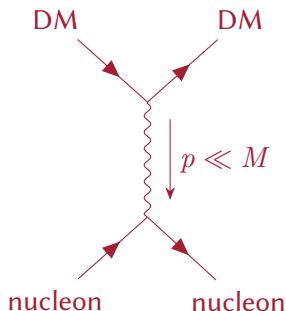
WHERE IS STRONGLY INTERACTING DARK MATTER?



$$Q = 0$$

collision time $\in [10^{10}, 10^{22}]$ years in
the interstellar medium

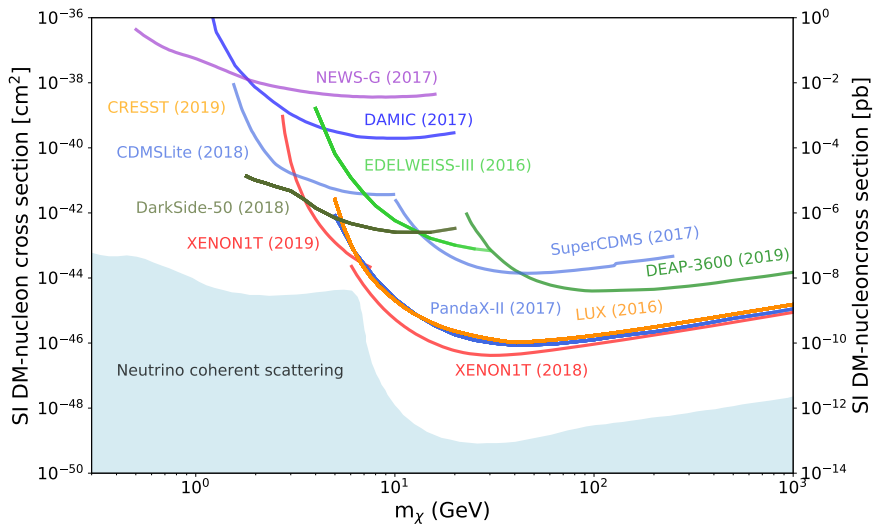
Spin-independent contact interaction



no annihilation or very little at the present epoch ($\bar{\chi}\chi \not\rightarrow NN$)

Similar analysis done with a dark photon ($m_{A'} = 15, 50 \text{ MeV}$).

BOUNDS ON WEAKLY INTERACTING MASSIVE PARTICLES



there are also bounds on DM-electron scattering, SD-proton and SD-neutron scattering

PARTICLE DETECTORS ARE BLIND TO STRONGLY INTERACTING DARK MATTER

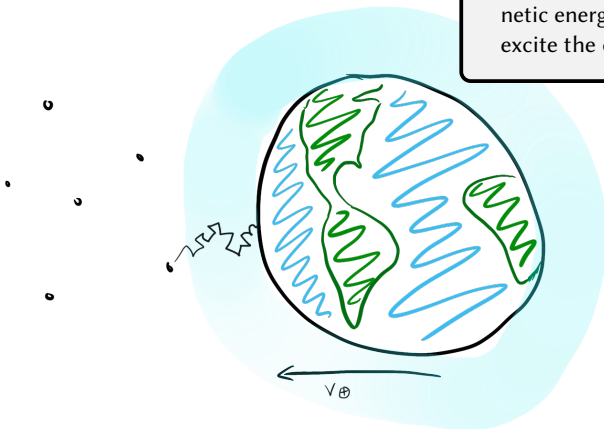
Scatters many times in the atmosphere and slows down

velocity $\sim \sqrt{\frac{kT}{m_\chi}}$, kinetic energy ~ 0.03 eV

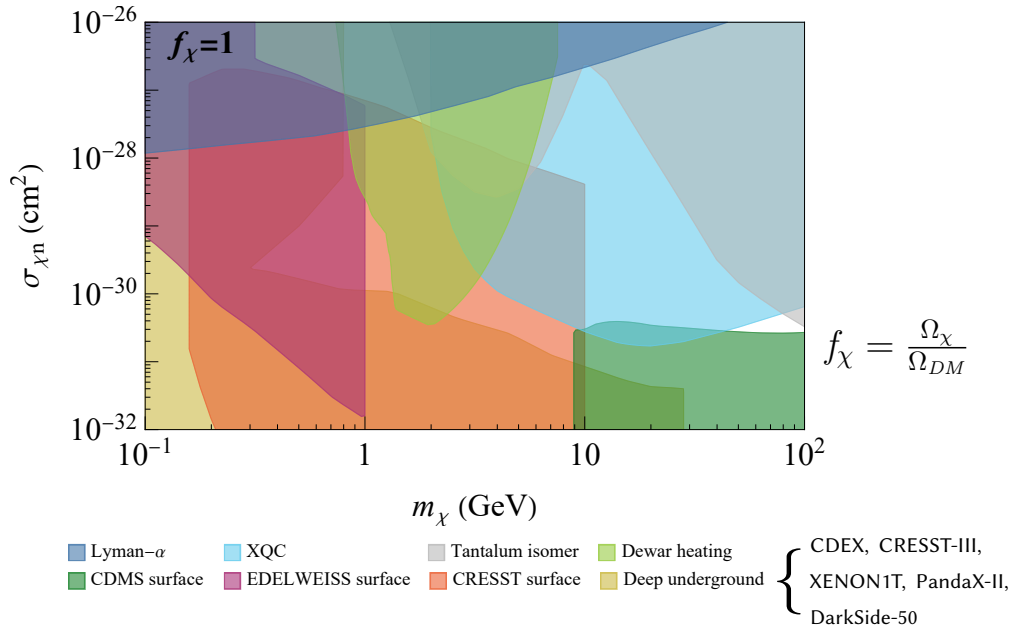
Current nuclear detection thresholds: $\mathcal{O}(\text{keV})$

Goodman & Witten

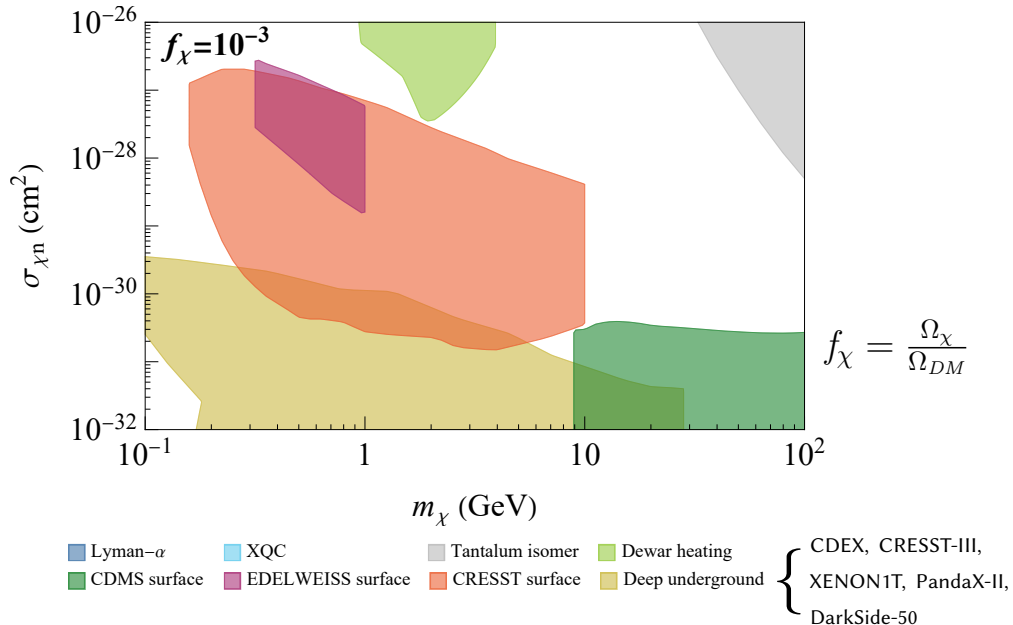
For values of $M_\chi \lesssim 100$ GeV, the stopping power of the atmosphere becomes important, and the χ particle probably could not be detected at sea level. Their kinetic energy would be too low to excite the detector [...] (1985)



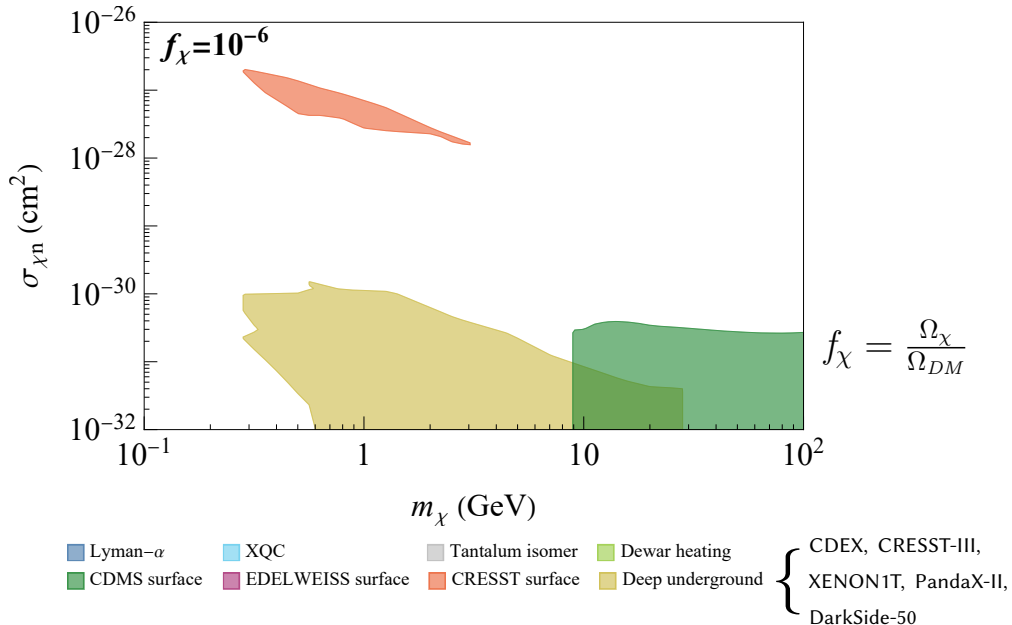
STRONGLY INTERACTING DARK MATTER IS COMPLETELY RULED OUT...



..., BUT NOT IF IT IS A SUB-COMPONENT!



..., BUT NOT IF IT IS A SUB-COMPONENT!



Should we care about strongly interacting dark matter?

Probably!

How can we detect it?

- | | | |
|--|---------------|---|
| ■ Too sparse ($f_\chi \sim 10^{-6}$, $\rho \sim 0.3$
GeV/cm ³ , $m_\chi \sim 1$ GeV) | $\Rightarrow$ | ■ accumulate in the Earth over its
lifetime |
| ■ Too slow (not energetic enough
to be detected) | | ■ accelerate with nuclear accelerators
and thermal sources |

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THERMAL POPULATION OF DARK MATTER IN THE EARTH

$$\frac{dN_{\text{DM}}}{dt} = \frac{\rho_{\text{DM}}}{m_{\text{DM}}} v_{\oplus} \pi R_{\oplus}^2 f_{\text{cap}} - N_{\text{DM}} f_{\text{loss}}$$

of DM intercepted by Earth

fraction of DM captured by gravity

$$f_{\text{cap}} = \frac{2(\ln[1 - \bar{f}_{\text{KE}}])^{1/2}}{(\pi \ln[v_{\text{es}}^2/v_{\oplus}^2])^{1/2}}$$

fraction of DM lost by Jeans escape

$$f_{\text{loss}} = \frac{3}{R_{\oplus} \bar{n}_{\text{DM}}} \frac{n_{\text{LSS}} v_{\text{LSS}}}{2\pi^{1/2}} \left(1 + \frac{v_{\text{es}}^2}{v_{\text{LSS}}^2}\right) \exp(-v_{\text{es}}^2/v_{\text{LSS}}^2)$$

gravity

evaporation

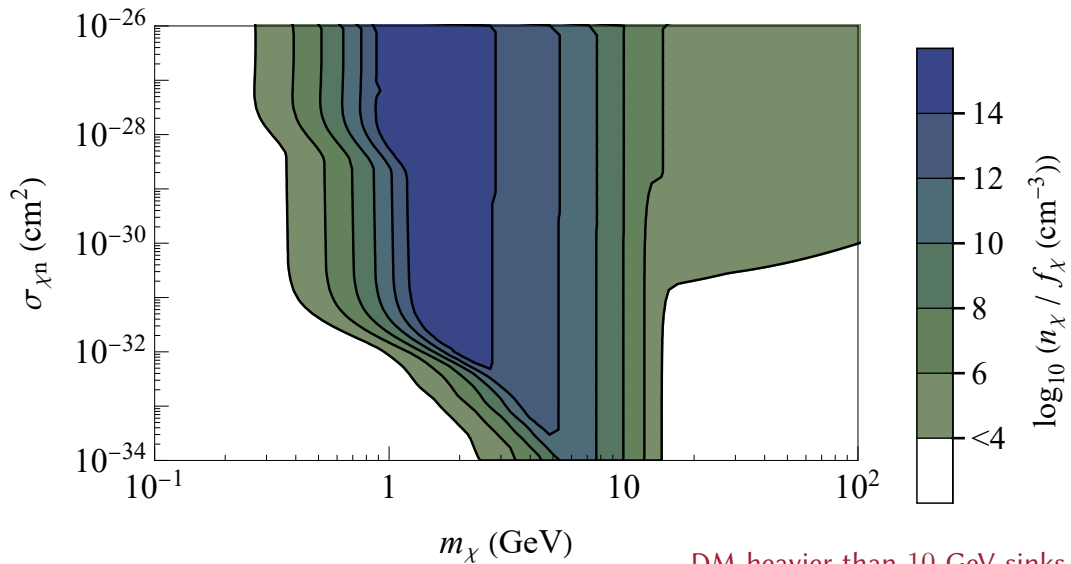
TRAFFIC JAM POPULATION OF DARK MATTER IN THE EARTH



infalling dark matter that has not yet reached the equilibrium distribution

$$v_z \frac{dv}{dz} \propto -n_N \sigma_{\chi n} \frac{v_N^3}{v}$$

A LARGE DENSITY NEAR THE EARTH SURFACE CAN BE OBTAINED



DM heavier than 10 GeV sinks,
lighter than 1 GeV evaporates

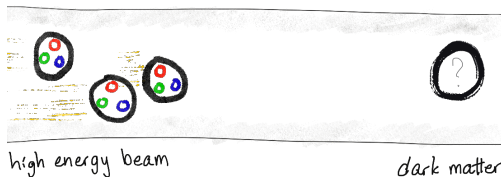
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ACCELERATING THERMALIZED DARK MATTER WITH BEAMS AND HEAT

We need to give energy to the dark matter:

Beams

1. Particle accelerators (~ 1 TeV)
2. Nuclear accelerators (~ 1 MeV)

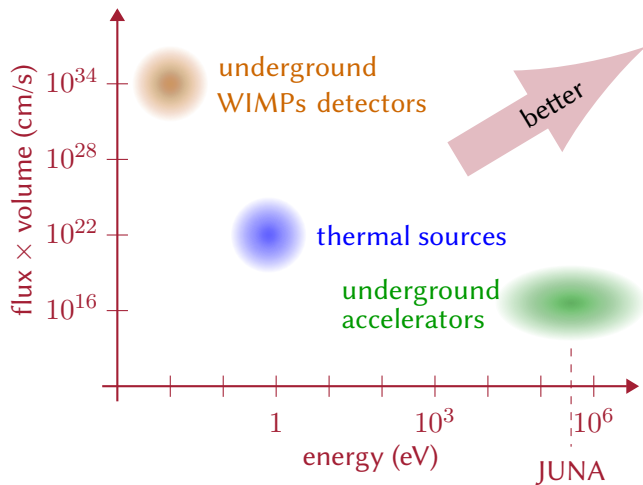


Heat

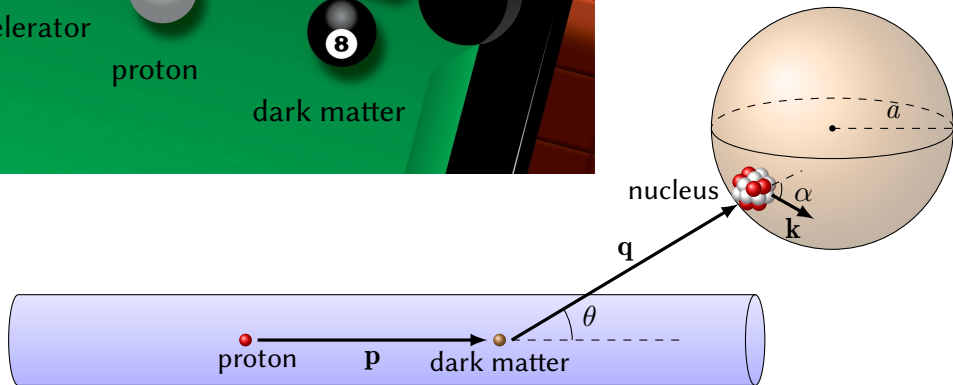
1. Light bulbs (~ 3000 K ~ 0.3 eV)
2. Tokamak ($\sim 10^8$ K ~ 10 keV)



OVERVIEW OF ACCELERATING TECHNIQUES

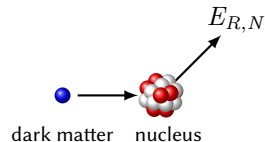
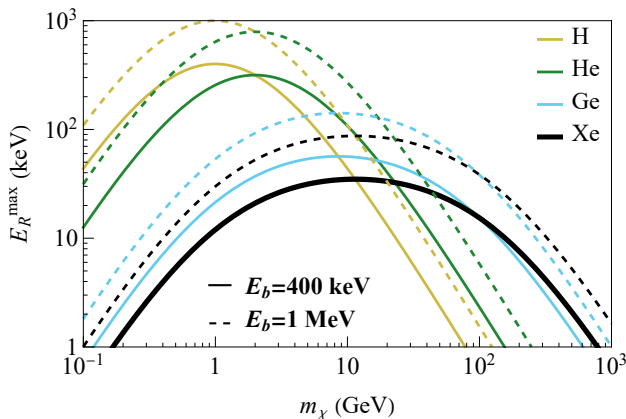


PLAYING POOL WITH DARK MATTER



A NUCLEAR RECOIL IS ACHIEVABLE WITH NUCLEAR ACCELERATORS

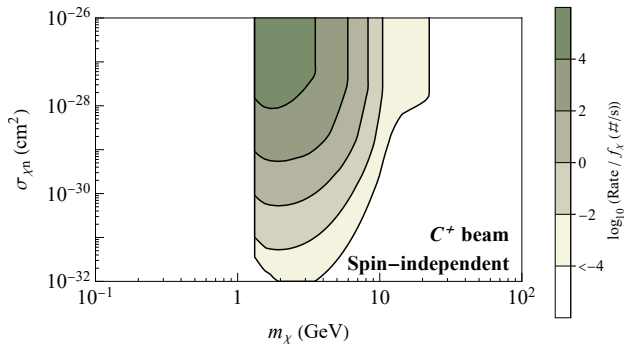
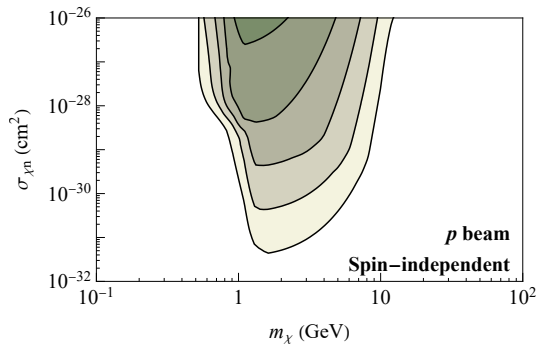
	LUNA-MV	JUNA	CASPAR
voltage (kV)	3500	400	1100
H current (mA)	1	10	0.25
He current (mA)	0.5	2	0.22
C current (mA)	0.15	—	—



THE RATE IS MEASURABLE IN A 10CM DIAMETER XENON DETECTOR

$$E_{\text{th}} = 5 \text{ keV}$$

$$R \sim n_{\chi} \times \underbrace{\left(\frac{I_b}{Q_b} L \sigma_{\chi b} \right)}_{\sim 10^{17} \text{ cm/s}} \times (\ell \sigma_{\chi N} n_N) P_{\text{sh}} P_{\text{th}}$$



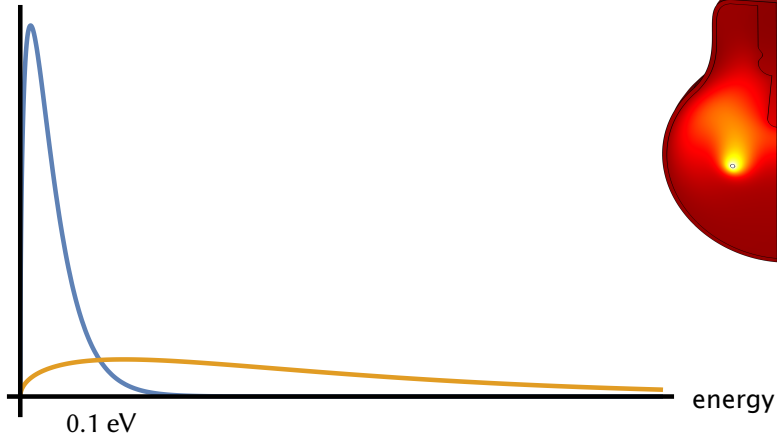
THE INCANDESCENT LIGHT BULB SETUP

Dark matter gas @ 300 K

Argon gas @ 3000 K

$$P(E) dE = \sqrt{\frac{E}{\pi T^3}} \exp\left(-\frac{E}{T}\right) dE$$

probability

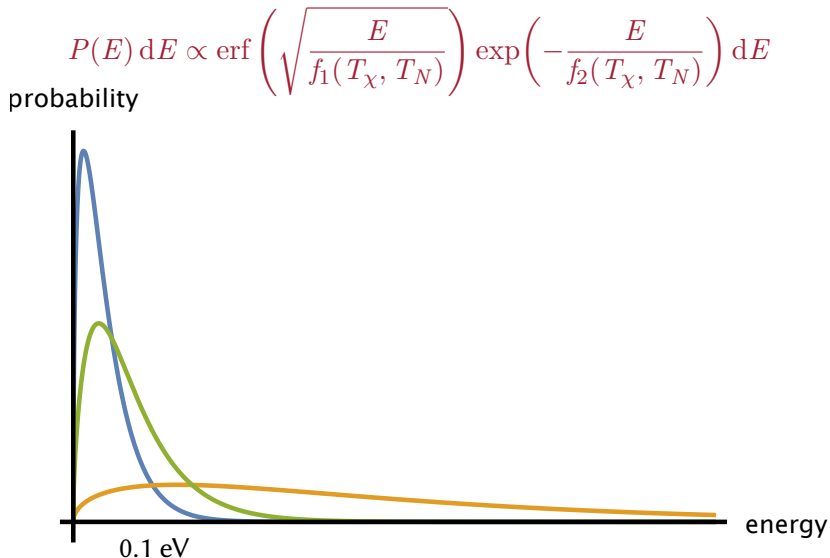


ENERGY ACHIEVABLE BY UP-SCATTERED DARK MATTER

Dark matter gas @ 300 K

Argon gas @ 3000 K

Dark matter gas scattering

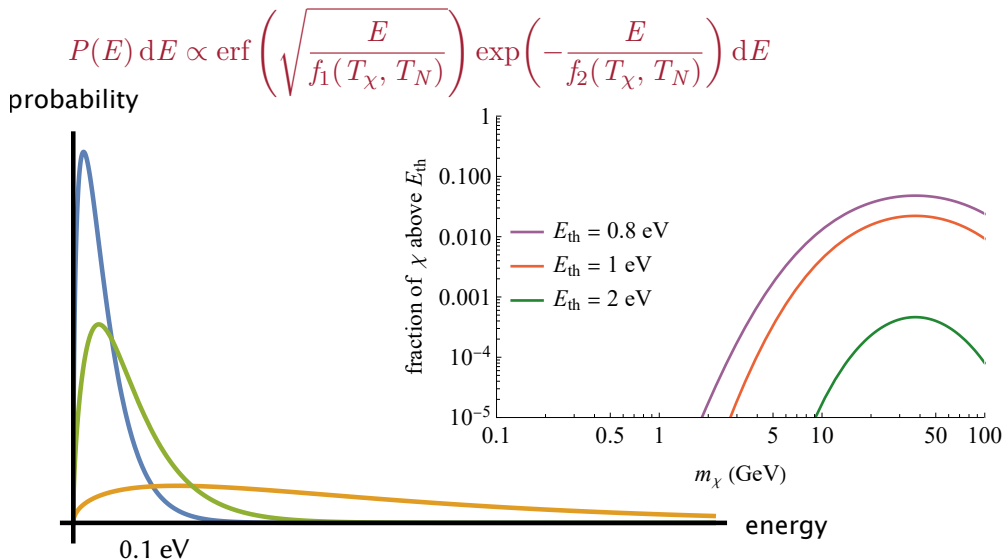


ENERGY ACHIEVABLE BY UP-SCATTERED DARK MATTER

Dark matter gas @ 300 K

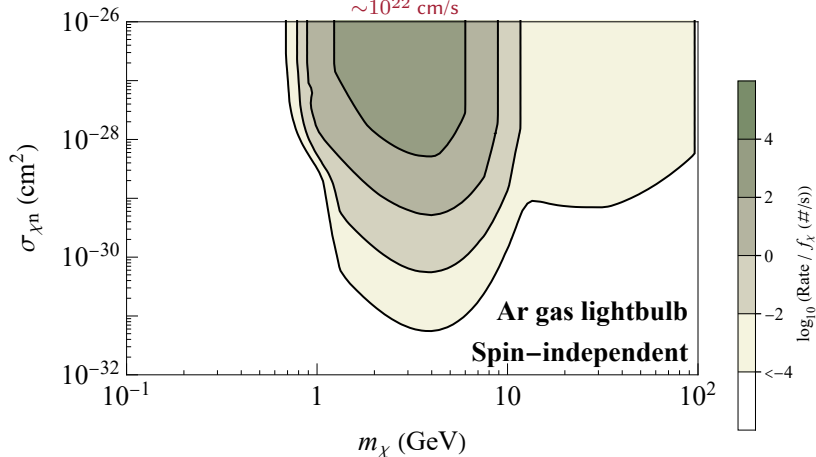
Argon gas @ 3000 K

Dark matter gas scattering

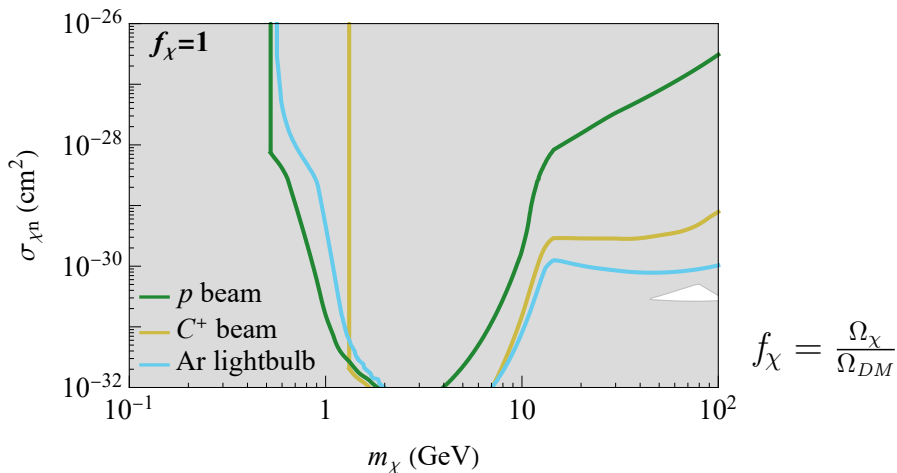


THE RATE IS MEASURABLE IN A 1G CO DETECTOR

$$E_{\text{th}} = 0.8 \text{ eV} \quad R \sim n_{\chi} \underbrace{\left(N_g \sqrt{\frac{T_g}{m_g}} \sigma_{\chi g} \right)}_{\sim 10^{22} \text{ cm/s}} \times \left(\frac{1}{4\pi d^2} \sigma_{\chi N} N_N \right) P_{\text{sh}} P_{\text{th}}$$

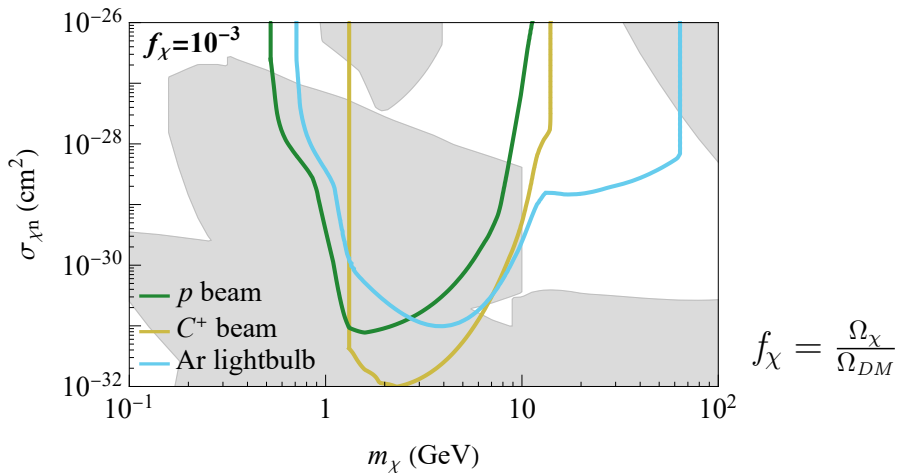


RATES IN A DETECTOR FOR $f_\chi = 1$



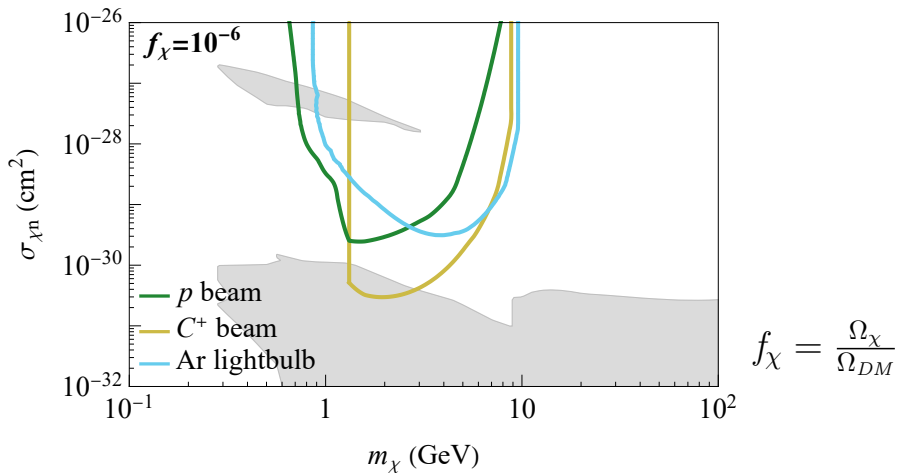
10 events per year, rate $\propto \sigma_{\chi n}^2$

RATES IN A DETECTOR FOR $f_\chi = 10^{-3}$



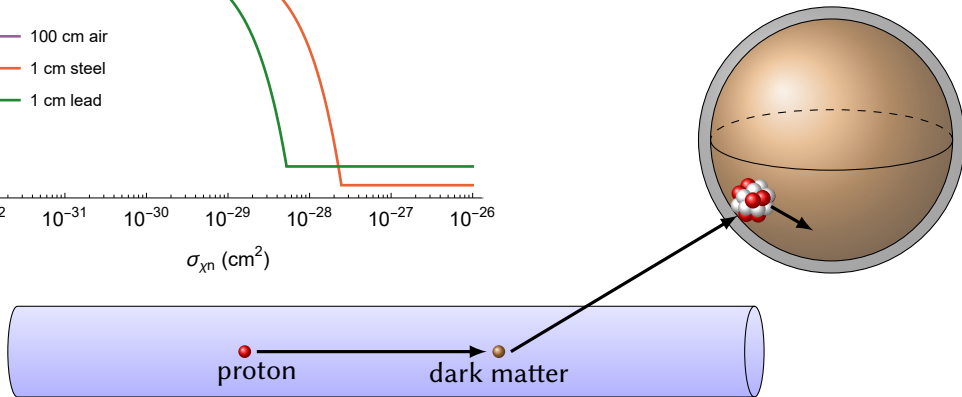
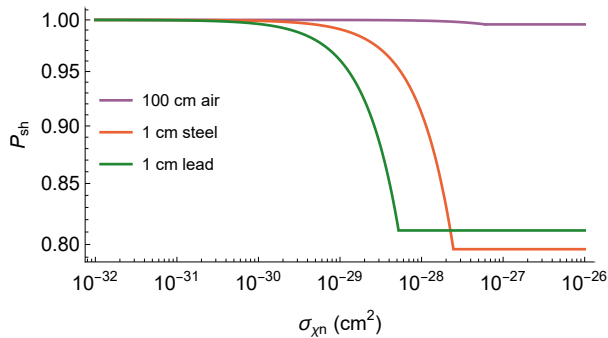
10 events per year, rate $\propto \sigma_{\chi n}^2$

RATES IN A DETECTOR FOR $f_\chi = 10^{-6}$

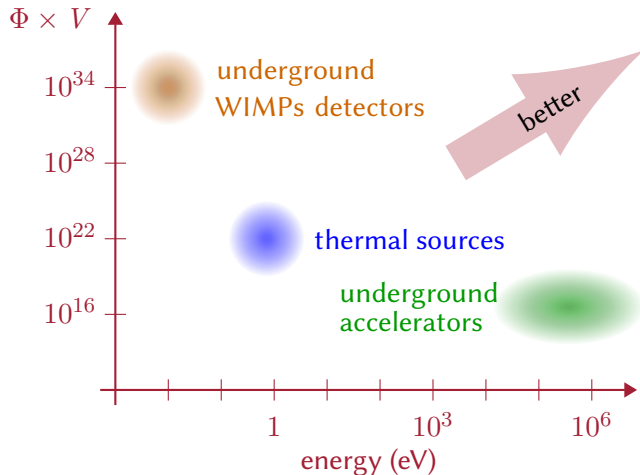


10 events per year, rate $\propto \sigma_{\chi n}^2$

$$P_{\text{sh}} \sim P_{\text{sh, air}} \times P_{\text{sh, lead}} \times \dots$$



elements to improve	effect on rate in detector
shielding analysis	↓
collimation of the beam	↑
V and T of gas	↑



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SUMMARY

Strongly interacting dark matter:

- Could exist as a fraction of the total dark matter ✓
- Direct detection methods cannot probe sub-components ✕
- Interacts with protons and neutrons ✓
- Overdense ✓
- Very slow ✕

Accelerate it through up-scattering with nuclear accelerators (possible with current detector thresholds) or thermal sources, such as light bulbs (possible in the near future...?)



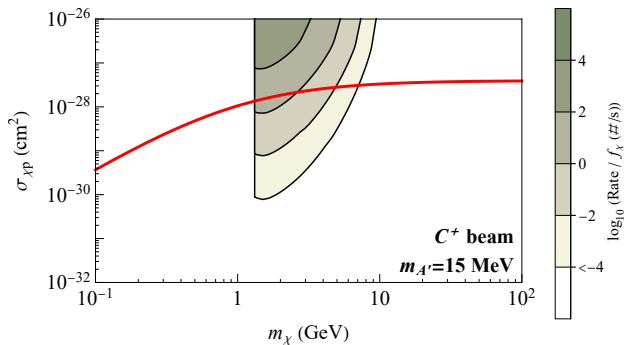
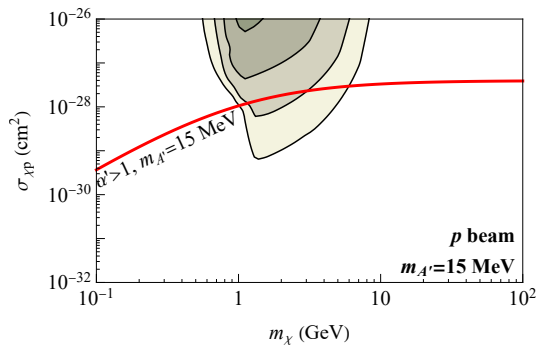
millicharge dark matter:

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APPENDIX

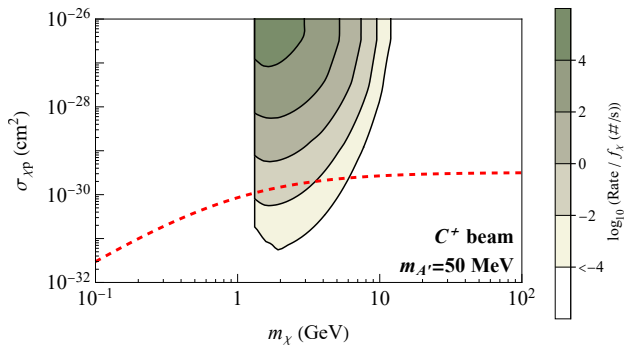
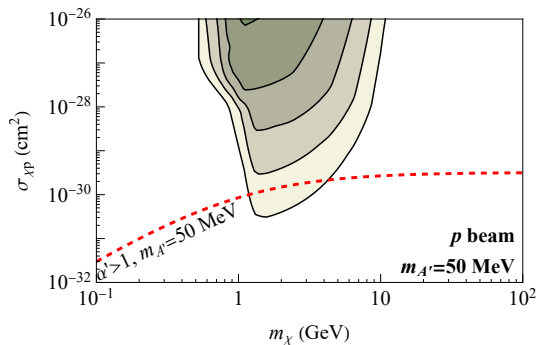
15 MeV DARK PHOTON

need $\alpha' \leq 1$ to maintain perturbative consistency



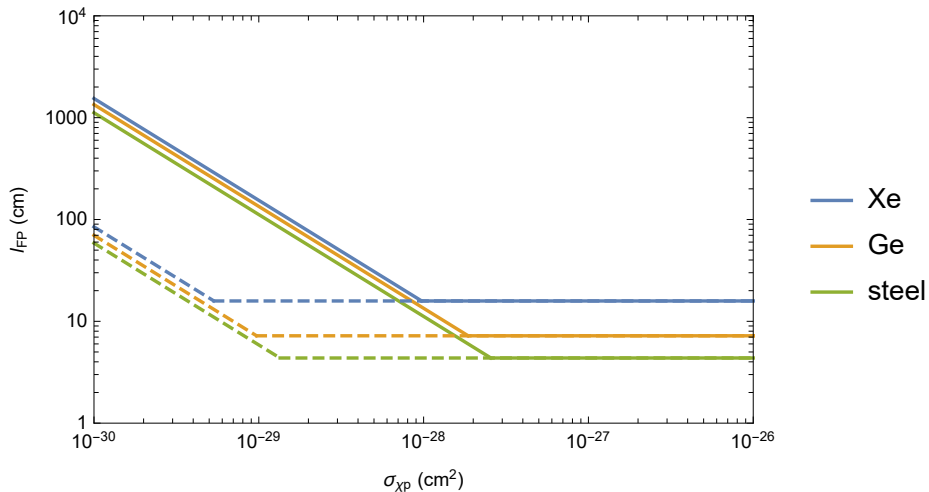
50 MeV DARK PHOTON

need $\alpha' \leq 1$ to maintain perturbative consistency



MEAN FREE PATH OF DARK MATTER IN MATERIALS

$l_{FP} = (n_T \sigma_{\chi p})^{-1}$, saturates at the geometric cross-section $4\pi(1.2 \text{ fm } A^{1/3})^2$



(for a particle of mass 1 GeV (full) and 10 GeV (dashed) up-scattered by a 400 keV beam)